

Notes: b-Contact Structures on Symplectic Hyperboloids

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The McGehee Transformation

If we have a vector field on $M := T^*\mathbb{R}^+ \times T^*\mathbb{S}^{n-1}$, then we can look at the push-forward via τ_{McG}^{-1} of this to the two components of $X \setminus Z$ and try to extend it to Z to get a b^3 -vector field.

Q: Does this push forward always extend to a b^3 -vector field on X ?

A: No. (for example the radial vector field)

Q: Does the push forward of a Liouville vector field extend to a b^3 -Liouville vector field on X ?

A: Yes, assuming it extends to a b^3 -vector field in the first place.

If we have a hypersurface $S \subset M$ which is closed as a subset. Then we can define:

$$S_{\text{McG}} := \overline{\tau_{\text{McG}}^{-1}(S)} \subset X$$

(Be aware that $\tau_{\text{McG}}^{-1}(S)$ is by definition a subset of $X \setminus Z$, but we take it's closure in X .)

Q: Is this extension always a hypersurface in X ?

A: No.

Q: Is the extension of a contact type hypersurface b^3 -contact type?

A: No, even if it is a hypersurface in X the b^3 -contact type property can still fail.

Theorem

From now on we assume that S is a hypersurface in $T^*\mathbb{R}^n$ given by $S = H^{-1}(0)$ and H is given by: $H(q, p) = q^t B p - 1$ where B is a (non-singular) $n \times n$ -matrix.

Then S is of contact type, since all non-degenerate quadratics are transverse to the radial vector field.

Next, we can look at it as a hypersurface in M and thus at S_{McG} as we've seen before.

Thm: If $B + B^t$ is positive definite, then S_{McG} is a (smooth) hypersurface in X and of b^3 -contact type.

Proof:

Fix $\phi: O \subset \mathbb{S}^{n-1} \subset \mathbb{R}^n \rightarrow U \subset \mathbb{R}^{n-1}$ an arbitrary chart and we denote its inverse by ψ .

Then we can write the splitting as $q = r \cdot \psi$.

Since we want to preserve the structure of the cotangent bundle, we get an induced map on the coordinates on the cotangent fiber. Thus we have:

$$p = \frac{\partial r}{\partial q} P_r + \frac{\partial \phi}{\partial q} \eta$$

Now one can simply calculate: $\frac{\partial r}{\partial q_i} = \frac{q_i}{r} = \psi_i$. The matrix $\frac{\partial \phi}{\partial q}$ we can of course not calculate explicitly,

because ϕ is arbitrary, but we can always say that it is $\frac{1}{r} \Psi$ for some matrix Ψ , which is independent of r . Thus we have the formula:

$$p = \psi P_r + \frac{1}{r} \Psi \eta$$

Now we can express H in terms of the n-spherical coordinates, this gives us:

