

Shape and Topology optimization applied to Compact Heat Exchangers

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Von Karmann Institute

Compact Heat Exchangers in Additive Manufacturing – 2021, April 27th



ETH zürich

Topology optimization

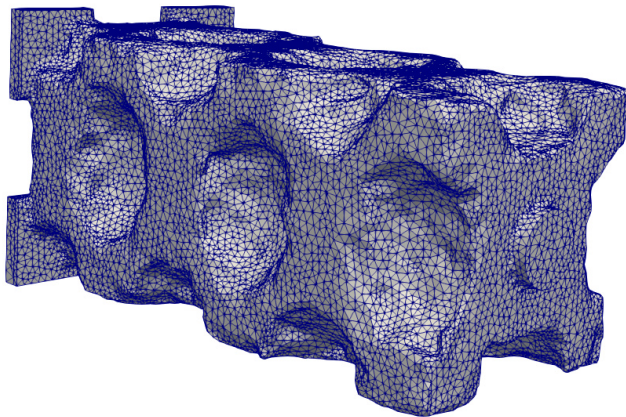


Figure: Optimization of the rigidity of a mechanical structure subject to flexural load

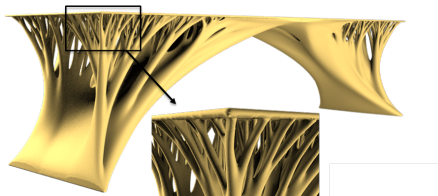
State of the art for 3D topology optimization



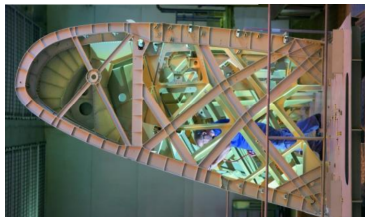
(a) Siemens (2017)



(b) APWorks (2016)



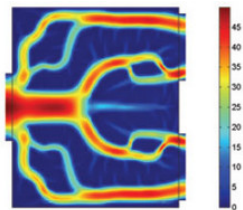
(c) M2DO (Kambampati et. al. 2018)



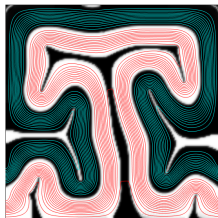
(d) AIRBUS (2010)

State of the art for 3D topology optimization

For thermal-fluid systems it is still an active research field.



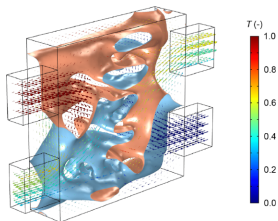
(a) Dede (2009, Toyota)



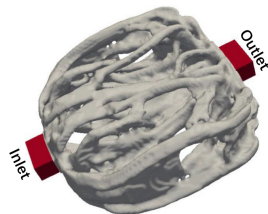
(b) Papazoglou (2015, TU Delft)

Figure: Fluid pipes optimized for convective heat transfer.

State of the art for 3D topology optimization



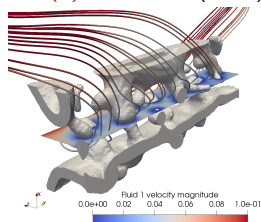
(a) Kobayashi et. al. (2020)



(b) Yu et. al. (2020)



(c) Savier (2019, United Technologies)



(d) Hoghoj et. al. (2020)

Figure: Fluid pipes optimized for convective heat transfer with density methods.

State of the art for 3D topology optimization

The objective today: shape and topologically optimized heat exchangers with the method of Hadamard and body-fitted meshes.

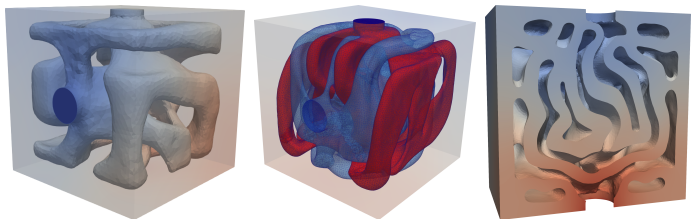
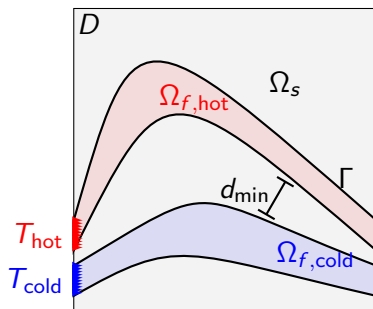


Figure: Topology optimized heat exchanger devices with the method of Hadamard and a body-fitted mesh evolution algorithm. Figures from ¹.

¹Feppon et al., *Body-fitted topology optimization of 2D and 3D fluid-to-fluid heat exchangers* (2021)

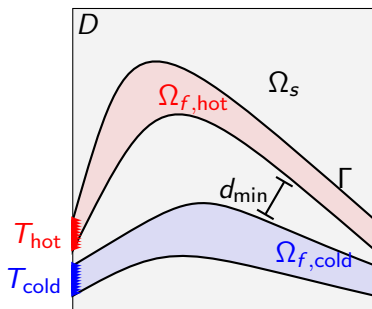
Outline

1. Formulation of the optimal heat exchanger design problem
 - ▶ Physical modelling
 - ▶ Non-mixing constraint
2. Shape and Topology optimization with the method of Hadamard
 - ▶ Shape derivatives
 - ▶ Treatment of geometric constraints
3. Numerical Topology optimization
 - ▶ Null space optimization algorithm
 - ▶ Body fitted mesh evolution method for numerical shape updates
4. Numerical Results



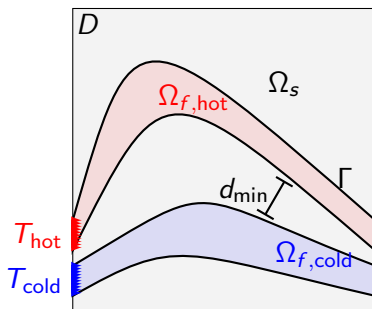
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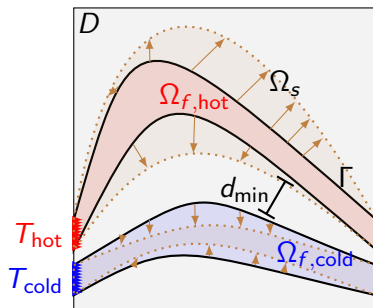
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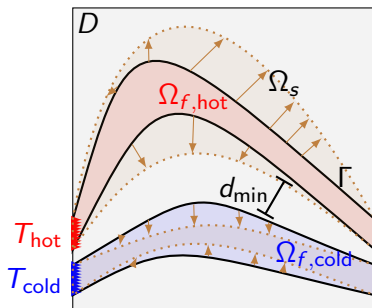
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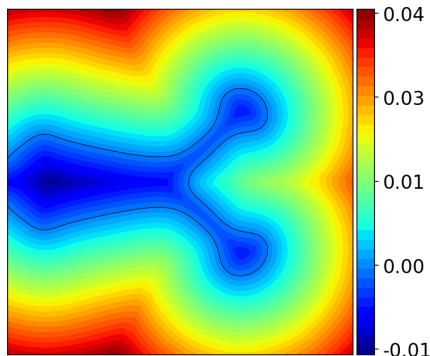
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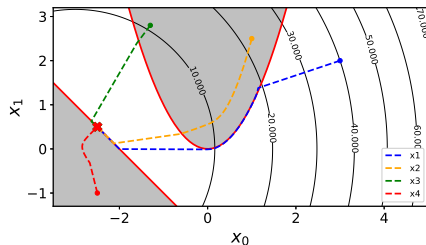
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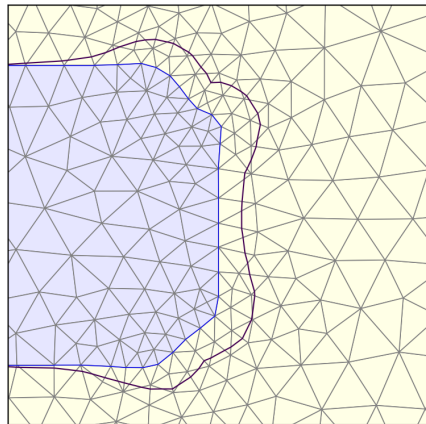
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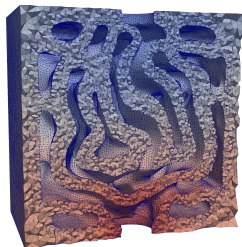
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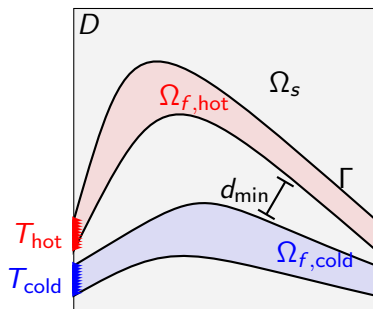
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Outline

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 - Physical modelling



Problem at hand

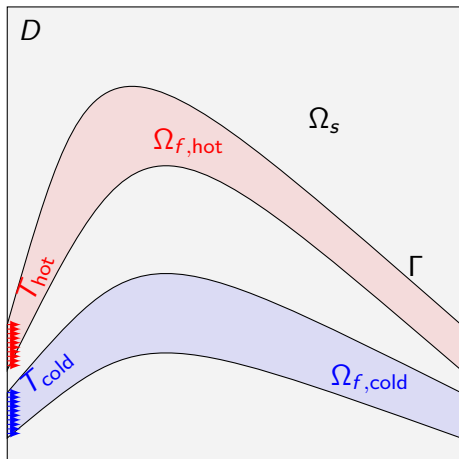
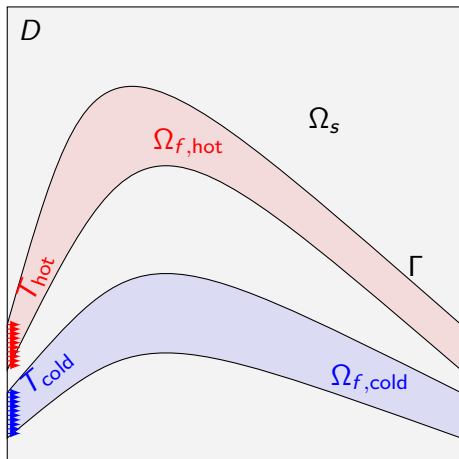


Figure: Settings of the heat exchanger topology optimization problem.

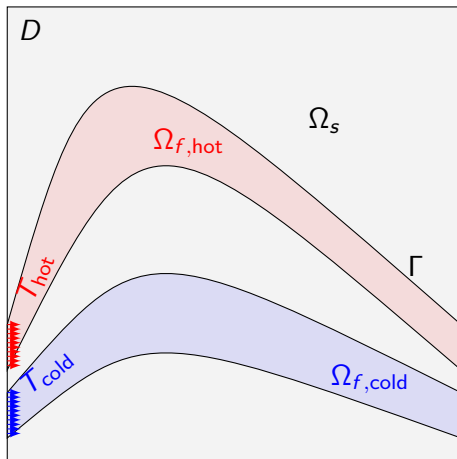
Problem at hand



- Navier-Stokes flows in the **hot** and **cold** phases $\Omega_{f,hot}$ and $\Omega_{f,cold}$.

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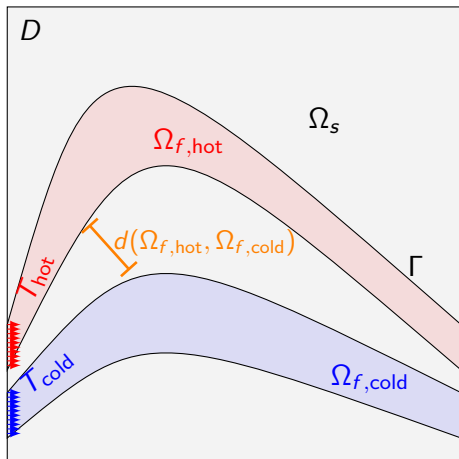
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- ▶ Navier-Stokes flows in the **hot** and **cold** phases $\Omega_{f,hot}$ and $\Omega_{f,cold}$.
- ▶ Thermal convection in the fluid phase $\Omega_f = \Omega_{f,hot} \cup \Omega_{f,cold}$.
- ▶ Thermal diffusion in Ω_s and Ω_f with conductivities $k_s \gg k_f$.

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- ▶ Non-penetration constraint:

$$d(\Omega_{f,hot}, \Omega_{f,cold}) \geq d_{min}.$$

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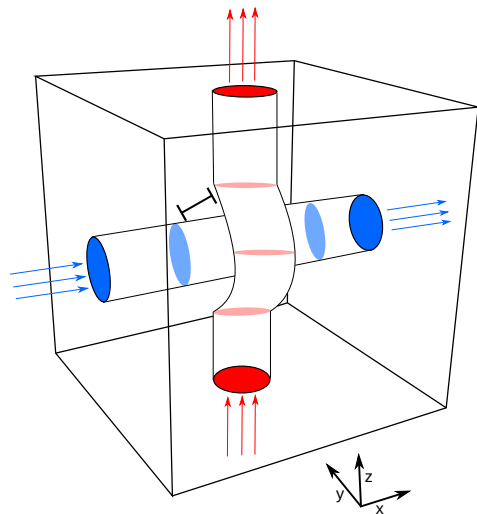


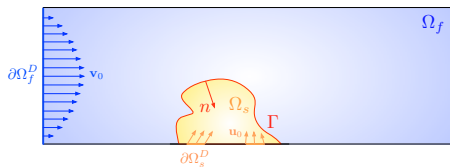
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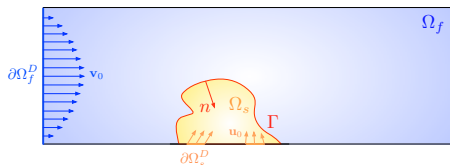
$$d(\Omega_{f,hot}, \Omega_{f,cold}) \geq d_{min}.$$

- ▶ In 3D!

The coupled physics model



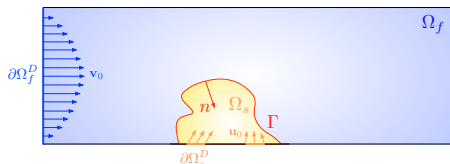
The coupled physics model



- Incompressible Navier-Stokes system for the velocity and pressure (\mathbf{v}, p) in Ω_f

$$\begin{aligned} -\operatorname{div}(\sigma_f(\mathbf{v}, p)) + \rho \nabla \mathbf{v} \mathbf{v} &= \mathbf{f}_f \text{ in } \Omega_f \\ \operatorname{div}(\mathbf{v}) &= 0 \text{ in } \Omega_f \end{aligned}$$

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- Convection-diffusion for the temperature T in Ω_f and Ω_s :

$$\begin{aligned} -\operatorname{div}(k_f \nabla T_f) + \rho \mathbf{v} \cdot \nabla T_f &= Q_f \quad \text{in } \Omega_f \\ -\operatorname{div}(k_s \nabla T_s) &= Q_s \quad \text{in } \Omega_s \end{aligned}$$

Physical solvers

- ▶ The body-fitted approach uses discretization meshes with the fluid interface explicitly discretized.

Physical solvers

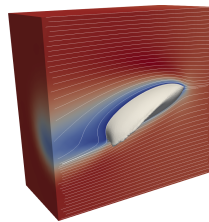
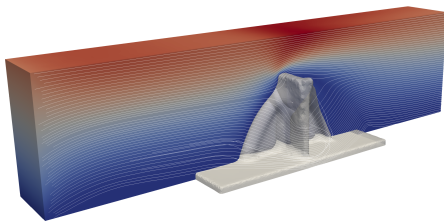
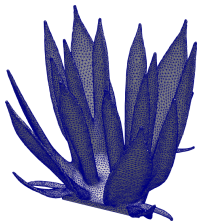
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For 3D applications, absolute need of parallel computing.

- ▶ We rely on finite element discretization of the weak formulation and a Newton scheme for Navier-Stokes problem
- ▶ We Use Domain Decomposition and adapted preconditioners for solving finite element problems : all FEM related operations are achieved in parallel with FreeFEM.

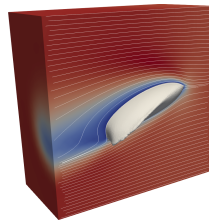
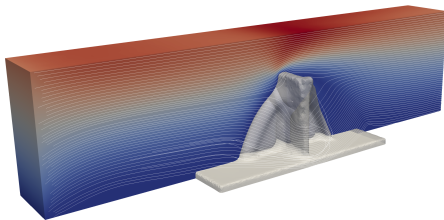
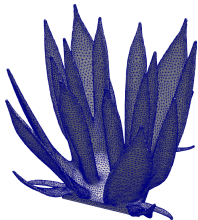


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- ▶ We Use Domain Decomposition and adapted preconditioners for solving finite element problems : all FEM related operations are achieved in parallel with FreeFEM.
- ▶ Our examples feature fluid FEM problems on meshes up to 5 millions of Tetrahedra with 30 CPUs.

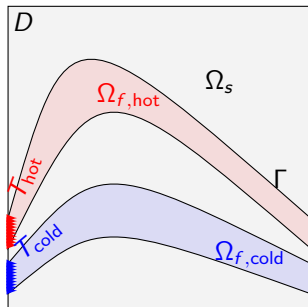


The optimal design problem

- Heat exchanged:

$$W(\Omega_f, \mathbf{v}(\Omega_f), T(\Omega_f))$$

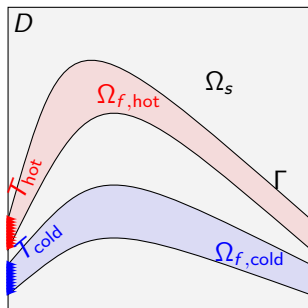
$$:= \int_{\partial\Omega_{f,\text{hot}}} \rho c_p T \mathbf{v} \cdot \mathbf{n} dy - \int_{\partial\Omega_{f,\text{cold}}} \rho c_p T \mathbf{v} \cdot \mathbf{n} dy,$$



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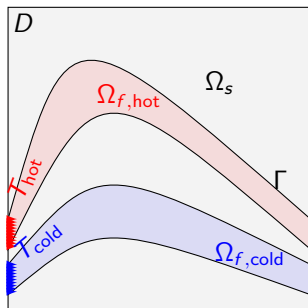
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- Pressure drop:

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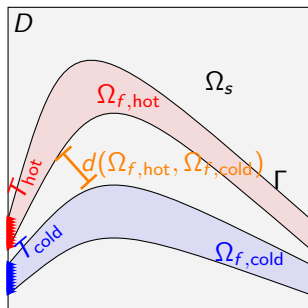
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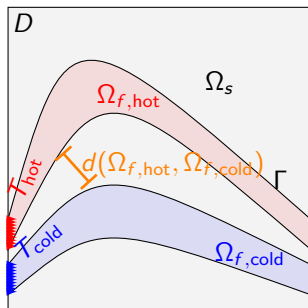
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$$d(\Omega_{f,cold}, \Omega_{f,hot}) := \inf_{\substack{x \in \Omega_{f,cold} \\ y \in \Omega_{f,hot}}} |x - y| \geq d_{min}.$$

The optimal design problem

The shape optimization problem:

$$\begin{aligned} & \max_{\Gamma = \overline{\Omega_f} \cap \overline{\Omega_s}} && W(\Omega_f, \mathbf{v}(\Omega_f), T(\Omega_f)) \\ & \text{s.t.} && \begin{cases} \text{DP}(\Omega_{f,\text{cold}}, p(\Omega_f)) \leq \text{DP}_0 \\ \text{DP}(\Omega_{f,\text{hot}}, p(\Omega_f)) \leq \text{DP}_0 \\ d(\Omega_{f,\text{cold}}, \Omega_{f,\text{hot}}) \geq d_{\min}. \end{cases} \end{aligned}$$

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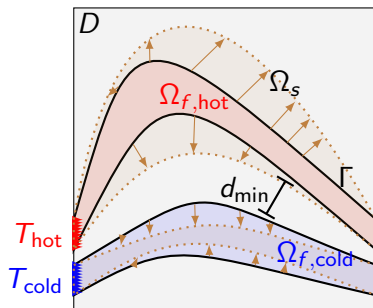
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Optionally, mass constraints on fluids (or on solid):

$$\text{Vol}(\Omega_{f,\text{cold}}) \leq \text{Vol}_0, \quad \text{Vol}(\Omega_{f,\text{hot}}) \leq \text{Vol}_0.$$

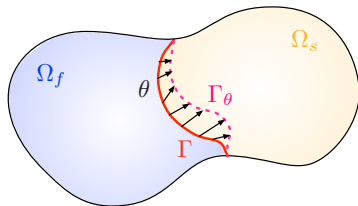
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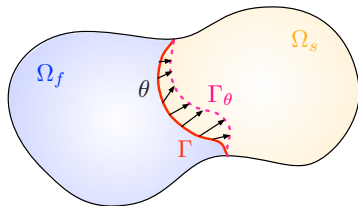
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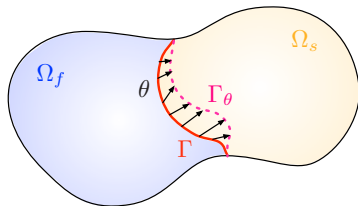
$$\min_{\Gamma} J(\Gamma)$$



$$\Gamma_\theta = (I + \theta)\Gamma, \text{ with } \theta \in W_0^{1,\infty}(D, \mathbb{R}^d), \|\theta\|_{W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d)} < 1.$$

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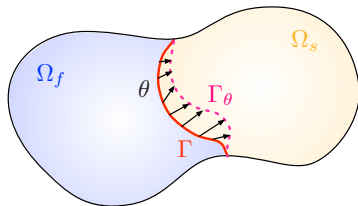


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$$J(\Gamma_{\theta}) = J(\Gamma) + \frac{dJ}{d\theta}(\theta) + o(\theta), \quad \text{with } \frac{|o(\theta)|}{\|\theta\|_{W^{1,\infty}(D, \mathbb{R}^d)}} \xrightarrow{\theta \rightarrow 0} 0.$$

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Under suitable regularity assumptions, Hadamard structure theorem holds:

$$\frac{dJ}{d\theta}(\Omega)(\theta) = \int_{\partial\Omega} v_J(\Omega) \theta \cdot n dy$$

for some $v_J(\Omega) \in L^1(\partial\Omega)$.

Shape derivatives

Proposition

Let $F(\Omega_f, T(\Omega_f), \mathbf{v}(\Omega_f), p(\Omega_f))$ an arbitrary cost function. If F has continuous partial derivatives, then $\Omega_f \mapsto F(\Omega_f, \mathbf{u}(\Omega_f), T(\Omega_f), \mathbf{v}(\Omega_f), p(\Omega_f))$ is shape differentiable and the shape derivative reads²:

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Shape derivatives

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This is the volume expression of the shape derivative. There exists also a surface expression of the form $DF(\boldsymbol{\theta}) = \int_{\Gamma} v_F \boldsymbol{\theta} \cdot \mathbf{n} dy$.

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Two adjoint terms corresponding either of the two physics

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Depends on adjoint states S and (\mathbf{w}, q) solved in inverse cascade.

Shape derivatives

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Adjoint system

Adjoint states in variational formulation fed with **partial derivatives** :

$$\int_{\Omega_s} k_s \nabla S \cdot \nabla S' dx + \int_{\Omega_f} (k_f \nabla S \cdot \nabla S' + \rho c_p S \mathbf{v} \cdot \nabla S') dx = \frac{\partial F}{\partial T}(S).$$

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$$\forall (\mathbf{w}', q') \in V_{\mathbf{v}, p}$$

$$\int_{\Omega_f} \left(\sigma_f(\mathbf{w}, q) : \nabla \mathbf{w}' + \rho \mathbf{w} \cdot \nabla \mathbf{w}' \cdot \mathbf{v} + \rho \mathbf{w} \cdot \nabla \mathbf{v} \cdot \mathbf{w}' - q' \operatorname{div}(\mathbf{w}) \right) dx = \int_{\Omega_f} -\rho c_p S \nabla T \cdot \mathbf{w}' dx + \frac{\partial F}{\partial (\mathbf{v}, p)}(\mathbf{w}', q'),$$

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This can be implemented once for all and allows for easy changes of objective functions.

Adjoint system

This allows to compute the shape derivatives of the heat transfer and of the pressure drop.

$$\begin{aligned} & \max_{\Gamma = \overline{\Omega_f} \cap \overline{\Omega_s}} && W(\Omega_f, \mathbf{v}(\Omega_f), T(\Omega_f)) \\ & \text{s.t.} && \begin{cases} DP(\Omega_{f,\text{cold}}, p(\Omega_f)) \leq DP_0 \\ DP(\Omega_{f,\text{hot}}, p(\Omega_f)) \leq DP_0 \\ d(\Omega_{f,\text{cold}}, \Omega_{f,\text{hot}}) \geq d_{\min}. \end{cases} \end{aligned}$$

Adjoint system

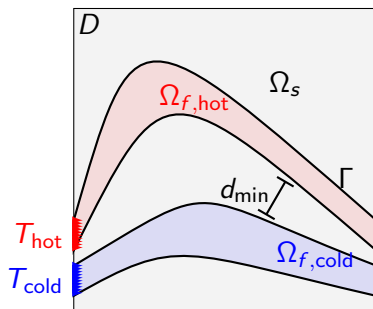
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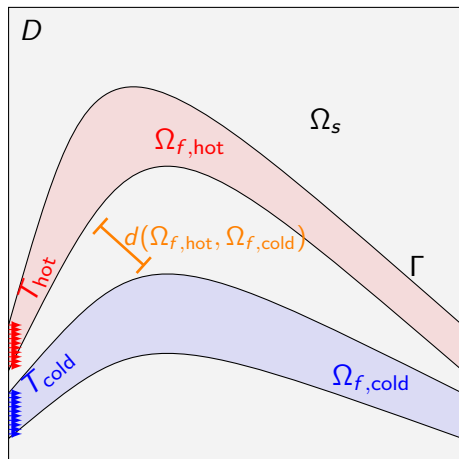
Geometric constraints need a special treatment.

Outline

1. Formulation of the optimal heat exchanger design problem
 - ▶ Physical modelling
 - ▶ Non-mixing constraint



Outline

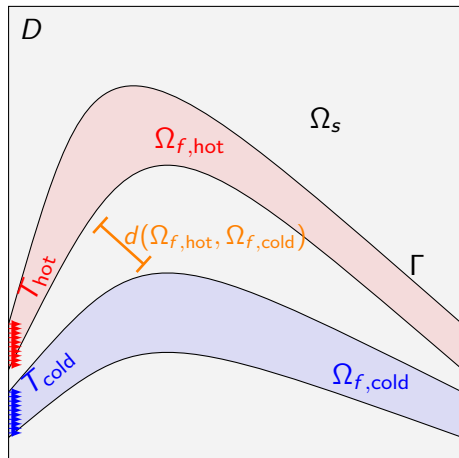


Non-penetration constraint:

$$d(\Omega_{f,hot}, \Omega_{f,cold}) \geq d_{min}.$$

Figure: Settings of the heat exchanger topology optimization problem .

Outline



Non-penetration constraint:

$$d(\Omega_{f,hot}, \Omega_{f,cold}) \geq d_{min}.$$

We enforce it by imposing

$$\forall x \in \Omega_{f,cold}, d_{\Omega_{f,hot}}(x) \geq d_{min},$$

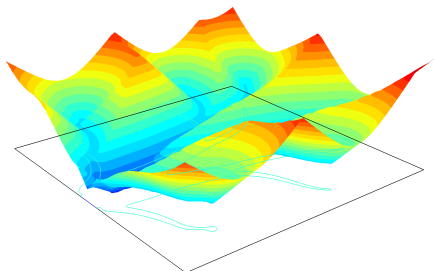
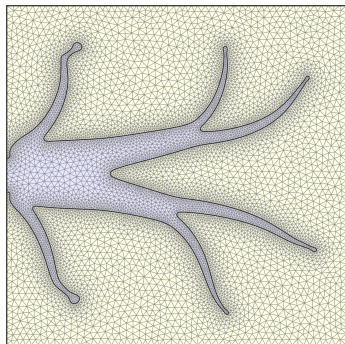
where $d_{\Omega_{f,hot}}$ is the signed distance function to the domain $\Omega_{f,hot}$.

Figure: Settings of the heat exchanger topology optimization problem .

The signed distance function

The signed distance function d_Ω to the domain $\Omega \subset D$ is defined by:

$$\forall x \in D, d_\Omega(x) = \begin{cases} -\min_{y \in \partial\Omega} \|y - x\| & \text{if } x \in \Omega, \\ \min_{y \in \partial\Omega} \|y - x\| & \text{if } x \in D \setminus \Omega. \end{cases}$$



The signed distance function

This allows to compute the shape derivatives of the heat transfer and of the pressure drop.

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Geometric constraints need a special treatment.

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$$d(\Omega_{f,\text{cold}}, \Omega_{f,\text{hot}}) \geq d_{\min}$$

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$$d(\Omega_{f,\text{cold}}, \Omega_{f,\text{hot}}) \geq d_{\min} \Leftrightarrow \forall y \in \partial\Omega_{f,\text{cold}}, d_{\Omega_{f,\text{hot}}}(y) \geq d_{\min}$$

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This is equivalent to

$$\left\| \frac{1}{d_{\Omega_{f,\text{hot}}}} \right\|_{L^\infty(\partial\Omega_{f,\text{cold}})}^{-1} \geq d_{\min}$$

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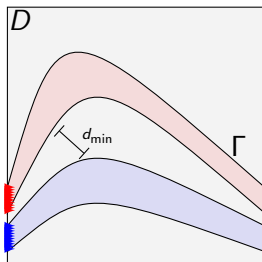
$$\left\| \frac{1}{d_{\Omega_{f,\text{hot}}}} \right\|_{L^\infty(\partial\Omega_{f,\text{cold}})}^{-1} \geq d_{\min}$$

We approximate the infinity norm with an averaged p -norm:

$$P_{\text{cold} \rightarrow \text{hot}}(\Omega_f) \geq d_{\min},$$

$$\text{with } P_{\text{cold} \rightarrow \text{hot}}(\Omega_f) := \left\| \frac{1}{d_{\Omega_{f,\text{hot}}}} \right\|_{L^p(\partial\Omega_{f,\text{cold}})}^{-1} = \left(\int_{\partial\Omega_{f,\text{cold}}} \frac{1}{|d_{\Omega_{f,\text{hot}}}|^p} ds \right)^{-\frac{1}{p}}.$$

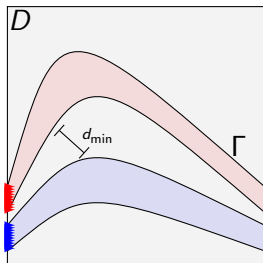
2D heat exchangers



Heat exchanger problem with limited pressure loss and non-mixing constraint:

$$\begin{aligned} \min_{\Gamma} \quad & J(\Omega_f) = - \left(\int_{\Omega_{f,cold}} \rho c_p \mathbf{v} \cdot \nabla T dx - \int_{\Omega_{f,hot}} \rho c_p \mathbf{v} \cdot \nabla T dx \right) \\ \text{s.c.} \quad & \begin{cases} DP(\Omega_f) = \int_{\partial\Omega_f^D} p ds - \int_{\partial\Omega_f^N} p ds \leq DP_0 \\ P_{\text{cold} \rightarrow \text{hot}}(\Omega_f) \geq d_{\min}. \end{cases} \end{aligned}$$

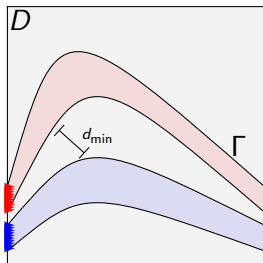
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What is the shape derivative of $P_{cold \rightarrow hot}(\Omega_f)$?

Shape derivatives of geometric constraints

This reduces to the setting of computing the shape derivative of $d_{\Omega_f, \text{hot}}$, $P_{\text{hot} \leftrightarrow \text{cold}}(\Omega_f)$ with:

$$P_{\text{hot} \leftrightarrow \text{cold}}(\Omega_f) := \int_D j(d_{\Omega_f, \text{hot}}) dx.$$

Shape derivatives of geometric constraints

The shape derivative of $P_{hot \leftrightarrow cold}(\Omega_f)$ is given by³:

$$P'_{hot \leftrightarrow cold}(\Omega)(\theta) = \int_{\partial\Omega_{f,hot}} u(y) \theta \cdot \mathbf{n} dy$$

$$\text{with } u(y) = - \int_{z \in \text{ray}(y)} j'(d_{\Omega_{f,hot}}(z)) \prod_{1 \leq i \leq n-1} (1 + \kappa_i(y) d_{\Omega_{f,hot}}(z)) dz, \quad \forall y \in \partial\Omega.$$

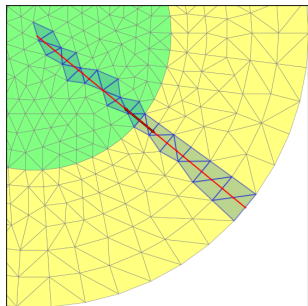
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The computation of $u(y)$ requires a priori integration along the normal rays and the computation of curvatures $\kappa_i(y)$.

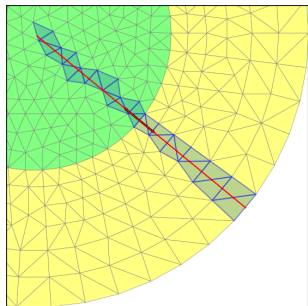
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It turns out that it is possible to compute u without integrating along the rays⁴:

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Proposition

Let $\hat{u} \in V_\omega$ be the solution to the variational problem

$$\forall v \in V_\omega, \int_{\partial\Omega_{f,hot}} \hat{u} v ds + \int_D \omega (\nabla d_{\Omega_{f,hot}} \cdot \nabla \hat{u}) (\nabla d_{\Omega_{f,hot}} \cdot \nabla v) dx = - \int_D j'(d_{\Omega_{f,hot}}) v dx.$$

Then $u(y) = \hat{u}(y)$ for any $y \in \partial\Omega_{f,hot}$.

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Then $u(y) = \hat{u}(y)$ for any $y \in \partial\Omega_{f,hot}$.

- ▶ This variational problem can easily be solved with FEM in 2-D and 3D
- ▶ This allows to handle conveniently geometric constraints (e.g. maximum thickness, minum distance, etc. . .) in 2D and 3D level set based topology optimization.

⁴Feppon, Allaire, and Dapogny, *A variational formulation for computing shape derivatives of geometric constraints along rays* (2019)

Shape derivatives of geometric constraints

A comparison with an analytic case:

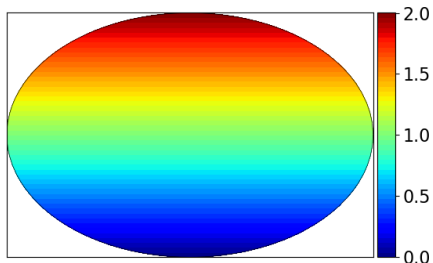
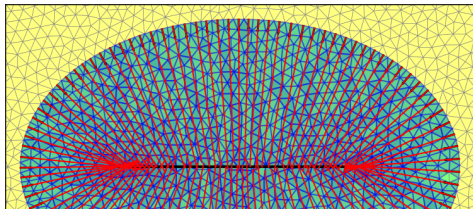
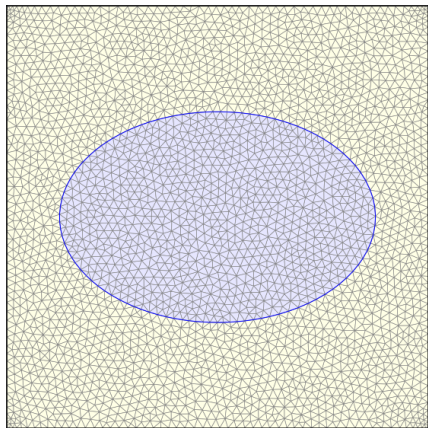
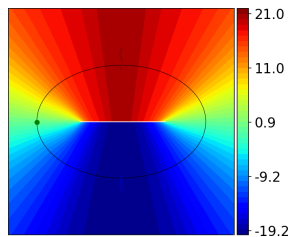


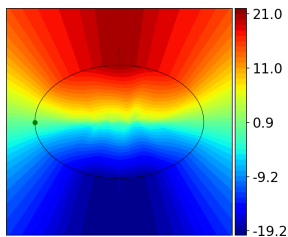
Figure: A prescribed $-j'(d_{\Omega}(x))$.

Shape derivatives of geometric constraints

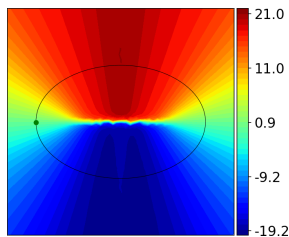
The weight ω needs to vanish near the skeleton (medial axis).



(a) Mesh \mathcal{T}' (manually removed skeleton), $\omega = 1$

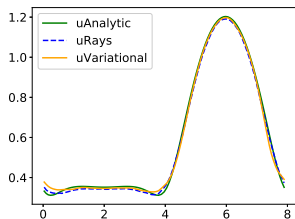


(b) Mesh \mathcal{T} , $\omega = 1$.

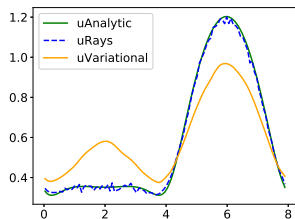


(c) Mesh \mathcal{T} ,
 $\omega = 2/(1 + |\Delta d_{\Omega}|^{3.5})$

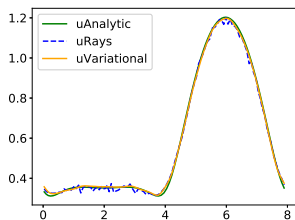
Shape derivatives of geometric constraints



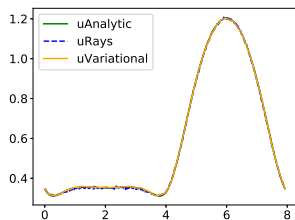
(a) Mesh \mathcal{T}' , $\omega = 1$



(b) Mesh \mathcal{T} , $\omega = 1$



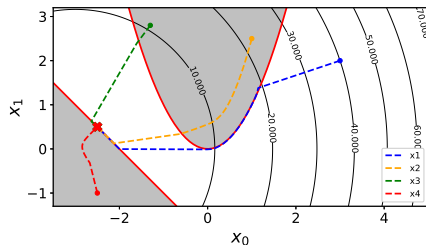
(c) Mesh \mathcal{T} , $\omega = 2/(1 + |\Delta d_{\Omega}|^{3.5})$



(d) Finer mesh \mathcal{T} , $\omega = 2/(1 + |\Delta d_{\Omega}|^{3.5})$

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 - ▶ Null space optimization algorithm



- Our goal: solve constrained shape optimization problems

$$\begin{aligned} \min_{\Gamma} \quad & J(\Gamma, \mathbf{v}(\Gamma), p(\Gamma), T(\Gamma)) \\ \text{s.t.} \quad & g_i(\Gamma, \mathbf{v}(\Gamma), p(\Gamma), T(\Gamma)) = 0, \quad 1 \leq i \leq p. \\ & h_i(\Gamma, \mathbf{v}(\Gamma), p(\Gamma), T(\Gamma)) \leq 0, \quad 1 \leq i \leq q \end{aligned}$$

with *arbitrary* functionals J, g_i, h_i ;

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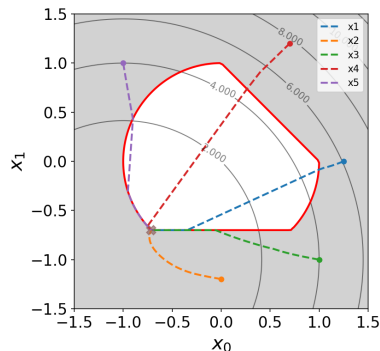
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- ▶ g_i and h_i represent industrial specification constraints (mass, pressure drop...)

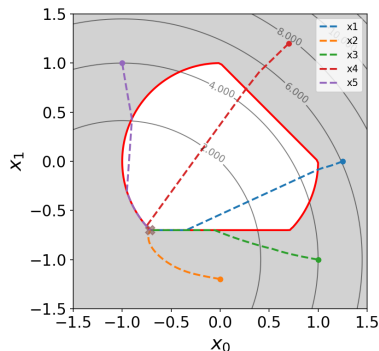
Null space optimization algorithm

- Nonlinear constrained optimization on manifolds with a moderate number of constraints



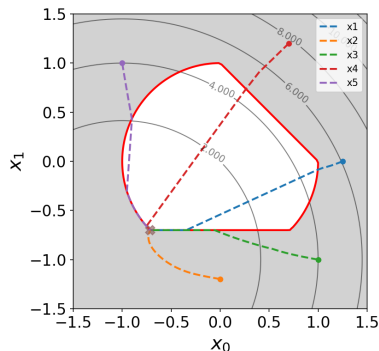
Null space optimization algorithm

- ▶ Nonlinear constrained optimization on manifolds with a moderate number of constraints
- ▶ Generalization of the unconstrained gradient flow: no hard tuning of parameters



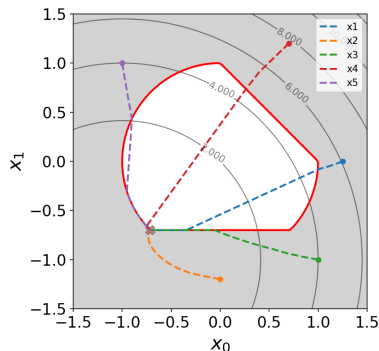
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Null space optimization algorithm

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- ▶ Gradual corrections of unfeasible initializations
- ▶ Adapted to the infinite dimensional setting of the method of Hadamard



Null space optimization algorithm

For the exposure of our method, let us consider

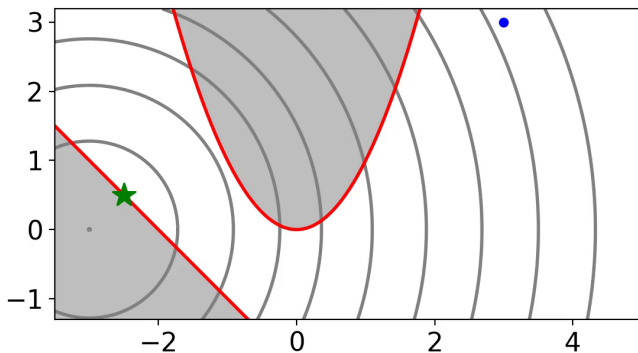
$$\begin{aligned} \min_{x \in \mathbb{R}^n} \quad & J(x) \\ \text{s.t.} \quad & \begin{cases} \mathbf{g}(x) = 0 \\ \mathbf{h}(x) \leq 0, \end{cases} \end{aligned}$$

with $J : \mathbb{R}^n \rightarrow \mathbb{R}$, $\mathbf{g} : \mathbb{R}^n \rightarrow \mathbb{R}^p$ and $\mathbf{h} : \mathbb{R}^n \rightarrow \mathbb{R}^q$ Fréchet differentiable.

Null space optimization algorithm

$$\min_{(x_1, x_2) \in \mathbb{R}^2} J(x_1, x_2) = x_1^2 + (x_2 + 3)^2$$

$$\text{s.t.} \begin{cases} h_1(x_1, x_2) = -x_1^2 + x_2 & \leq 0 \\ h_2(x_1, x_2) = -x_1 - x_2 - 2 & \leq 0 \end{cases}$$



Null space optimization algorithm

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$$\dot{x} = -(I - D\mathbf{g}^T(D\mathbf{g}D\mathbf{g}^T)^{-1}D\mathbf{g})\nabla J(x)$$

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$$\dot{x} = -\alpha_J(I - D\mathbf{g}^T(D\mathbf{g}D\mathbf{g}^T)^{-1}D\mathbf{g})\nabla J(x) - \alpha_C D\mathbf{g}^T(D\mathbf{g}D\mathbf{g}^T)^{-1}\mathbf{g}(x)$$

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$\mathbf{g}(x(t)) = \mathbf{g}(x(0))e^{-\alpha_C t}$ and $J(x(t))$ decreases if $\mathbf{g}(x(t)) = 0$.

Null space optimization algorithm

For *both* equality constraints $\mathbf{g}(x) = 0$ and inequality constraints $\mathbf{h}(x) \leq 0$, we consider:

$$\dot{x} = -\alpha_J \xi_J(x(t)) - \alpha_C \xi_C(x(t))$$

with

$$\xi_J(x) := (I - \mathbf{D}\mathbf{C}_{\hat{I}(x)}^T (\mathbf{D}\mathbf{C}_{\hat{I}(x)} \mathbf{D}\mathbf{C}_{\hat{I}(x)}^T)^{-1} \mathbf{D}\mathbf{C}_{\hat{I}(x)}) (\nabla J(x))$$

$$\xi_C(x) = \mathbf{D}\mathbf{C}_{\tilde{I}(x)}^T (\mathbf{D}\mathbf{C}_{\tilde{I}(x)} \mathbf{D}\mathbf{C}_{\tilde{I}(x)}^T)^{-1} \mathbf{C}_{\tilde{I}(x)}(x),$$

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$$\xi_C(x) = \text{DC}_{\tilde{I}(x)}^T (\text{DC}_{\tilde{I}(x)} \text{DC}_{\tilde{I}(x)}^T)^{-1} \mathbf{C}_{\tilde{I}(x)}(x),$$

$\tilde{I}(x)$ the set of violated constraints:

$$\tilde{I}(x) = \{i \in \{1, \dots, q\} \mid h_i(x) \geq 0\}.$$

$$\mathbf{C}_{\tilde{I}(x)} = \left[\mathbf{g}(x) \mid (h_i(x))_{i \in \tilde{I}(x)} \right]^T$$

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$\hat{I}(x) \subset \tilde{I}(x)$ is an “optimal” subset of the active or violated constraints which can be computed by mean of a dual subproblem.

$$\hat{I}(x) := \{i \in \tilde{I}(x) \mid \mu_i^*(x) > 0\}.$$

$$\mathbf{C}_{\hat{I}(x)} = \left[\mathbf{g}(x) \quad \mid \quad (h_i(x))_{i \in \hat{I}(x)} \right]^T$$

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The transpose \cdot^T operator encompasses the regularization and extension step of the shape derivative.

Null space optimization algorithm

We can prove:

1. Constraints are asymptotically satisfied:

$$\mathbf{g}(x(t)) = e^{-\alpha c t} \mathbf{g}(x(0)) \text{ and } \mathbf{h}_{\tilde{l}(x(t))} \leq e^{-\alpha c t} \mathbf{h}(x(0))$$

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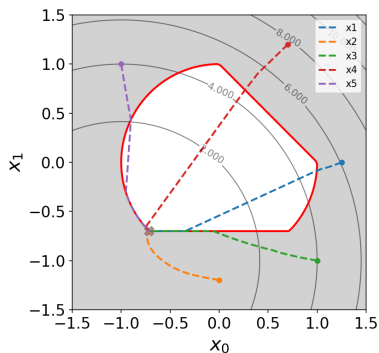
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The algorithm can be adapted to the infinite-dimensional shape optimization framework.

Null space optimization algorithm

Try it yourself! Open source implementation⁵:



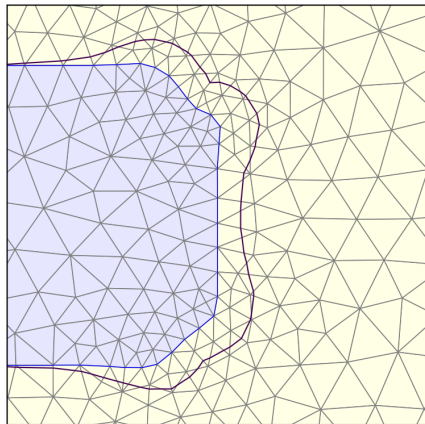
<https://gitlab.com/florian.feppon/null-space-optimizer>

```
pip install nullspace_optimizer
```

⁵Feppon, Allaire, and Dapogny, *Null space gradient flows for constrained optimization with applications to shape optimization* (2019)

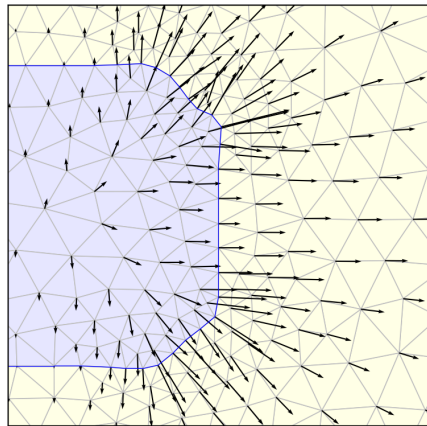
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Body-fitted meshes

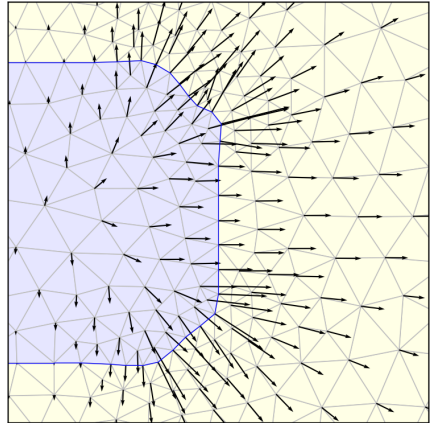
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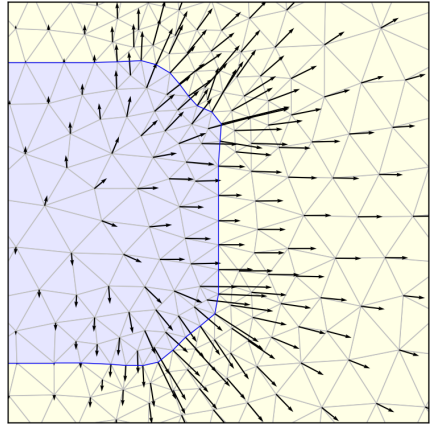
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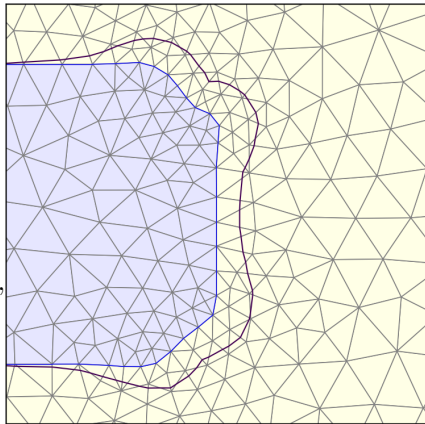
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$$\begin{cases} \frac{\partial \phi}{\partial t}(t, \mathbf{x}) + \boldsymbol{\theta}(t, \mathbf{x}) \cdot \nabla \phi(t, \mathbf{x}) = 0, \\ \phi(0, \mathbf{x}) = \phi^n(\mathbf{x}), \end{cases}$$

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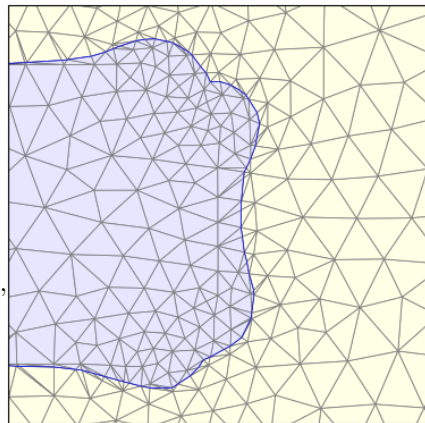
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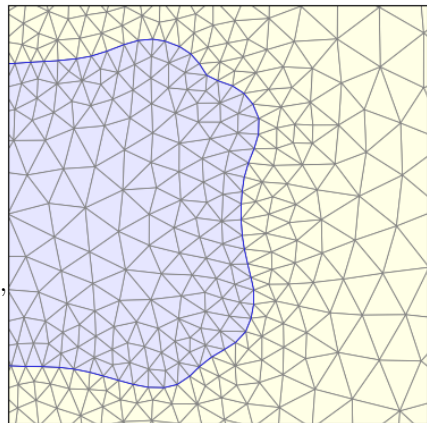
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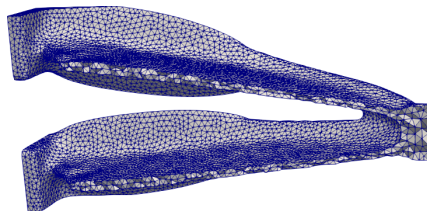
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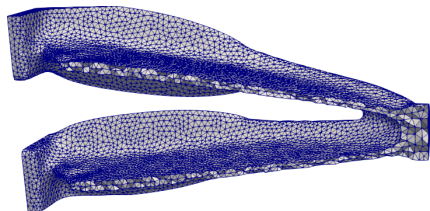
⁶Allaire, Dapogny, and Frey, *Shape optimization with a level set based mesh evolution method* (2014)

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Remark: Mesh adaptation and Isosurface discretization in Mmg is still sequential. A future release of (Par)Mmg will allow to remesh in parallel.

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3D thermal diffusion

Maximization of heat conduction:

$$\begin{aligned} \min_{\Omega_f \subset D} \quad & \int_D T dx \\ \text{s.c.} \quad & \int_{\Omega_f} dx \leq V_0 \end{aligned}$$

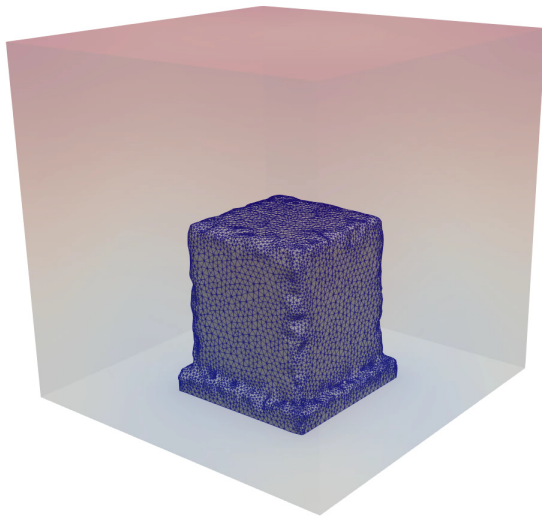
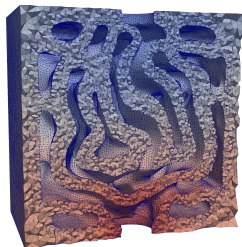


Figure: Thermal diffusion

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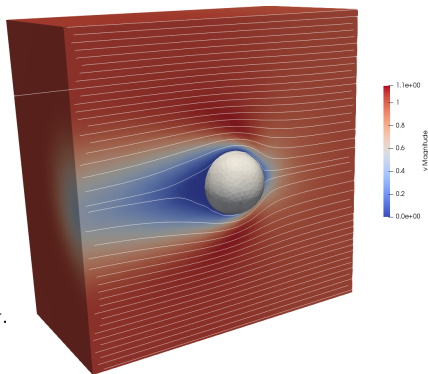


Lift-Drag optimization

$$\begin{aligned} \min \quad & -\text{Lift}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) \\ \text{s.c.} \quad & \left\{ \begin{array}{l} \text{Drag}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) \leq \text{DRAG}_0 \\ \text{Vol}(\Omega_f) = V_0 \\ \mathbf{X}(\Omega_s) := \frac{1}{|\Omega_s|} \int_{\Omega_s} \mathbf{x} dx = \mathbf{x}_0, \end{array} \right. \end{aligned}$$

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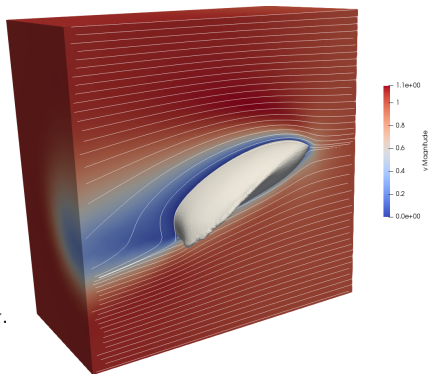


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$$\text{Lift}(\Gamma, \mathbf{v}(\Gamma), p(\Gamma)) := - \int_{\Gamma} \mathbf{e}_y \cdot \sigma_f(\mathbf{v}, p) \mathbf{n} ds,$$

$$\text{Drag}(\Gamma, \mathbf{v}(\Gamma), p(\Gamma)) := \int_{\Omega_f} \sigma_f(\mathbf{v}, p) : \nabla \mathbf{v} dx.$$

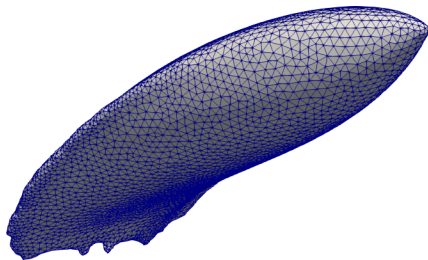


Lift-Drag optimization

$$\begin{aligned} \min \quad & -\text{Lift}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) \\ \text{s.c.} \quad & \left\{ \begin{array}{l} \text{Drag}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) \leq \text{DRAG}_0 \\ \text{Vol}(\Omega_f) = V_0 \\ \mathbf{X}(\Omega_s) := \frac{1}{|\Omega_s|} \int_{\Omega_s} \mathbf{x} dx = \mathbf{x}_0, \end{array} \right. \end{aligned}$$

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$$\text{Drag}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) := \int_{\Omega_f} \sigma_f(\mathbf{v}, \rho) : \nabla \mathbf{v} dx.$$

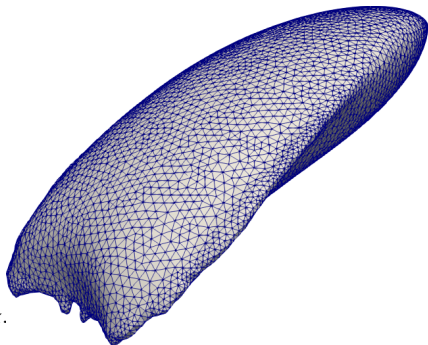


Lift-Drag optimization

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$$\text{Lift}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) := - \int_{\Gamma} \mathbf{e}_y \cdot \sigma_f(\mathbf{v}, \rho) \mathbf{n} ds,$$

$$\text{Drag}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) := \int_{\Omega_f} \sigma_f(\mathbf{v}, \rho) : \nabla \mathbf{v} dx.$$

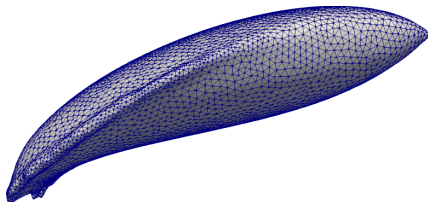


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$$\text{Lift}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) := - \int_{\Gamma} \mathbf{e}_y \cdot \sigma_f(\mathbf{v}, \rho) \mathbf{n} ds,$$

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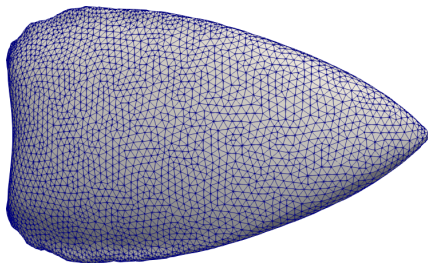


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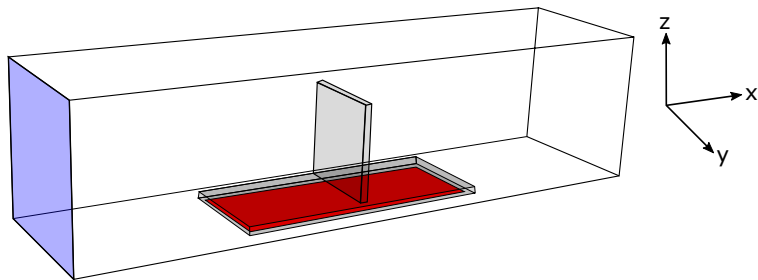
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$$\text{Drag}(\Gamma, \mathbf{v}(\Gamma), \rho(\Gamma)) := \int_{\Omega_f} \sigma_f(\mathbf{v}, \rho) : \nabla \mathbf{v} dx.$$



Fluid-structure interaction

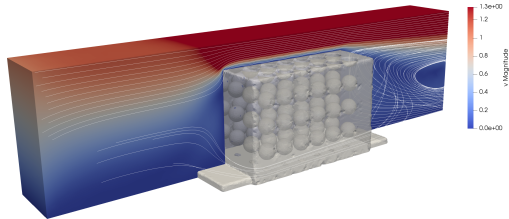
Minimization of the rigidity of a supporting structure subject to the pressure of an incoming flow.



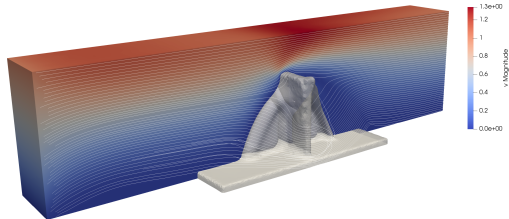
$$\begin{aligned} \min \quad & \int_{\Omega_s} A e(\mathbf{u}) : e(\mathbf{u}) dx \\ \text{s.c.} \quad & \text{Vol}(\Omega_s) = \text{Vol}_{\text{target}}. \end{aligned}$$

Fluid-structure interaction

Minimization of the rigidity of a supporting structure subject to the pressure of an incoming flow.



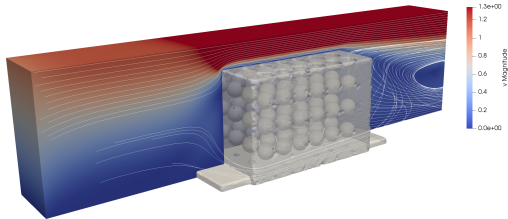
(a) Initial shape



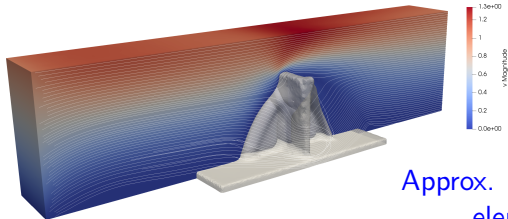
(b) Optimized shape

Fluid-structure interaction

Minimization of the rigidity of a supporting structure subject to the pressure of an incoming flow.



(a) Initial shape



(b) Optimized shape

Approx. 2 millions of elements.

Fluid-structure interaction

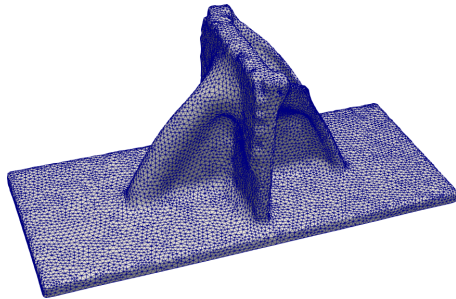


Figure: Optimized shape.

Fluid-structure interaction

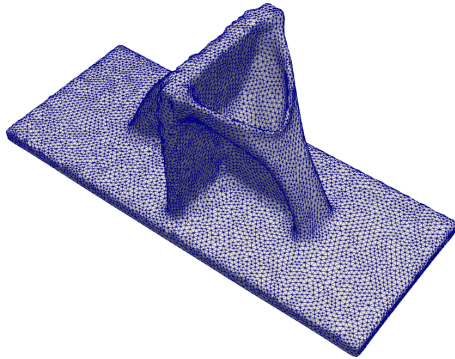


Figure: Optimized shape.

Fluid-structure interaction

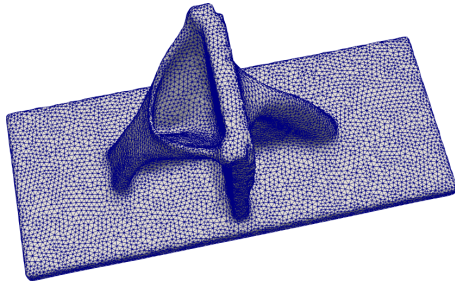


Figure: Optimized shape.

Fluid-structure interaction

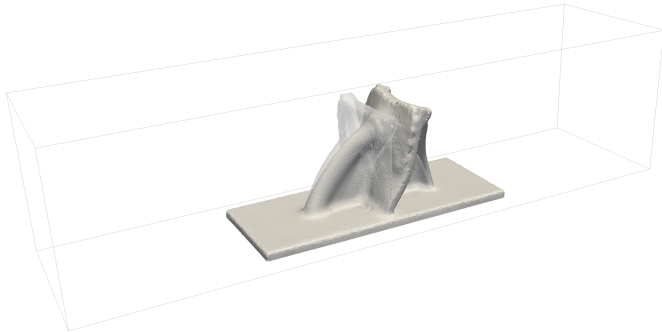
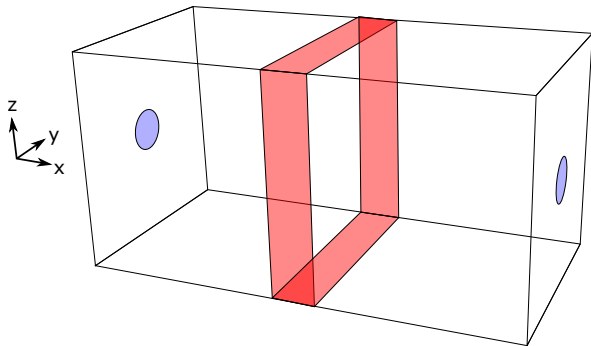


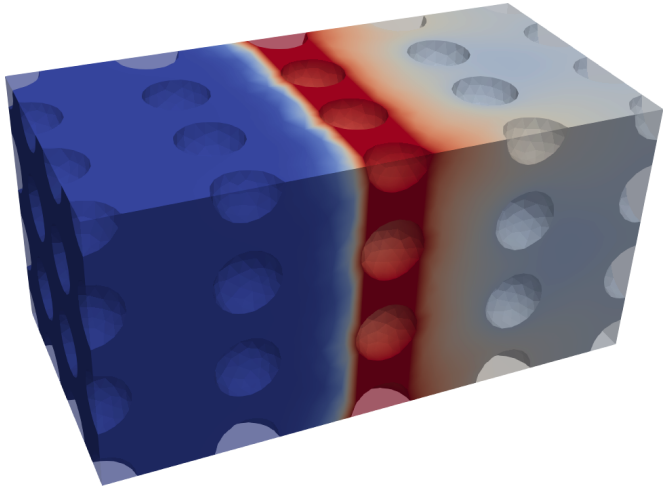
Figure: Elastic deformation.

3D convective heat transfer

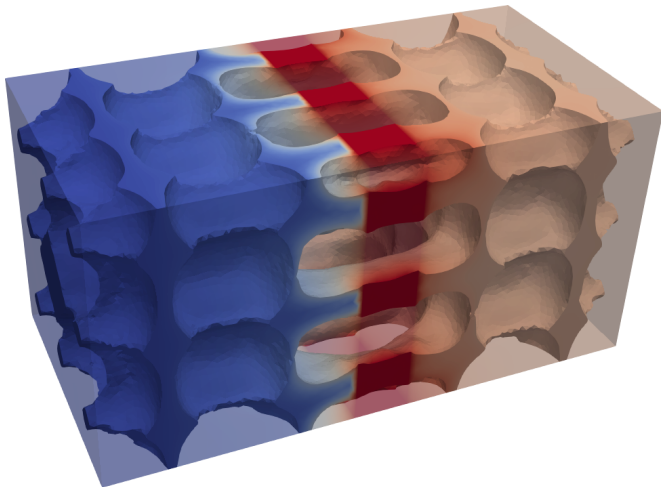


$$\begin{aligned} \min_{\Gamma} \quad & J(\Gamma, \mathbf{v}(\Gamma), T(\Gamma)) := - \int_{\Omega_f} \rho c_p \mathbf{v} \cdot \nabla T \, dx \\ \text{s.t.} \quad & \left\{ \begin{aligned} \text{DP}(\rho(\Gamma)) &:= \int_{\partial\Omega_f^{\text{in}}} \rho \, ds - \int_{\partial\Omega_f^{\text{N}}} \rho \, ds \leq \text{DP}_T \\ \text{Vol}(\Omega_f) &= V_T. \end{aligned} \right. \end{aligned}$$

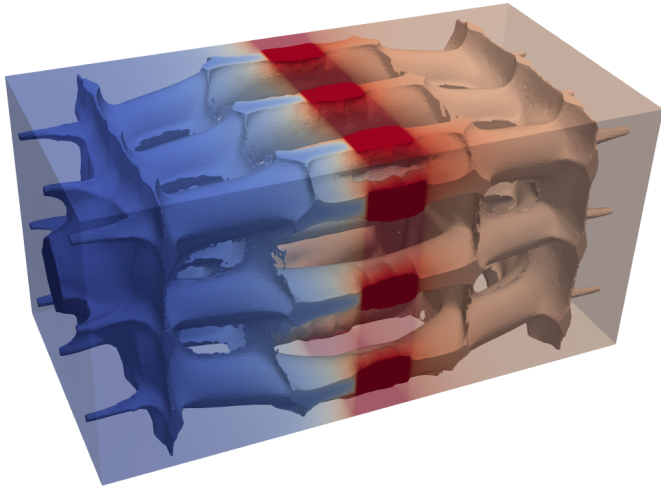
3D convective heat transfer



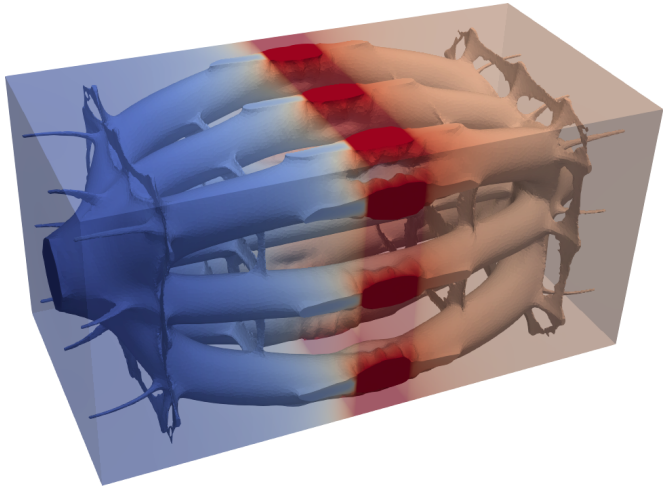
3D convective heat transfer



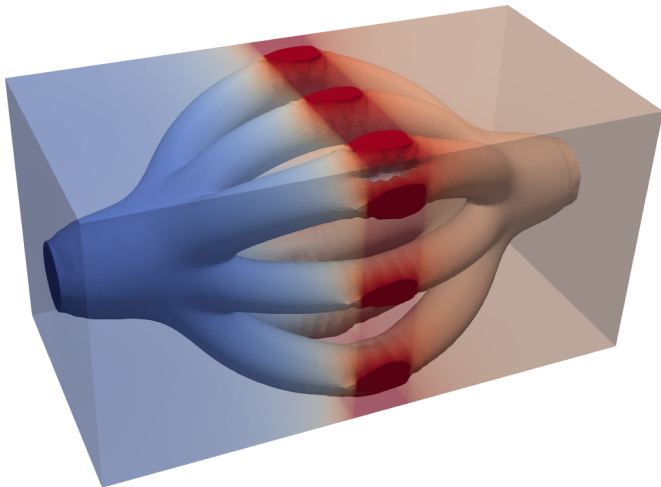
3D convective heat transfer



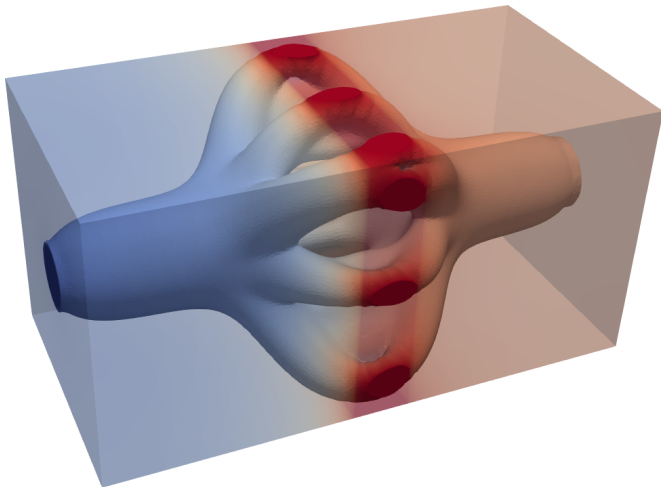
3D convective heat transfer



3D convective heat transfer



3D convective heat transfer



3D convective heat transfer

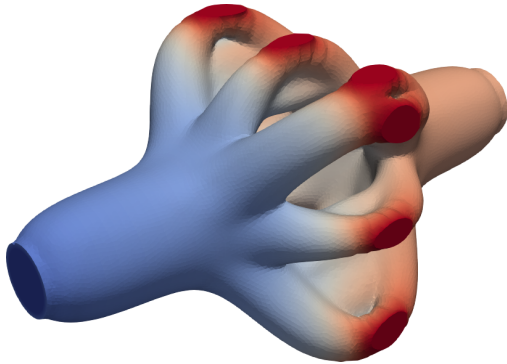


Figure: Optimized design.

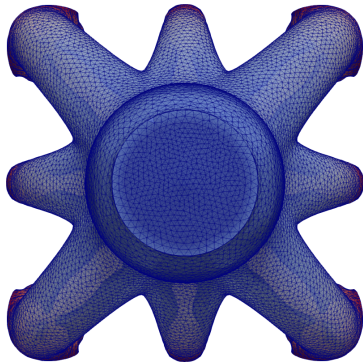


Figure: Optimized design.

3D convective heat transfer

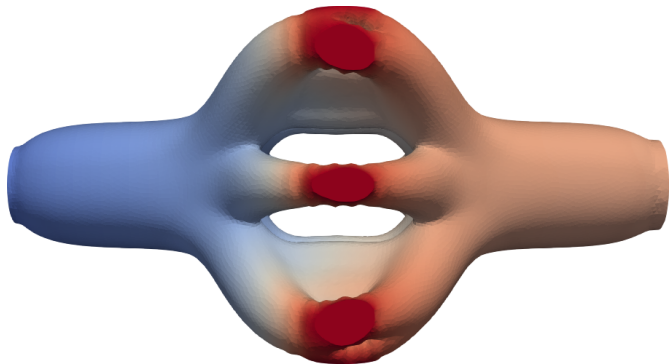


Figure: Optimized design.

3D convective heat transfer

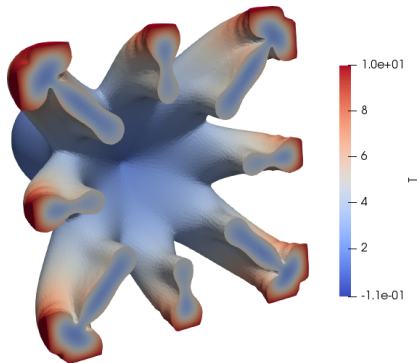


Figure: Optimized design.

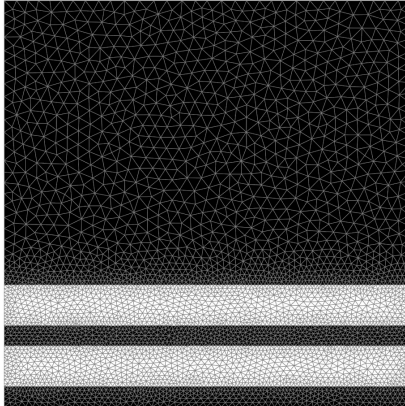
Fluid-to-fluid heat exchangers

This allows to compute the shape derivatives of the heat transfer and of the pressure drop.

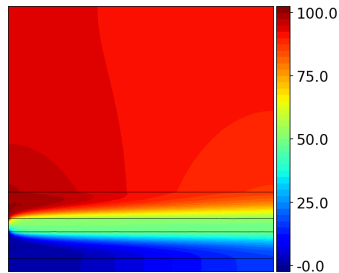
$$\begin{aligned} & \max_{\Gamma = \overline{\Omega_f} \cap \overline{\Omega_s}} && W(\Omega_f, \mathbf{v}(\Omega_f), T(\Omega_f)) \\ & \text{s.t.} && \begin{cases} DP(\Omega_{f,\text{cold}}, p(\Omega_f)) \leq DP_0 \\ DP(\Omega_{f,\text{hot}}, p(\Omega_f)) \leq DP_0 \\ d(\Omega_{f,\text{cold}}, \Omega_{f,\text{hot}}) \geq d_{\min}. \end{cases} \end{aligned}$$

2D counter-current Heat exchanger

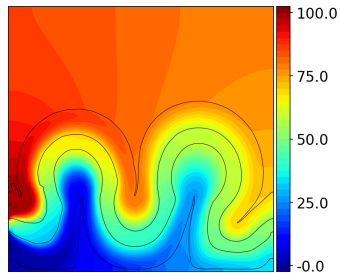
Iteration 0



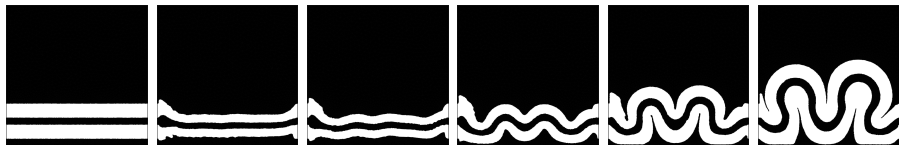
2D Heat Exchangers with non-mixing constraint



(a) Initial temperature



(b) Final temperature.



(c) Intermediate iterations 0, 8, 20, 50, 88 et 200.

2D Heat Exchangers with non-mixing constraint

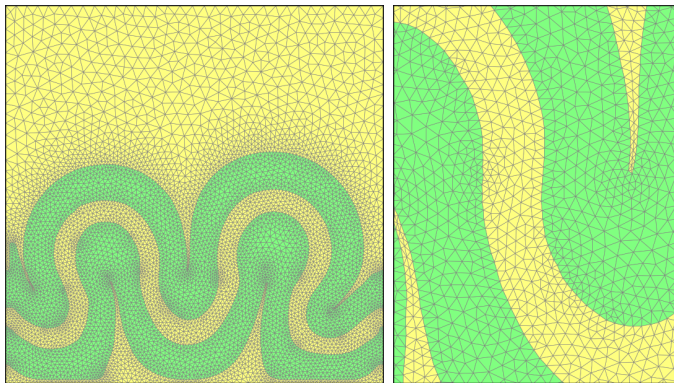
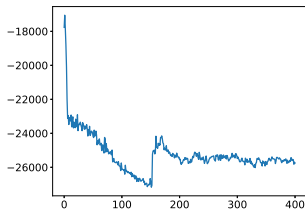
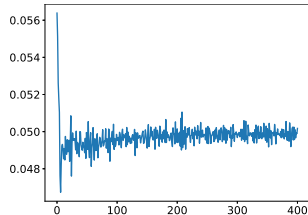


Figure: Zoom on the optimized mesh.

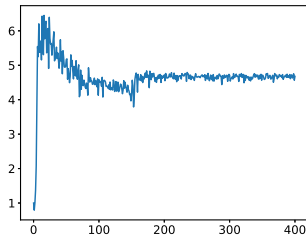
2D Heat Exchangers with non-mixing constraint



(a) Objective function $J(\Omega_f)$.



(b) Distance constraint $P_{\text{cold} \rightarrow \text{hot}}$ for the non-mixing constraint.



(c) Pressure drop constraint $DP(\Omega_f)/DP(\Omega_f^0)$.

3D fluid-to-fluid heat exchanger

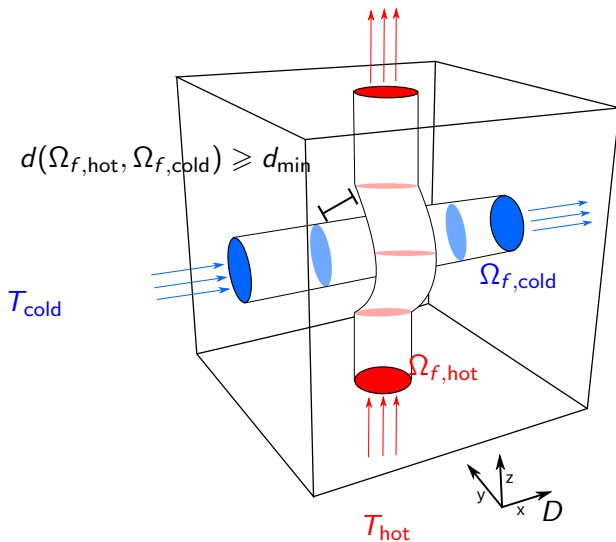


Figure: Schematic of the 3D setting.

3D fluid-to-fluid heat exchanger

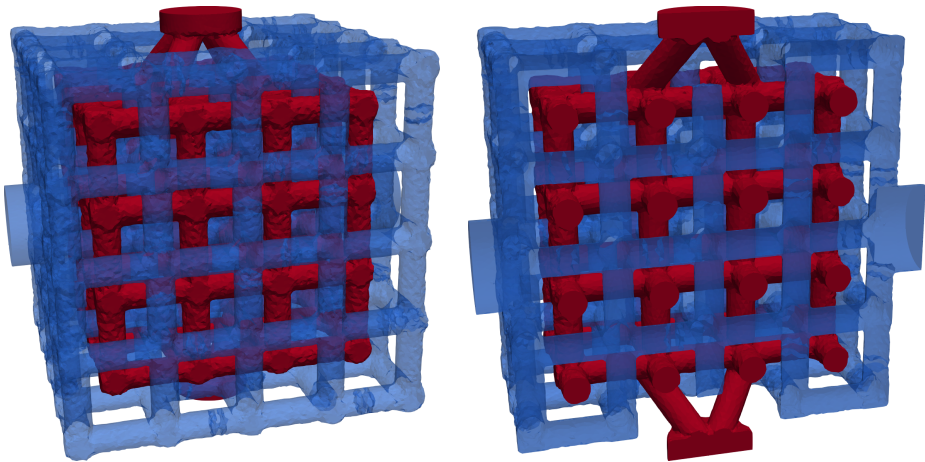
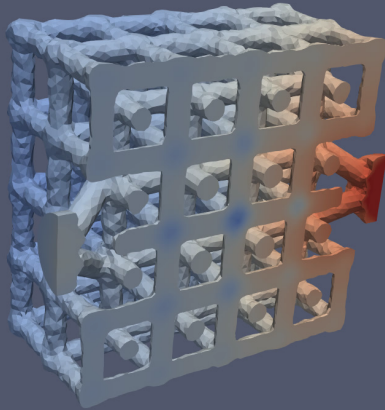
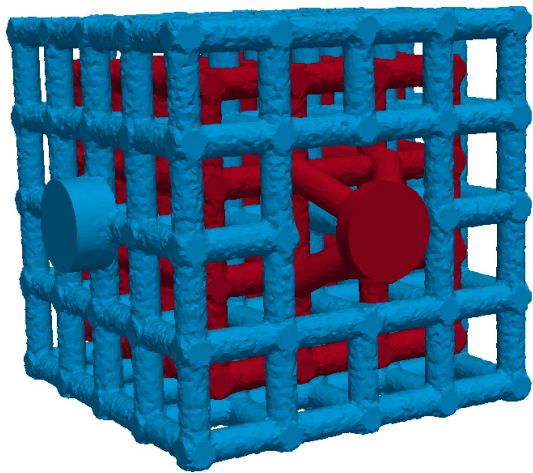


Figure: Initial distribution of fluid considered for the 3D heat exchanger test case.

3D fluid-to-fluid heat exchanger



3D fluid-to-fluid heat exchanger



3D fluid-to-fluid heat exchanger

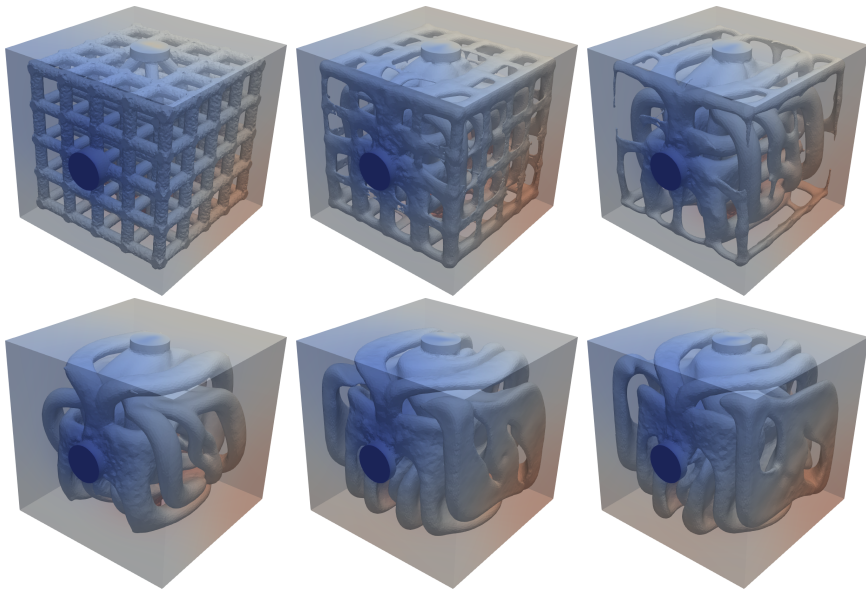
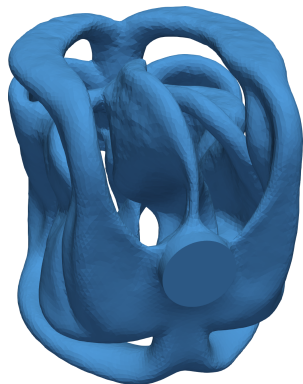
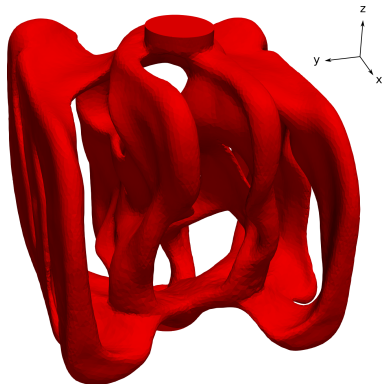


Figure: Intermediate iterations.

3D fluid-to-fluid heat exchanger



(a) Cold phase



(b) Hot phase

Figure: Separate plots of the topologically optimized cold and hot fluid phases in the configuration $d_{\min} = 0.04$.

3D fluid-to-fluid heat exchanger

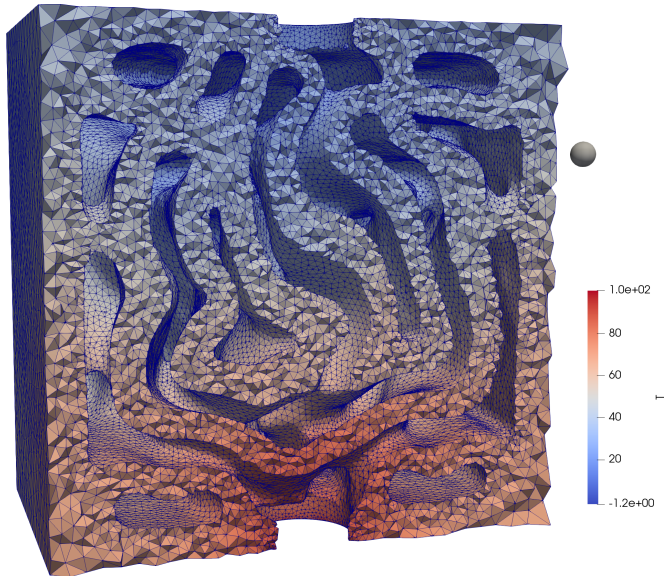
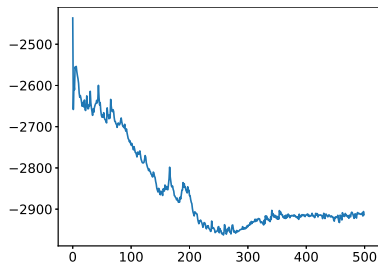
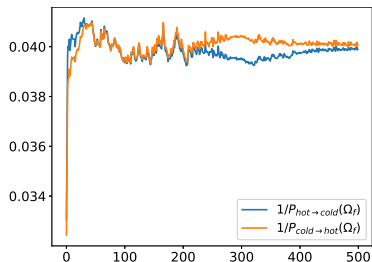


Figure: Cut of the resulting solid domain

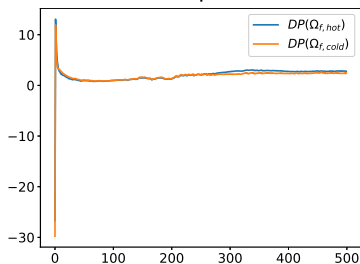
3D fluid-to-fluid heat exchanger



(a) Objective function (opposite of the heat exchanged).



(b) Averaged distance between the two phases.



(c) Pressure drop constraint.

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 1. nonlinearities are differentiable

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 2. adjoint equations (linearized transpose) can be solved numerically

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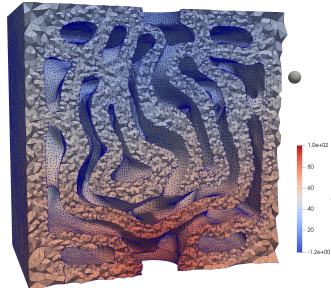
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- ▶ Other topology optimization approaches, such as homogenization based, could lead to alternative methods for generating complex design.

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We expect turbulence can be treated in this manner.

- ▶ 3D remeshing is a bottle neck. Parallel remeshing will substantially reduce computational times.
- ▶ Other topology optimization approaches, such as homogenization based, could lead to alternative methods for generating complex design.
- ▶ In contrast with density based methods, the body fitted approach allows to treat explicitly the non-mixing constraint, and is compatible in principle with non-intrusive solvers.

Many thanks for your attention!

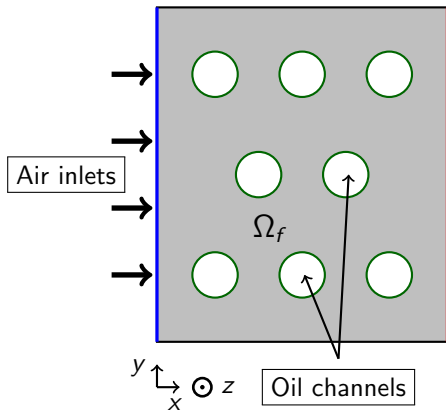


Appendix : an alternative, simple 2D heat exchanger model.

2D Air oil heat exchanger

Safran Aeroboosters case study:

$$T = T_{oil} \text{ on } \Gamma.$$
$$T_{air} < T_{oil}.$$



2D Air oil heat exchanger

- Optimization problem:

$$\begin{aligned} \min_{\Omega_f \subset D} \quad & J(\Omega_f) := - \int_{\Omega_f} \rho c_p \mathbf{v} \cdot \nabla T \, dx \\ \text{s. c.} \quad & \text{DP}(\Omega_f) := \int_{\partial\Omega_{f,in}} p \, ds - \int_{\partial\Omega_{f,out}} p \, ds \leq \text{DP}_0. \end{aligned}$$

2D Air oil heat exchanger

- Optimization problem:

$$\begin{aligned} \min_{\Omega_f \subset D} \quad & J(\Omega_f) := - \int_{\Omega_f} \rho c_p \mathbf{v} \cdot \nabla T \, dx \\ \text{s.c.} \quad & \text{DP}(\Omega_f) := \int_{\partial\Omega_{f,in}} p \, ds - \int_{\partial\Omega_{f,out}} p \, ds \leq \text{DP}_0. \end{aligned}$$

- We consider an alternative formulation to impose a minimum thickness constraint on the oil channels.

$$\begin{aligned} \min_{\Omega_f \subset D} \quad & E(\Omega_f) := - \int_{D \setminus \Omega_f} d_{\Omega_f}^2 \max(-d_{\Omega_f} + d_{\min}/2, 0)^2 \, dx \\ \text{s.c.} \quad & \begin{cases} \text{DP}(\Omega_f) \leq \text{DP}_0 \\ J(\Omega_f) \leq J_0. \end{cases} \end{aligned}$$

2D Air oil heat exchanger

Results :

Test case 7

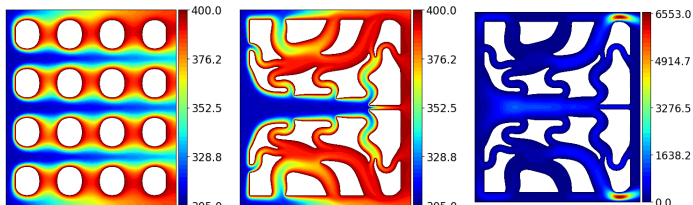
$$\|\mathbf{v}_0\|_\infty = 10$$

$$DP_0 = 1300$$

$$J_{final} = 4086$$

$$DP_{final} = 1308$$

Without
minimum
thickness.



Test case 8

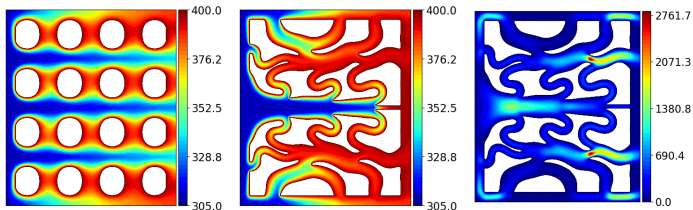
$$\|\mathbf{v}_0\|_\infty = 10$$

$$DP_0 = 1300$$

$$J_{final} = 4168$$

$$DP_{final} = 1188$$

With mini-
mum thick-
ness.



(a) T initial

(b) Optimized T

(c) Optimized $\|\mathbf{v}\|$

2D Air oil heat exchanger

Results :

Test case 9

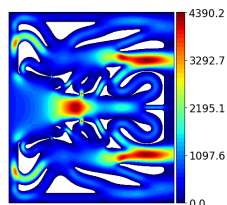
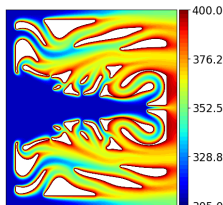
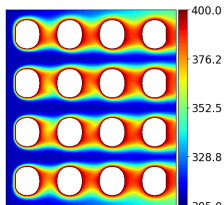
$$\|\mathbf{v}_0\|_\infty = 25$$

$$DP_0 = 1030$$

$$J_{final} = 7667$$

$$DP_{final} = 968$$

Without
minimum
thickness.



Test case 10

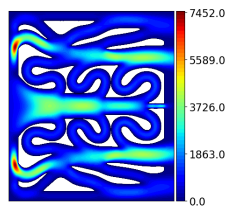
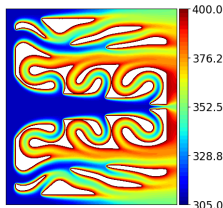
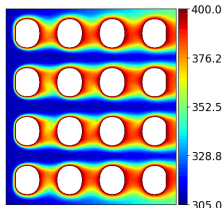
$$\|\mathbf{v}_0\|_\infty = 25$$

$$DP_0 = 1030$$

$$J_{final} = 7508$$

$$DP_{final} = 1112$$

With mini-
mum thick-
ness.



(a) T initial

(b) Optimized T

(c) Optimized $\|\mathbf{v}\|$

2D Air oil heat exchanger

Results :

Test case 11

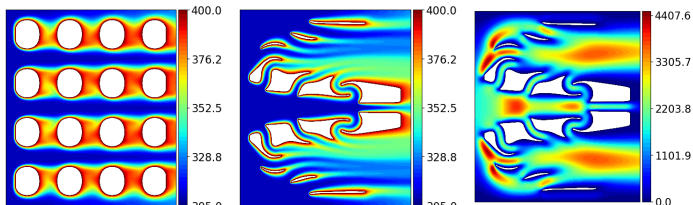
$$\|\mathbf{v}_0\|_\infty = 40$$

$$DP_0 = 475$$

$$J_{final} = 5731$$

$$DP_{final} = 479$$

Without
minimum
thickness.



Test case 12

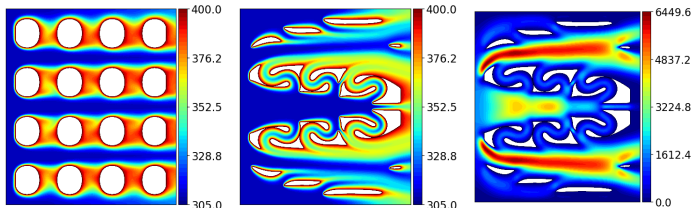
$$\|\mathbf{v}_0\|_\infty = 40$$

$$DP_0 = 475$$

$$J_{final} = 6847$$

$$DP_{final} = 524$$

With mini-
mum thick-
ness.



(a) T initial

(b) Optimized T

(c) Optimized $\|\mathbf{v}\|$

2D Air oil heat exchanger

Results :

Test case 16

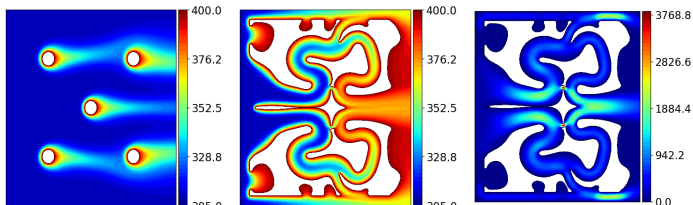
$$\|\mathbf{v}_0\|_\infty = 40$$

$$DP_0 = 475$$

$$J_{final} = 5731$$

$$DP_{final} = 479$$

Without
minimum
thickness.



Test case 17

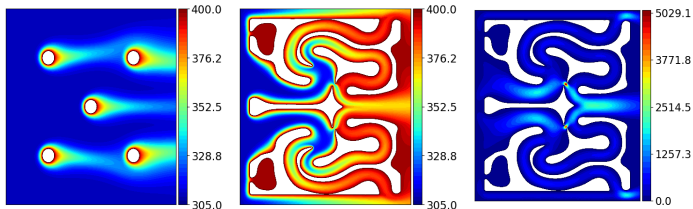
$$\|\mathbf{v}_0\|_\infty = 40$$

$$DP_0 = 475$$

$$J_{final} = 6847$$

$$DP_{final} = 524$$

With mini-
mum thick-
ness.



(a) T initial

(b) Optimized T

(c) Optimized $\|\mathbf{v}\|$