Body-fitted topology optimization of 2D and 3D fluid-to-fluid heat exchangers

Florian Feppon

PhD advisors : Grégoire Allaire, Charles Dapogny

Safran Tech advisors : Julien Cortial, Felipe Bordeu.

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- Non-penetration constraint:

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In 3D!

The boundary variation method of Hadamard





The boundary variation method of Hadamard



 $\Gamma_{\boldsymbol{\theta}} = (\boldsymbol{I} + \boldsymbol{\theta}) \Gamma, \text{ with } \boldsymbol{\theta} \in W^{1,\infty}_0(D, \mathbb{R}^d), \ ||\boldsymbol{\theta}||_{W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d)} < 1.$

The boundary variation method of Hadamard



$$J(\Gamma_{\boldsymbol{\theta}}) = J(\Gamma) + \frac{\mathrm{d}J}{\mathrm{d}\boldsymbol{\theta}}(\boldsymbol{\theta}) + o(\boldsymbol{\theta}), \quad \text{ with } \frac{|o(\boldsymbol{\theta})|}{||\boldsymbol{\theta}||_{W^{1,\infty}(D,\mathbb{R}^d)}} \xrightarrow{\boldsymbol{\theta} \to 0} 0.$$















Coupled physics system



Coupled physics system



► Incompressible Navier-Stokes system for the velocity and pressure (\mathbf{v}, p) in Ω_f $-\operatorname{div}(\sigma_f(\mathbf{v}, p)) + \rho \nabla \mathbf{v} \mathbf{v} = \mathbf{f}_f$ in Ω_f

• Convection-diffusion for the temperature T in Ω_f and Ω_s :

$$-\operatorname{div}(k_f \nabla T_f) + \rho \mathbf{v} \cdot \nabla T_f = Q_f \quad \text{in } \Omega_f$$
$$-\operatorname{div}(k_s \nabla T_s) = Q_s \quad \text{in } \Omega_s$$

Boundary conditions on $\Gamma = \partial \Omega_f$:

$$\begin{cases} T_f = T_s \text{ on } \Gamma \\ k_f \nabla T_f \cdot \boldsymbol{n} = k_s \nabla T_s \cdot \boldsymbol{n} \text{ on } \Gamma \\ \boldsymbol{v} = 0 \text{ on } \Gamma. \end{cases}$$

Proposition

Let $J(\Gamma, T, \mathbf{v}, p)$ an arbitrary cost function. If J has continuous partial derivatives, then $\Gamma \mapsto J(\Gamma, \mathbf{u}(\Gamma), T(\Gamma), \mathbf{v}(\Gamma), p(\Gamma))$ is shape differentiable and the shape derivative reads¹:

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}\theta} \Big[J(\Gamma_{\theta}, \mathbf{v}(\Gamma_{\theta}), p(\Gamma_{\theta}), T(\Gamma_{\theta}), \mathbf{u}(\Gamma_{\theta})) \Big](\theta) \\ &= \overline{\frac{\partial \mathfrak{J}}{\partial \theta}}(\theta) + \int_{\Gamma} (f_{f} \cdot \mathbf{w} - \sigma_{f}(\mathbf{v}, p) : \nabla \mathbf{w} + \mathbf{n} \cdot \sigma_{f}(\mathbf{w}, q) \nabla \mathbf{v} \cdot \mathbf{n} + \mathbf{n} \cdot \sigma_{f}(\mathbf{v}, p) \nabla \mathbf{w} \cdot \mathbf{n})(\theta \cdot \mathbf{n}) \mathrm{d}s \\ &+ \int_{\Gamma} \left(k_{s} \nabla T_{s} \cdot \nabla S_{s} - k_{f} \nabla T_{f} \cdot \nabla S_{f} + Q_{f} S_{f} - Q_{s} S_{s} - 2k_{s} \frac{\partial T_{s}}{\partial \mathbf{n}} \frac{\partial S_{s}}{\partial \mathbf{n}} + 2k_{f} \frac{\partial T_{f}}{\partial \mathbf{n}} \frac{\partial S_{f}}{\partial \mathbf{n}} \right) (\theta \cdot \mathbf{n}) \mathrm{d}s \end{split}$$

¹Feppon et al., Shape optimization of a coupled thermal fluid–structure problem in a level set mesh evolution framework (2019)

We rely on body fitted meshes^{2,3}.



^{$^{2}}Allaire, Dapogny, and Frey, Shape optimization with a level set based mesh evolution method (2014)$ </sup>

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 Fluid-Solid interface Γ exactly captured, no need of physics interpolation because no porous regions.



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Null space optimization algorithm

 Nonlinear constrained optimization on manifolds with a moderate number of constraints



Open source implementation⁴: https://gitlab.com/florian.feppon/null-space-optimizer pip install nullspace_optimizer

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Null space optimization algorithm

- Nonlinear constrained optimization on manifolds with a moderate number of constraints
- Generalization of the unconstrained gradient flow: no hard tuning of parameters



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Null space optimization algorithm

- Nonlinear constrained optimization on manifolds with a moderate number of constraints
- Generalization of the unconstrained gradient flow: no hard tuning of parameters
- Adapted to the infinite dimensional setting of the method of Hadamard



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Parallel computing

 Use of Domain Decomposition and adapted preconditioners for solving finite element problems : all FEM related operations are achieved in parallel⁵.



⁵Feppon et al., *Topology optimization of thermal fluid–structure systems using body-fitted meshes and parallel computing* (2020)

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- Use of Domain Decomposition and adapted preconditioners for solving finite element problems : all FEM related operations are achieved in parallel⁵.
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Parallel computing

- Use of Domain Decomposition and adapted preconditioners for solving finite element problems : all FEM related operations are achieved in parallel⁵.
- We solve fluid FEM problems on meshes up to 4.8 millions of Tetrahedra with 30 CPUs.
- Mesh adaptation and Isosurface discretization is still sequential. A future release of (Par)Mmg will allow to do it in parallel.



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Non mixing constraint



Non-penetration constraint:

 $d(\Omega_{f,\text{hot}},\Omega_{f,\text{cold}}) \geqslant d_{\min}.$

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Non mixing constraint



Non-penetration constraint: $d(\Omega_{f,\text{hot}},\Omega_{f,\text{cold}}) \ge d_{\min}$. We enforce it by imposing $\forall x \in \Omega_{f, \text{cold}}, \ d_{\Omega_{f, \text{hot}}}(x) \geqslant d_{\min},$ where $d_{\Omega_{f,hot}}$ is the signed

distance function to the domain $\Omega_{f,hot}$.

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The signed distance function

The signed distance function d_{Ω} to the domain $\Omega \subset D$ is defined by:

$$orall x \in D, \ d_{\Omega}(x) = \left\{ egin{array}{c} -\min_{y\in\partial\Omega} ||y-x|| & ext{if } x\in\Omega, \ \min_{y\in\partial\Omega} ||y-x|| & ext{if } x\in Dackslash \Omega. \end{array}
ight.$$





2D heat exchangers



Heat exchanger problem with limited pressure loss and non-mixing constraint:

$$\min_{\Gamma} \qquad J(\Omega_{f}) = -\left(\int_{\Omega_{f,cold}} \rho c_{\rho} \boldsymbol{v} \cdot \nabla T \, \mathrm{d}x - \int_{\Omega_{f,hot}} \rho c_{\rho} \boldsymbol{v} \cdot \nabla T \, \mathrm{d}x\right) \\ s.c. \begin{cases} \mathsf{DP}(\Omega_{f}) = \int_{\partial \Omega_{f}^{D}} \rho \mathrm{d}s - \int_{\partial \Omega_{f}^{N}} \rho \mathrm{d}s \leq \mathsf{DP}_{0} \\ Q_{hot \leftrightarrow cold}(\Omega_{f}) \geqslant d_{\min}. \end{cases}$$

2D heat exchangers



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$$\min_{\Gamma} \qquad J(\Omega_f) = -\left(\int_{\Omega_{f,cold}} \rho c_{\rho} \boldsymbol{v} \cdot \nabla T dx - \int_{\Omega_{f,hot}} \rho c_{\rho} \boldsymbol{v} \cdot \nabla T dx\right)$$

s.c.
$$\begin{cases} \mathsf{DP}(\Omega_f) = \int_{\partial \Omega_f^D} \rho ds - \int_{\partial \Omega_f^N} \rho ds \leq \mathsf{DP}_0 \\ Q_{hot \leftrightarrow cold}(\Omega_f) = \int_D j(d_{\Omega_{f,hot}}) dx \geq d_{\min}. \end{cases}$$

For instance

$$Q_{hot\leftrightarrow cold}(\Omega_f) := \left(\int_{\Omega_{f,cold}} \frac{1}{|d_{\Omega_{f,hot}}|^p} \mathrm{d}x \right)^{-\frac{1}{p}} \simeq \left| \left| \frac{1}{d_{\Omega_{f,hot}}} \right| \right|_{L^{\infty}(\Omega_{f,cold})}^{-1}$$

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$$Q_{hot\leftrightarrow cold}(\Omega_f) := \left(\int_{\Omega_{f,cold}} \frac{1}{|d_{\Omega_{f,hot}}|^p} \mathrm{d}x \right)^{-\frac{1}{p}} \simeq \left| \left| \frac{1}{d_{\Omega_{f,hot}}} \right| \right|_{L^{\infty}(\Omega_{f,cold})}^{-1}$$

This reduces to the setting of computing the shape derivative of some penalty functional $Q_{hot\leftrightarrow cold}(\Omega_f)$ with:

$$Q_{hot\leftrightarrow cold}(\Omega_f) := \int_D j(\mathbf{d}_{\Omega_{f,hot}}) \mathrm{d} x.$$

The shape derivative of $Q_{hot\leftrightarrow cold}(\Omega_f)$ is given by⁶:

•

$$Q'_{hot\leftrightarrow cold}(\Omega)(\boldsymbol{ heta}) = \int_{\partial\Omega_{f,hot}} u(y) \ \boldsymbol{ heta} \cdot \boldsymbol{n} \, \mathrm{d}y$$

with
$$u(y) = -\int_{z\in\operatorname{ray}(y)} j'(d_{\Omega_{f,hot}}(z)) \prod_{1\leq i\leq n-1} (1+\kappa_i(y)d_{\Omega_{f,hot}}(z)) \mathrm{d} z, \qquad \forall y\in\partial\Omega.$$

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The computation of u(y) requires a priori integration along the normal rays and the computation of curvatures $\kappa_i(y)$.

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It turns out that it is possible to compute u without integrating along the rays⁷:

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Let $\hat{u} \in V_{\omega}$ be the solution to the variational problem

$$\forall v \in V_{\omega}, \int_{\partial\Omega_{f,hot}} \hat{u}v \mathrm{d}s + \int_{D} \omega (\nabla d_{\Omega_{f,hot}} \cdot \nabla \hat{u}) (\nabla d_{\Omega_{f,hot}} \cdot \nabla v) \mathrm{d}x = -\int_{D} j' (d_{\Omega_{f,hot}})v \mathrm{d}x.$$

Then $u(y) = \hat{u}(y)$ for any $y \in \partial \Omega_{f,hot}$.

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- ▶ This variational problem can easily be solved with FEM in 2D and 3D
- This allows to handle conveniently geometric constraints (e.g. maximum thickness, minum distance, etc...) in 2D and 3D level set based topology optimization.

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2D Heat Exchangers with non-mixing constraint



(a) Initial temperature



(b) Final temperature.



(c) Intermediate iterations 0, 8, 20, 50, 88 et 200.



Figure: Schematic of the 3D setting.



Figure: Initial distribution of fluid considered for the 3D heat exchanger test case.







Figure: Intermediate iterations.



(a) Cold phase

(b) Hot phase

Figure: Separate plots of the topologically optimized cold and hot fluid phases in the configuration $d_{\min} = 0.04$.



Figure: Cut of the resulting solid domain

Many thanks for your attention!

