Introduction to FreeFEM. Numerical shape updates

Florian Feppon

Spring 2022 - Seminar for Applied Mathematics

ETH zürich

1. Finite Element Meshes - MMg

- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

1. Finite Element Meshes - MMg

2. Introduction to FreeFEM

3. Numerical shape updates

- 3.1 Moving mesh method
- 3.2 The level set method
- 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM

3. Numerical shape updates

- 3.1 Moving mesh method
- 3.2 The level set method
- 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

1. Finite Element Meshes - MMg

- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

Let Ω be a polyhedral connected open set of \mathbb{R}^N . A triangular/tetrahedral mesh is a set \mathcal{T} of non-degenerate *N*-simplices $(K_i)_{1 \le i \le n}$ which verify

1. $K_i \subset \overline{\Omega}$ and $\overline{\Omega} = \bigcup_{i=1}^n K_i$,



Let Ω be a polyhedral connected open set of \mathbb{R}^N . A triangular/tetrahedral mesh is a set \mathcal{T} of non-degenerate *N*-simplices $(K_i)_{1 \le i \le n}$ which verify

1. $K_i \subset \overline{\Omega}$ and $\overline{\Omega} = \bigcup_{i=1}^n K_i$,

2. for any $1 \le i, j \le n$, $K_i \cap K_j$ is a simplex whose vertices are also vertices of K_i and K_j .



Let Ω be a polyhedral connected open set of \mathbb{R}^N . A triangular/tetrahedral mesh is a set \mathcal{T} of non-degenerate *N*-simplices $(K_i)_{1 \le i \le n}$ which verify

1. $K_i \subset \overline{\Omega}$ and $\overline{\Omega} = \bigcup_{i=1}^n K_i$,

2. for any $1 \le i, j \le n$, $K_i \cap K_j$ is a simplex whose vertices are also vertices of K_i and K_j .



Let Ω be a polyhedral connected open set of \mathbb{R}^N . A triangular/tetrahedral mesh is a set \mathcal{T} of non-degenerate *N*-simplices $(K_i)_{1 \le i \le n}$ which verify

1. $K_i \subset \overline{\Omega}$ and $\overline{\Omega} = \bigcup_{i=1}^n K_i$,

2. for any $1 \le i, j \le n$, $K_i \cap K_j$ is a simplex whose vertices are also vertices of K_i and K_j .



Good meshes leading to accurate FEM analysis are without "flat" triangles.

Given a mesh \mathcal{T} with simplices $(K_i)_{1 \leq i \leq n}$, we call the space of \mathbb{P}_k -finite elements the finite-dimensional set

 $\mathcal{V}_h := \{ p \in \mathcal{C}(\bar{\Omega}) \, | \, p_{|K_i} \text{ is a polynomial of degree less than } k. \}$

Remark 1

A \mathbb{P}_1 function is piecewise linear on each triangle/tetrahedron and is characterized by its values at the node of the mesh.

One of the most easiest way to generate meshes is to use **level set** functions and **remeshing**.

Level set functions

Given a computational domain D, we can represent a subset $\Omega \subset D$ as the negative subdomain of a scalar "level set" function $\phi : D \to \mathbb{R}$:

$$\begin{cases} \phi(x) < 0 & \text{If } x \in \Omega, \\ \phi(x) = 0 & \text{If } x \in \partial\Omega, \\ \phi(x) > 0 & \text{If } x \in D \backslash \Omega. \end{cases}$$



▶ It is rather difficult to mesh arbitrary domain.

- It is rather difficult to mesh arbitrary domain.
- It is possible to mesh domains with explicitly parameterized boundaries



Figure: A triangular mesh with two elliptic boundaries

- It is rather difficult to mesh arbitrary domain.
- It is possible to mesh domains with explicitly parameterized boundaries



Figure: A triangular mesh with two elliptic boundaries

It is possible to improve the "quality" of meshes using metric tensors.



Figure: Improvement of the quality of a tetrahedral mesh (Figure from Dapogny, Dobrzynski and Frey, 2014).

It is possible to mesh negative subdomains of level set functions.



Figure: Meshing of a circular subdomain according to a level set function

It is possible to mesh negative subdomains of level set functions.



Figure: Meshing of a circular subdomain according to a level set function

We can do this using the library ${\tt Mmg}$

Elementary operations on sets can be easily performed with level-set representations:

 $\blacktriangleright \ \Omega_1 \cup \Omega_2 \quad \longleftrightarrow \quad \min(\phi_1, \phi_2)$

Elementary operations on sets can be easily performed with level-set representations:

- $\blacktriangleright \ \Omega_1 \cup \Omega_2 \quad \longleftrightarrow \quad \min(\phi_1, \phi_2)$
- $\blacktriangleright \ \Omega_1 \cap \Omega_2 \quad \longleftrightarrow \quad \mathsf{max}(\phi_1, \phi_2)$

Elementary operations on sets can be easily performed with level-set representations:

- $\blacktriangleright \ \Omega_1 \cup \Omega_2 \quad \longleftrightarrow \quad \mathsf{min}(\phi_1, \phi_2)$
- $\blacktriangleright \ \Omega_1 \cap \Omega_2 \quad \longleftrightarrow \quad \mathsf{max}(\phi_1, \phi_2)$
- $\blacktriangleright D \backslash \Omega \quad \longleftrightarrow \quad -\phi$

Elementary operations on sets can be easily performed with level-set representations:

- $\blacktriangleright \ \Omega_1 \cup \Omega_2 \quad \longleftrightarrow \quad \mathsf{min}(\phi_1, \phi_2)$
- $\blacktriangleright \ \Omega_1 \cap \Omega_2 \quad \longleftrightarrow \quad \mathsf{max}(\phi_1, \phi_2)$
- $\blacktriangleright D \backslash \Omega \quad \longleftrightarrow \quad -\phi$

Elementary operations on sets can be easily performed with level-set representations:

- $\blacktriangleright \ \Omega_1 \cup \Omega_2 \quad \longleftrightarrow \quad \mathsf{min}(\phi_1, \phi_2)$
- $\blacktriangleright \ \Omega_1 \cap \Omega_2 \quad \longleftrightarrow \quad \mathsf{max}(\phi_1, \phi_2)$

$$\blacktriangleright D \backslash \Omega \quad \longleftrightarrow \quad -\phi$$

Exercise: write the operation on level set which, given a material distribution $\Omega_f \subset D$ and a non optimizable region $\omega \subset D$, returns a new set Ω_f where the part $\Omega_f \cap \omega$ has been replaced with a new prescribed distribution \mathcal{X} .



Figure: Enforcing non-optimizable regions ω : the distribution of material Ω inside the domain ω should match exactly the red set \mathcal{X} .

1. Finite Element Meshes - MMg

2. Introduction to FreeFEM

3. Numerical shape updates

- 3.1 Moving mesh method
- 3.2 The level set method
- 3.3 Body-fitted meshes



FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics
- It allows to perform advanced operations on FEM structures



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics
- It allows to perform advanced operations on FEM structures
- It is interfaced with other powerful libraries, notably PETSc (linear algebra), MPI (parallel computing).



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics
- It allows to perform advanced operations on FEM structures
- It is interfaced with other powerful libraries, notably PETSc (linear algebra), MPI (parallel computing).
- there are regular fixes and updates



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics
- It allows to perform advanced operations on FEM structures
- It is interfaced with other powerful libraries, notably PETSc (linear algebra), MPI (parallel computing).
- there are regular fixes and updates



- FreeFEM is a powerful PDE solver (developped by F. Hecht, F. Nataf, P.-H. Tournier).
- It allows to solve PDEs with Finite Elements by writing their variational formulation in a spirit close to the mathematics
- It allows to perform advanced operations on FEM structures
- It is interfaced with other powerful libraries, notably PETSc (linear algebra), MPI (parallel computing).
- there are regular fixes and updates

FreeFEM allows in principle to solve large 3D FEM problems with millions of tetrahedral elements.

Some drawbacks:

- FreeFEM might not be devoid of bugs
- the documentation is sometimes incomplete
- it is not interfaced with more communly used industrial software (using CAD descriptions)

FreeFEM can be a bit delicate to install. Please refer to the page
 https://people.math.ethz.ch/~ffeppon/topopt_course/install_freefem.html
for the installation !

Let us write a solver for the heat conduction problem



Figure: A bi-material distribution of two conductive media with conductivity k_s and k_v . $\begin{cases} -\operatorname{div}(k_f \nabla T_f) = Q_f \\ -\operatorname{div}(k_s \nabla T_s) = Q_s \\ T = T_0 \end{cases}$ in Ω_f in Ω_s

$$T = T_0$$
 on $\partial \Omega^D_T$

$$-k_f \frac{\partial T_f}{\partial \boldsymbol{n}} = h \qquad \text{on } \partial \Omega_T^N \cap \partial \Omega_f$$

$$-k_s \frac{\partial I_s}{\partial \mathbf{n}} = h \qquad \text{on } \partial \Omega^N_T \cap \partial \Omega_s$$

$$T_f = T_s$$
 on Γ

$$-k_f \frac{\partial T_f}{\partial \boldsymbol{n}} = -k_s \frac{\partial T_s}{\partial \boldsymbol{n}} \qquad \text{on } \boldsymbol{\Gamma},$$
The variational formulation reads find $T \in T_0 + V_T(\Gamma)$ such that, for any $S \in V_T(\Gamma)$,

$$\int_{\Omega_s} k_s \nabla T \cdot \nabla S \mathrm{d}x + \int_{\Omega_f} k_f \nabla T \cdot \nabla S \mathrm{d}x = \int_{\Omega_s} Q_s S \mathrm{d}x + \int_{\Omega_f} Q_f S \mathrm{d}x + \int_{\partial \Omega_T^N} h S \mathrm{d}s.$$

where

$$V_{\mathcal{T}}(\Gamma) = \{ S \in H^1(D) \, | \, S = 0 \text{ on } \partial \Omega^D_{\mathcal{T}} \},\$$

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM

3. Numerical shape updates

- 3.1 Moving mesh method
- 3.2 The level set method
- 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

 $\min_{\Omega} J(\Omega)$

and that we need to do a Hadamard's shape update:

 $\Omega_{n+1} = (I + \theta_n)\Omega_n$

for a current shape Ω_n and vector field $\theta_n \in W^{1,\infty}(D,\mathbb{R}^d)$.

 $\min_{\Omega} J(\Omega)$

and that we need to do a Hadamard's shape update:

 $\Omega_{n+1} = (I + \theta_n)\Omega_n$

for a current shape Ω_n and vector field $\theta_n \in W^{1,\infty}(D,\mathbb{R}^d)$.

lf Ω_n has a mesh discretization, one can try **nodal displacements**

 $\min_{\Omega} J(\Omega)$

and that we need to do a Hadamard's shape update:

$$\Omega_{n+1} = (I + \theta_n)\Omega_n$$

for a current shape Ω_n and vector field $\theta_n \in W^{1,\infty}(D,\mathbb{R}^d)$.

- lf Ω_n has a mesh discretization, one can try **nodal displacements**
- lf Ω_n is described by a level set, one can solve an **advection equation**

 $\min_{\Omega} J(\Omega)$

and that we need to do a Hadamard's shape update:

$$\Omega_{n+1} = (I + \theta_n)\Omega_n$$

for a current shape Ω_n and vector field $\theta_n \in W^{1,\infty}(D,\mathbb{R}^d)$.

- If Ω_n has a mesh discretization, one can try **nodal displacements**
- lf Ω_n is described by a level set, one can solve an advection equation
- If Ω_n is described as a meshed subdomain, then one can use a hybrid method coupling the level set method and remeshing.

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

Very simple algorithm:

 $x_i \leftarrow x_i + \theta(x_i)$ for all nodes x_i



Figure: Discretization of a 2-d domain Ω in a simplicial mesh and its deformation by application of a displacement vector field θ .

- In general yields poor quality meshes.
- \blacktriangleright A refinement of the method: construct an extension $\widetilde{\theta}$ such that

$$\begin{cases} -\operatorname{div}(A\nabla\widetilde{\theta}) = 0 \text{ in } D, \\ \widetilde{\theta} = \theta \text{ on } \Gamma, \\ \widetilde{\theta} = 0 \text{ on } \partial\Omega\backslash\Gamma \end{cases}$$

for some positive definite matrix A, where Γ is the deformed interface and $\partial \Omega \backslash \Gamma$ are fixed interfaces.

- In general yields poor quality meshes.
- A refinement of the method: construct an extension $\widetilde{ heta}$ such that

$$\begin{cases} -\operatorname{div}(A\nabla\widetilde{\theta}) = 0 \text{ in } D, \\ \widetilde{\theta} = \theta \text{ on } \Gamma, \\ \widetilde{\theta} = 0 \text{ on } \partial\Omega\backslash\Gamma \end{cases}$$

for some positive definite matrix A, where Γ is the deformed interface and $\partial\Omega\backslash\Gamma$ are fixed interfaces.

Then do

$$x_i \leftarrow x_i + \widetilde{\theta}(x_i).$$

- In general yields poor quality meshes.
- \blacktriangleright A refinement of the method: construct an extension $\widetilde{ heta}$ such that

$$\begin{cases} -\operatorname{div}(A\nabla\widetilde{\theta}) = 0 \text{ in } D, \\ \widetilde{\theta} = \theta \text{ on } \Gamma, \\ \widetilde{\theta} = 0 \text{ on } \partial\Omega\backslash\Gamma \end{cases}$$

for some positive definite matrix A, where Γ is the deformed interface and $\partial\Omega\backslash\Gamma$ are fixed interfaces.

Then do

$$x_i \leftarrow x_i + \widetilde{\theta}(x_i).$$

This method does not allow to treat topological changes and yield very poor quality meshes after a few iterations.



Figure: Moving mesh method, figures from Allaire 2007.

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes

If the domain Ω is represented by a level set function:

$$\begin{cases} \phi(x) < 0 & \text{ If } x \in \Omega, \\ \phi(x) = 0 & \text{ If } x \in \partial\Omega, \\ \phi(x) > 0 & \text{ If } x \in D \backslash \Omega. \end{cases}$$

then the motion of a domain $\Omega(t)$ in *D* according to a vector field $\theta(x)$ can be captured by the solution $\phi(t, x)$ of the **advection equation**:

If the domain Ω is represented by a level set function:

$$\begin{cases} \phi(x) < 0 & \text{ If } x \in \Omega, \\ \phi(x) = 0 & \text{ If } x \in \partial\Omega, \\ \phi(x) > 0 & \text{ If } x \in D \backslash \Omega. \end{cases}$$

then the motion of a domain $\Omega(t)$ in *D* according to a vector field $\theta(x)$ can be captured by the solution $\phi(t, x)$ of the **advection equation**:

Indeed one can show that $\phi(x(t)) = \phi_0(x_0)$ for any trajectory x(t) satisfying

$$\begin{cases} \dot{x} = \theta(x) \\ x(0) = x_0. \end{cases}$$
(1)

If the domain Ω is represented by a level set function:

$$\begin{cases} \phi(x) < 0 & \text{If } x \in \Omega, \\ \phi(x) = 0 & \text{If } x \in \partial\Omega, \\ \phi(x) > 0 & \text{If } x \in D \backslash \Omega. \end{cases}$$

then the motion of a domain $\Omega(t)$ in *D* according to a vector field $\theta(x)$ can be captured by the solution $\phi(t, x)$ of the **advection equation**:

Indeed one can show that $\phi(x(t)) = \phi_0(x_0)$ for any trajectory x(t) satisfying

$$\begin{cases} \dot{x} = \theta(x) \\ x(0) = x_0. \end{cases}$$
 (1)

The new domain is then $\Omega(1) = \{x | \phi(1, x) < 0\} = \Phi^{\theta}(\Omega)$ where Φ^{θ} is the flow map diffeomorphism associated with eq. (1).

Remark: we don't have $\Phi^{\theta}(\Omega) = (I + \theta)\Omega$

Remark: we don't have $\Phi^{\theta}(\Omega) = (I + \theta)\Omega$ However one can show that for small θ ,

 $\Phi^{\boldsymbol{\theta}} = \boldsymbol{I} + \boldsymbol{\theta} + O(||\boldsymbol{\theta}||^2),$

Remark: we don't have $\Phi^{\theta}(\Omega) = (I + \theta)\Omega$ However one can show that for small θ ,

 $\Phi^{\theta} = I + \theta + O(||\theta||^2),$

hence the first order asymptotic shape calculus works with Φ^{θ} .

The power of the level set method:

capture easily topological changes on fixed meshes:



The power of the level set method:

capture easily topological changes on fixed meshes:



The power of the level set method:

capture easily topological changes on fixed meshes:



Its main drawback:

one needs to use an interpolating physical model if one relies only on a fixed meshes where the physical interfaces are only captured implicitly

the linear elasticity equations are interpolated by using an "ersatz material" for representing void:

$$\begin{cases} -\operatorname{div}(Ae(\boldsymbol{u})) = \boldsymbol{f}_{s} \text{ in } \Omega_{s} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = 0 \text{ on } \Gamma \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} & \longrightarrow \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases} \begin{cases} -\operatorname{div}(A(\Omega_{s})e(\boldsymbol{u})) = \boldsymbol{f} \text{ in } D \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases}$$

with $A(\Omega_s) = A \mathbb{1}_{\Omega_s} + \epsilon I \mathbb{1}_{D \setminus \Omega_s}$.

the linear elasticity equations are interpolated by using an "ersatz material" for representing void:

$$\begin{cases} -\operatorname{div}(Ae(\boldsymbol{u})) = \boldsymbol{f}_{s} \text{ in } \Omega_{s} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = 0 \text{ on } \Gamma \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases} \xrightarrow{} \begin{cases} -\operatorname{div}(A(\Omega_{s})e(\boldsymbol{u})) = \boldsymbol{f} \text{ in } D \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases}$$

with $A(\Omega_s) = A \mathbb{1}_{\Omega_s} + \epsilon I \mathbb{1}_{D \setminus \Omega_s}$.

Physically, "void" is replaced with a very soft material.

the linear elasticity equations are interpolated by using an "ersatz material" for representing void:

$$\begin{cases} -\operatorname{div}(Ae(\boldsymbol{u})) = \boldsymbol{f}_{s} \text{ in } \Omega_{s} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = 0 \text{ on } \Gamma \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases} \xrightarrow{} \begin{cases} -\operatorname{div}(A(\Omega_{s})e(\boldsymbol{u})) = \boldsymbol{f} \text{ in } D \\ \boldsymbol{u} = \boldsymbol{u}_{0} \text{ on } \partial \Omega_{s}^{D} \\ Ae(\boldsymbol{u}) \cdot \boldsymbol{n} = \boldsymbol{g} \text{ on } \partial \Omega_{s}^{N} \end{cases}$$

with $A(\Omega_s) = A1_{\Omega_s} + \epsilon I 1_{D \setminus \Omega_s}$.

- Physically, "void" is replaced with a very soft material.
- Since the physical model is different, different formulas for the shape derivative must be implemented.

the Navier Stokes equations are interpolated by using the Brinkmann porous flow model:

$$\begin{pmatrix} -\operatorname{div}(\sigma_f(\boldsymbol{v},\boldsymbol{p})) + \rho \nabla \boldsymbol{v} \, \boldsymbol{v} = \boldsymbol{f}_f & \text{in } \Omega_f \\ \operatorname{div}(\boldsymbol{v}) = 0 & \text{in } \Omega_f \\ \boldsymbol{v} = \boldsymbol{v}_0 & \text{on } \partial \Omega_f^D \\ \sigma_f(\boldsymbol{v},\boldsymbol{p})\boldsymbol{n} = 0 & \text{on } \partial \Omega_f^N \\ \boldsymbol{v} = \boldsymbol{0} & \text{on } \Gamma, \end{cases}$$

the Navier Stokes equations are interpolated by using the Brinkmann porous flow model:

$$\begin{cases} -\operatorname{div}(\sigma_f(\boldsymbol{v},\boldsymbol{p})) + \rho \nabla \boldsymbol{v} \, \boldsymbol{v} + \boldsymbol{\alpha}(\Omega_f) \, \boldsymbol{v} = \boldsymbol{f}_f & \text{in } D \\ \operatorname{div}(\boldsymbol{v}) = 0 & \text{in } D \\ \boldsymbol{v} = \boldsymbol{v}_0 & \text{on } \partial \Omega_f^D \\ \sigma_f(\boldsymbol{v},\boldsymbol{p}) \boldsymbol{n} = 0 & \text{on } \partial \Omega_f^N \end{cases}$$

with $\alpha(\Omega_f) = tgv \mathbf{1}_{D \setminus \Omega_f}$ for some large value tgv.

Physically, the solid material is replaced with a slightly porous material.

when using the advection equation, topological changes are the outcome of some numerical diffusion.

when using the advection equation, topological changes are the outcome of some numerical diffusion. These are not mathematically captured by the notion of shape derivatives and happen "fortunately" in this method

- when using the advection equation, topological changes are the outcome of some numerical diffusion. These are not mathematically captured by the notion of shape derivatives and happen "fortunately" in this method
- "true topological changes" are permitted when solving the Hamilton-Jacobi equation:

$$\frac{\partial \phi}{\partial t} + \mathbf{v} |\nabla \phi| = \mathbf{0}$$

and considering viscosity solutions.

- when using the advection equation, topological changes are the outcome of some numerical diffusion. These are not mathematically captured by the notion of shape derivatives and happen "fortunately" in this method
- "true topological changes" are permitted when solving the Hamilton-Jacobi equation:

$$\frac{\partial \phi}{\partial t} + \mathbf{v} |\nabla \phi| = \mathbf{0}$$

and considering viscosity solutions.

Appropriate numerical schemes are required, which are quite delicate to implement on triangular meshes.

- when using the advection equation, topological changes are the outcome of some numerical diffusion. These are not mathematically captured by the notion of shape derivatives and happen "fortunately" in this method
- "true topological changes" are permitted when solving the Hamilton-Jacobi equation:

$$\frac{\partial \phi}{\partial t} + \mathbf{v} |\nabla \phi| = \mathbf{0}$$

and considering viscosity solutions.

Appropriate numerical schemes are required, which are quite delicate to implement on triangular meshes.

The advection equation can be solved on triangular/tetrahedral meshes with the method of characteristics.

- when using the advection equation, topological changes are the outcome of some numerical diffusion. These are not mathematically captured by the notion of shape derivatives and happen "fortunately" in this method
- "true topological changes" are permitted when solving the Hamilton-Jacobi equation:

$$\frac{\partial \phi}{\partial t} + \mathbf{v} |\nabla \phi| = \mathbf{0}$$

and considering viscosity solutions.

Appropriate numerical schemes are required, which are quite delicate to implement on triangular meshes.

- The advection equation can be solved on triangular/tetrahedral meshes with the method of characteristics.
- A common practice is to initialize the level-set to the signed distance function to the domain.
The signed distance function to a bounded open domain $\Omega \subset \mathbb{R}^d$ is the function

$$d_{\Omega} \,:\, \mathbb{R}^d o \mathbb{R}$$

defined for any $x \in \mathbb{R}^d$ by

$$d_{\Omega}(x) = \begin{cases} -\inf_{y \in \partial \Omega} ||x - y|| & \text{ if } x \in \Omega, \\ \inf_{y \in \partial \Omega} ||x - y|| & \text{ if } x \notin \Omega. \end{cases}$$



(a) Meshed subdomain $\Omega \subset D$ (in blue) of a computational domain D.



(b) Isocontours of the signed distance function d_{Ω} .

Figure: Example of signed distance function d_{Ω} numerically computed on a meshed domain.



Figure: 3-d plot of d_{Ω} .

A fundamental property: ∇d_{Ω} is an extension of the outward normal to d_{Ω} constant along the rays, in particular:

• $|\nabla d_{\Omega}(x)| = 1$ for any $x \in D$ where d_{Ω} is differentiable;

A fundamental property: ∇d_{Ω} is an extension of the outward normal to d_{Ω} constant along the rays, in particular:

- $|\nabla d_{\Omega}(x)| = 1$ for any $x \in D$ where d_{Ω} is differentiable;
- $\nabla d_{\Omega}(y) = \mathbf{n}(y)$ for any $y \in \partial \Omega$.

A fundamental property: ∇d_{Ω} is an extension of the outward normal to d_{Ω} constant along the rays, in particular:

- $|\nabla d_{\Omega}(x)| = 1$ for any $x \in D$ where d_{Ω} is differentiable;
- $\nabla d_{\Omega}(y) = \mathbf{n}(y)$ for any $y \in \partial \Omega$.

A fundamental property: ∇d_{Ω} is an extension of the outward normal to d_{Ω} constant along the rays, in particular:

- $|\nabla d_{\Omega}(x)| = 1$ for any $x \in D$ where d_{Ω} is differentiable;
- $\nabla d_{\Omega}(y) = \mathbf{n}(y)$ for any $y \in \partial \Omega$.

A fundamental property: ∇d_{Ω} is an extension of the outward normal to d_{Ω} constant along the rays, in particular:

- $|\nabla d_{\Omega}(x)| = 1$ for any $x \in D$ where d_{Ω} is differentiable;
- $\nabla d_{\Omega}(y) = \mathbf{n}(y)$ for any $y \in \partial \Omega$.

It also holds that $\Delta d_{\Omega} = \kappa$ on $\partial \Omega$ where κ is the mean curvature of $\partial \Omega$.





 There exist a number of numerical algorithms for computing d_Ω of a mesh subdomain (notably, the Fast Marching Method)

- There exist a number of numerical algorithms for computing d_Ω of a mesh subdomain (notably, the Fast Marching Method)
- From a given level set function $\phi_0(x)$, solving the **reinitialization** equation

$$\left\{egin{aligned} \partial_t \phi + \mathrm{sign}(\phi)(|
abla \phi|-1) &= 0 \ \phi(0,x) &= \phi_0(x). \end{aligned}
ight.$$

allows to transform ϕ_0 into the signed distance function to the domain $\Omega = \{x \mid \phi_0(x) < 0\}.$



(a) Initial design

(b) Intermediate design

(c) Final design

Figure: Topology optimization with the level set method

- 1. Finite Element Meshes MMg
- 2. Introduction to FreeFEM
- 3. Numerical shape updates
 - 3.1 Moving mesh method
 - 3.2 The level set method
 - 3.3 Body-fitted meshes



Domain interface Γ exactly captured.



- Domain interface Γ exactly captured.
- Mesh size control is possible.



- Domain interface Γ exactly captured.
- Mesh size control is possible.



- Domain interface Γ exactly captured.
- Mesh size control is possible.



- Domain interface Γ exactly captured.
- Mesh size control is possible.



- Domain interface Γ exactly captured.
- Mesh size control is possible.
- Less prone to numerical diffusion, allows to capture fine details.



change



(d) Final design

Figure: Level-set based mesh evolution method (figures from Dapogny et. al., 2013).