Lecture 12: shape optimization with geometric constraints.

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Given a Lipschitz domain Ω , we parameterize deformations of Ω by a continuous vector field θ :

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Figure: Deformation of a domain Ω with the method of Hadamard. A small vector field θ is used to deform Ω into $\Omega_{\theta} = (I + \theta)\Omega$.

Shape/Topology optimization is the mathematical art of generating shapes that best fulfill a proposed objective.

Generically, a design optimization problem arises under the form

$$egin{aligned} \min_{\Omega \subset D} & J(\Omega) \ s.t. egin{cases} G_i(\Omega) &= 0, & 1 \leq i \leq p \ H_j(\Omega) \leq 0, & 1 \leq j \leq q \end{aligned}$$

where

- Ω is an **open domain** sought to be optimized
- ► *J* is an **objective function** to minimize (corresponding to a measure of the performance)
- ► G_i and H_j are respectively p and q equality and inequality constraints (corresponding e.g. to industrial specifications to meet)

Today: how to take into account geometric constraints, e.g.:



minimum thickness

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- minimum thickness
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- minimum distance between to connected components



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- minimum thickness
- maximum thickness
- minimum distance between to connected components
- minimum angle with respect to a direction (overhang)



1. The signed distance function

- 2. Formulation of geometric constraints
- 3. Shape derivatives of geometric constraints
- 4. Numerical examples

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The signed distance function d_{Ω} to the domain $\Omega \subset D$ is defined by:

$$\forall x \in D, \ d_{\Omega}(x) = \begin{cases} -\min_{y \in \partial \Omega} ||y - x|| & \text{if } x \in \Omega, \\ \min_{y \in \partial \Omega} ||y - x|| & \text{if } x \in D \setminus \Omega. \end{cases}$$



The signed distance function

An example: a meshed subdomain $\Omega \subset D$



An example: the signed distance function d_{Ω} :



Definition 2 (Skeleton set and projection)

1. The set of points $x \in \mathbb{R}^d$ for which the minimization problem

$$\min_{y \in \partial \Omega} ||x - y|| \tag{1}$$

admits several minimizers is called the *skeleton* of Ω and is denoted by Σ .

Definition 2 (Skeleton set and projection)

For any x ∈ ℝ^d\Σ, the unique minimizer of eq. (1) is denoted p_{∂Ω}(x) and is called the (orthogonal) projection of x onto ∂Ω, in that case it holds

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The signed distance function



Proposition 1 (Differentiability of d_{Ω})

Assume Ω is a \mathcal{C}^1 domain with outward normal **n**.

• The signed distance function d_{Ω} is differentiable at any point $x \in \mathbb{R}^d \setminus \Sigma$, and it is not differentiable on Σ .

An example: the signed distance function d_{Ω} :



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- The gradient ∇d_{Ω} is an extension of the unit normal vector **n** to $\partial \Omega$ pointing outward Ω :

$$\forall x \in \mathbb{R}^d \setminus \Sigma, \, \nabla d_{\Omega}(x) = \boldsymbol{n}(p_{\partial \Omega}(x)).$$

An example: the gradient of the signed distance function ∇d_{Ω} :



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$$\forall x \in \mathbb{R}^d \setminus \Sigma, \, \nabla d_{\Omega}(x) = \mathbf{n}(p_{\partial \Omega}(x)).$$

• In particular, d_{Ω} solves the so-called "Eikonal" equation:

$$\left\{ egin{array}{l} ||
abla d_{\Omega}|| = 1 \ in \ \mathbb{R}^d ackslash \Sigma, \ d_{\Omega} = 0 \ on \ \partial\Omega. \end{array}
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The ray emerging from y is defined to be the one-dimensional segment

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Proposition 2

For any $y \in \partial \Omega$, define $\zeta_{-}(y)$ and $\zeta_{+}(y)$ the distance at which the ray hits the boundary of D or the skeleton:

$$\begin{aligned} \forall y \in \partial \Omega, \ \zeta_+(y) &= \sup\{s \ge 0 \mid \{y + t \nabla d_\Omega(y) \mid t \in [0, s)\} \cap (\overline{\Sigma} \cup \partial D) = \emptyset\}, \\ \forall y \in \partial \Omega, \ \zeta_-(y) &= \inf\{s \le 0 \mid \{y + t \nabla d_\Omega(y) \mid t \in (s, 0]\} \cap (\overline{\Sigma} \cup \partial D) = \emptyset\}. \end{aligned}$$

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Then we also have

$$ray(y) = \{y + sn(y) \mid \zeta_{-}(y) < s < \zeta_{+}(y)\}.$$

The signed distance function



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Geometric constraints are often point-wise constraints formulated from the signed distance function, e.g.

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- Sometimes, enforcing strictly the constraint might not be desirable because the feasible region becomes tight, or because it may prevent topological changes (such as minimum thickness).

A trade-off can be achieved by setting the value desired for the objective function as a constraint , e.g.:

$$egin{array}{ll} \min_{\Omega\subset D} & P(\Omega) \ s.t. & J(\Omega) \leq 0.9 \mathrm{J}_{opt} \end{array}$$

where J_{opt} would be the optimal value without the constraint.

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This means taking into account the averaged constraint

$$P_{\mathsf{max}}(\Omega) \leq d_{\mathsf{max}} ext{ with } P_{\mathsf{max}}(\Omega) := 2 \left(rac{1}{|\Omega|} \int_{\Omega} |d_{\Omega}|^p \mathrm{d}x
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Maximum thickness



(a) No maximum thickness constraint

(b) $d_{\rm max} = 0.07$.

Figure: Maximum thickness constraint for 2D arch.

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 - 1. ζ_{-} is not differentiable with respect to the shape
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 - 3. Enforcing the minimum thickness at all iterations would prevent topological changes to occur.
- It is better to rely on a more flexible formulation

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- In order to find a good compromise between thickness and the original optimization problem, we use the reformulation

$$\min_{\Omega \subset D} \quad P_{\min}(\Omega) \\ s.t. \quad J(\Omega) \leq \alpha J_{opt}$$

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the minimization will find shapes with good performances and which satisfy approximately the minimum thickness constraint.

Minimum thickness



Figure: Minimum thickness constraint for 2D cantilever.

Distance constraint:



An application: heat-exchangers:



Distance constraint:

$$d(\Omega_{f,\mathsf{hot}},\Omega_{f,\mathsf{cold}})\geqslant d_{\mathsf{min}}.$$

Figure: Minimum distance constraint for two-fluid heat-exchangers.

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Distance constraint:

 $d(\Omega_{f,hot},\Omega_{f,cold}) \geqslant d_{min}.$

We enforce it by imposing

 $\forall x \in \Omega_{f, cold}, \ d_{\Omega_{f, hot}}(x) \geqslant d_{\min},$

where $d_{\Omega_{f,hot}}$ is the signed distance function to the domain $\Omega_{f,hot}$.

$$\forall x \in \Omega_{f, cold}, \ d_{\Omega_{f, hot}}(x) \geqslant d_{\min},$$

This constraint can be equivalently formulated as

$$\left|\left|\frac{1}{d_{\Omega_{f,\mathsf{hot}}}}\right|\right|_{L^{\infty}(\Omega_{f,\mathsf{cold}})} \leq \frac{1}{d_{\mathsf{min}}} \Leftrightarrow \left|\left|\frac{1}{d_{\Omega_{f,\mathsf{hot}}}}\right|\right|_{L^{\infty}(\Omega_{f,\mathsf{cold}})}^{-1} \geqslant d_{\mathsf{min}}.$$

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We can approximate it by

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$$Q_{\mathsf{hot} o \mathsf{cold}}(\Omega_f) := \left(\int_{\Omega_{f,\mathsf{cold}}} rac{1}{|d_{\Omega_{f,\mathsf{hot}}}|^p} \mathrm{d} x
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$$\forall y \in \partial \Omega, \mathbf{n}(y) \cdot \mathbf{e}_y = \cos(\pi/2 + \theta) = -\sin(\theta) \ge -\sin(\beta).$$

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Remembering that $\mathbf{n}(y) = \nabla d_{\Omega(y)}$, we can formulate this in terms of d_{Ω} :

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This is equivalent to

$$\nabla d_{\Omega} \cdot \boldsymbol{e}_{y} + \sin \beta \geq 0,$$

for instance

$$P_{\beta}(\Omega) = 0 \text{ with } P_{\beta}(\Omega) := \int_{\Omega} \min \left(\nabla d_{\Omega} \cdot \boldsymbol{e}_{y} + \sin(\beta), 0 \right)^{2} \mathrm{d}x.$$

Maximum thickness:

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Maximum overhang:

$$P_{\beta}(\Omega) = 0 \text{ with } P_{\beta}(\Omega) := \int_{\Omega} \min \left(\nabla d_{\Omega} \cdot \boldsymbol{e}_{y} + \sin(\beta), 0 \right)^{2} \mathrm{d}x.$$

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$$P'(\Omega)(\boldsymbol{ heta}) = \int_{D\setminus\overline{\Sigma}} j'(d_{\Omega}(x))d'_{\Omega}(\boldsymbol{ heta})(x)\mathrm{d}x$$

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$$P'(\Omega)(\boldsymbol{\theta}) = \int_{D\setminus\overline{\Sigma}} j'(d_{\Omega}(x))d'_{\Omega}(\boldsymbol{\theta})(x)\mathrm{d}x$$

In order to compute the shape derivative of $P(\Omega)$, we need to compute the shape derivative $d'_{\Omega}(\theta)$.

For any $x \notin \Sigma$, the map $\theta \mapsto d_{(I+\theta)\Omega}(x)$ is Gâteaux-differentiable at θ as an application from $W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d)$ into R^d and its derivative reads

 $d'_{\Omega}(\boldsymbol{\theta})(x) = -\boldsymbol{\theta}(p_{\partial\Omega}(x)) \cdot \boldsymbol{n}(p_{\partial\Omega}(x)).$

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Equivalently, $d'_{\Omega}(\theta)$ is characterized by the boundary value problem



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$$P'(\Omega)(\boldsymbol{\theta}) = \int_{D\setminus\overline{\Sigma}} j'(d_{\Omega}(x))d'_{\Omega}(\boldsymbol{\theta})(x)\mathrm{d}x = -\int_{D\setminus\overline{\Sigma}} j'(d_{\Omega}(x))\boldsymbol{\theta}(\boldsymbol{p}_{\partial\Omega}(x))\cdot\boldsymbol{n}(\boldsymbol{p}_{\partial\Omega}(x))\mathrm{d}x.$$

The composition with the projection $p_{\partial\Omega}$ is not very easy to implement.

For any $oldsymbol{ heta} \in W^{1,\infty}(D,\mathbb{R}^d)$, we have

$$P'(\Omega)(\boldsymbol{\theta}) = -\int_{D\setminus\overline{\Sigma}} j'(d_{\Omega}(x))\boldsymbol{\theta}(p_{\partial\Omega}(x)) \cdot \boldsymbol{n}(p_{\partial\Omega}(x)) \mathrm{d}x$$

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$$= -\int_{\partial\Omega} \left(\int_{z \in ray(y)} j'(d_{\Omega}(z)) \prod_{i=1}^{d-1} (1 + \kappa_i(y)d_{\Omega}(z)) dz \right) \theta(y) \cdot \boldsymbol{n}(y) d\sigma(y)$$

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= $\int_{\partial\Omega} u(y)\theta(y) \cdot \boldsymbol{n}(y) d\sigma(y)$
with $u(y) = -\int_{z \in ray(y)} j'(d_{\Omega}(z)) \prod_{i=1}^{d-1} (1 + \kappa_i(y)d_{\Omega}(z)) dz.$

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Computing *u* requires:

1. Integrating along rays on the discretization mesh:





2. Estimating the principal curvatures $\kappa_i(y)$.

These two operations are quite delicate to implement !

It turns out that it is possible to compute u without integrating along the rays:

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Let $\hat{u} \in V_{\omega}$ be the solution to the variational problem

$$\forall \mathbf{v} \in V_{\omega}, \int_{\partial\Omega} \hat{u} \mathbf{v} \mathrm{d} \mathbf{s} + \int_{D} \omega (\nabla d_{\Omega} \cdot \nabla \hat{u}) (\nabla d_{\Omega} \cdot \nabla \mathbf{v}) \mathrm{d} \mathbf{x} = -\int_{D} j'(d_{\Omega}) \mathbf{v} \mathrm{d} \mathbf{x},$$

Then $u(y) = \hat{u}(y)$ for any $y \in \partial \Omega$.

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 $\omega>0$ is a weight that can be chosen rather arbitrarily.

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This variational problem can easily be solved with FEM in 2D and 3D!

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This variational problem can easily be solved with FEM in 2D and 3D!

Sketch of the proof.

Take $v = -d'_{\Omega}(\boldsymbol{\theta})$. Since $\nabla v \cdot \nabla d_{\Omega} = 0$ and $v = \boldsymbol{\theta} \cdot \boldsymbol{n}$ on $\partial \Omega$, one finds

$$\int_{\partial\Omega} \hat{\boldsymbol{u}} \boldsymbol{\theta} \cdot \boldsymbol{n} \mathrm{d}\boldsymbol{s} = \int_{D} j'(\boldsymbol{d}_{\Omega}) \boldsymbol{d}_{\Omega}'(\boldsymbol{\theta}) \mathrm{d}\boldsymbol{x}.$$

whence $\hat{u} = u$ on $\partial \Omega$.

Recap:

The shape derivative of a geometric constraint of the form

$$P(\Omega) = \int_D j(d_{\Omega}(x)) \mathrm{d}x$$

reads

$$P'(\Omega)(\boldsymbol{\theta}) = \int_D j'(\boldsymbol{d}_\Omega) \boldsymbol{d}'_\Omega(\boldsymbol{\theta}) \mathrm{d}x = \int_{\partial\Omega} u \boldsymbol{\theta} \cdot \boldsymbol{n} \mathrm{d}s.$$

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- 1. The signed distance function
- 2. Formulation of geometric constraints
- 3. Shape derivatives of geometric constraints
- 4. Numerical examples



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Optimized shape without maximum thickness constraint (max $d_{\Omega} = 0.36$).

Figure: MBB beam without thickness constraint



(c) Optimized shape with maximum thickness constraint.

Figure: MBB beam with maximum thickness constraint



(a) Optimized shape without minimum thickness constraint.

Figure: MBB beam without minimum thickness constraint



(b) Optimized shape with minimum thickness constraint $(d_{\min} = 0.1)$.

Figure: MBB beam with thickness constraint



(c) Optimized shape with minimum thickness constraint $(d_{\min} = 0.2)$.

Figure: MBB beam with minimum thickness constraint



Figure: Settings of the heat exchanger topology optimization problem .

Non-penetration constraint:

 $d(\Omega_{f, \mathsf{hot}}, \Omega_{f, \mathsf{cold}}) \geqslant d_{\mathsf{min}}.$

We enforce it by imposing

 $\forall x \in \Omega_{f, cold}, \ d_{\Omega_{f, hot}}(x) \geqslant d_{\min},$

where $d_{\Omega_{f,hot}}$ is the signed distance function to the domain $\Omega_{f,hot}$.





Figure: Schematic of the 3D setting.


Figure: Initial distribution of fluid considered for the 3D heat exchanger test case.







Figure: Separate plots of the topologically optimized cold and hot fluid phases in the configuration $d_{\min} = 0.04$.

Numerical examples



Figure: Cut of the resulting solid domain