

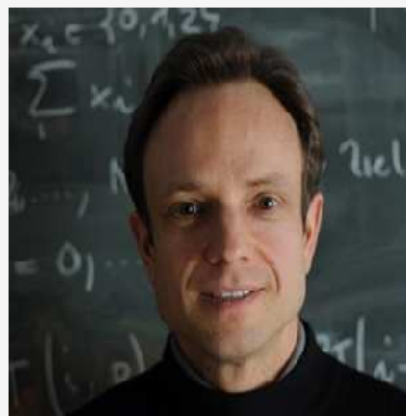
Course Video

Section 1.2.2: Electrostatic Fields

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(C) Seminar für Angewandte Mathematik, ETH Zürich



Prerequisites.

- Derivatives in 1D and 2D, in particular the gradient, see [Lecture → § 0.1.2.18]
- Piecewise continuous functions [Lecture → Section 0.1.2.3]
- Integration in 1D and 2D [Lecture → Section 0.1.2.5]

Dependency. Unit on [Lecture → Section 1.2.1] should be covered before.

Note: Possible minor *mismatch of video and tablet notes!*

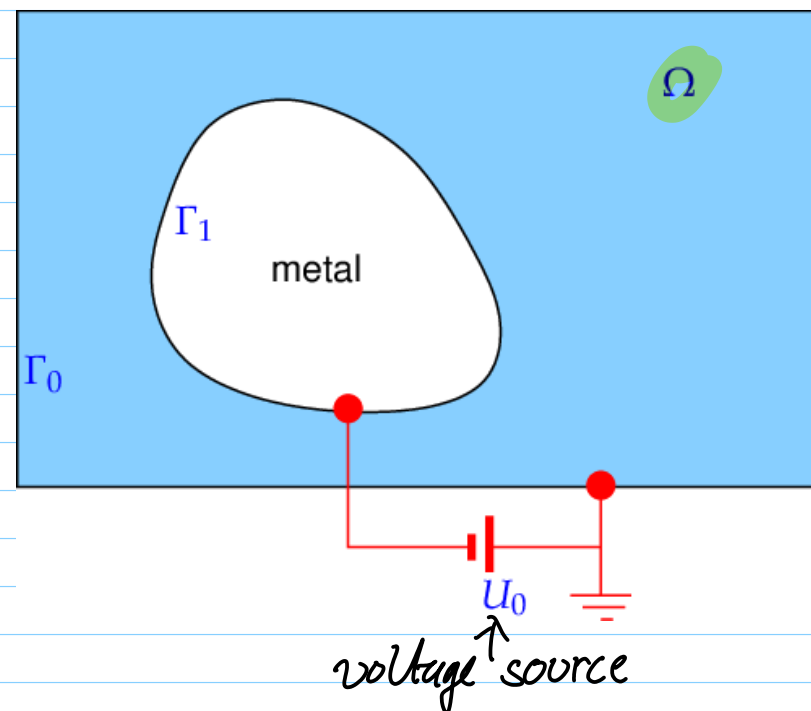
[Corrections and updates can be incorporated into tablet notes only]

I. Second-Order Scalar Elliptic Boundary Value Problems

1.2 Equilibrium Models: Examples

1.2.1. Elastic Membranes

1.2.2. Electrostatic Fields



← conducting box

Sought: **Electric field**

$$E: \Omega \subset \mathbb{R}^d \rightarrow \mathbb{R}^d$$

$$d = 3(2)$$

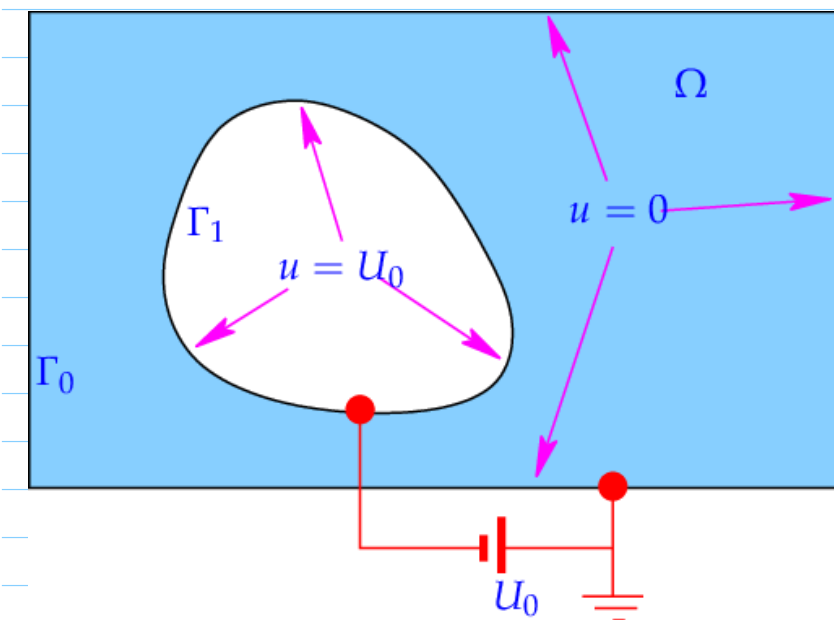
Maxwell equ.:

$$E = -\text{grad } \mu$$

scalar potential $\Omega \rightarrow \mathbb{R}$
(main unknown)



② BDC: Surfaces of conductor are equipotential surfaces,
($\mu \equiv \text{const}$)



$$\begin{aligned} u &= 0 \text{ on } T_0 \\ u &= U_0 \text{ on } T_1 \end{aligned}$$

▷ (Tentative) def.-spec.
 $\hat{V}_E = \{ u \in C^0(\overline{\Omega}) : \begin{aligned} &u = 0 \text{ on } T_0 \\ &u = U_0 \text{ on } T_1 \end{aligned} \}$
 ↑
 Meaningful gradient?

Physics:

Electromagnetic field energy: (electrostatic setting with scalar potential u)

$$J_E(u) = \frac{1}{2} \int_{\Omega} (\epsilon(x) \mathbf{grad} u(x)) \cdot \mathbf{grad} u(x) dx, \quad (1.2.2.6)$$

where $\epsilon : \Omega \mapsto \mathbb{R}^{3,3}$ is the **dielectric tensor**, $\epsilon(x)$ is symmetric, with units $[\epsilon] = \frac{\text{As}}{\text{Vm}}$.

$\underline{\epsilon} \hat{=}$ macroscopic material coefficient

$$(\epsilon(x) \mathbf{grad} u(x)) \cdot \mathbf{grad} u(x) = \sum_{i=1}^3 \sum_{j=1}^3 \epsilon_{ij}(x) \frac{\partial u}{\partial x_i}(x) \frac{\partial u}{\partial x_j}(x).$$

Essential: $J_E(u) \geq 0$

$\Rightarrow \underline{\epsilon}$ is **uniformly positive definite**

$$\exists 0 < \epsilon^- \leq \epsilon^+ < \infty: \epsilon^- \|z\|^2 \leq (\epsilon(x)z) \cdot z \leq \epsilon^+ \|z\|^2 \quad \forall z \in \mathbb{R}^3, \forall x \in \Omega. \quad (1.2.2.8)$$

[apply with $z := \mathbf{grad} u$]

Meaningful potentials must allow evaluation of $J_E(u)$

Configuration space for electrostatic field problems

As configuration space for scalar potentials we use the affine space

$$\hat{V}_E := \{ u \in C_{pw}^1(\overline{\Omega}), u \text{ satisfies (1.2.2.4)} \}. \quad (1.2.2.13)$$

Note: $\hat{V}_E \subset C^0(\overline{\Omega})$

Equilibrium condition for electrostatics

The physical field will attain minimal electromagnetic field energy

$$u_* = \underset{u \in \hat{V}_E}{\operatorname{argmin}} J_E(u). \quad (1.2.2.16)$$

③ Review questions 1.2.2.17

A:

Explain, why surfaces of conducting bodies are equipotential surfaces in electrostatics.

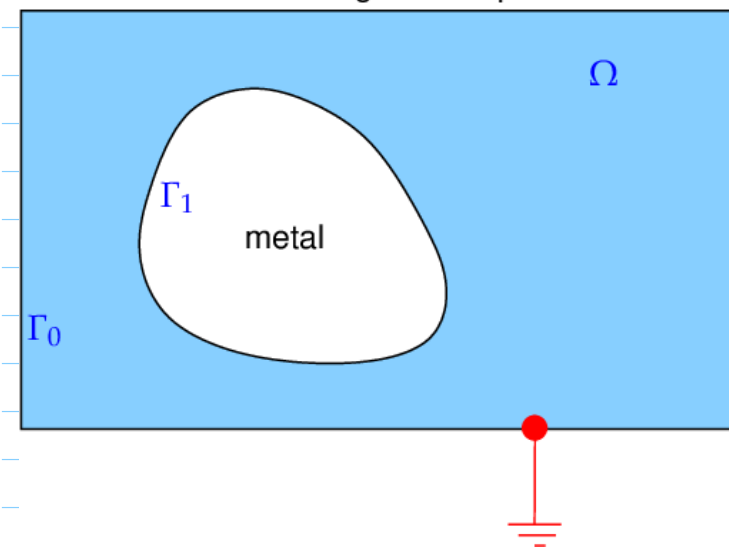
B:

What are the physical units for the electric scalar potential u , the electric field \mathbf{E} , the dielectric tensor ϵ , and the electromagnetic field energy

$$J_E(u) = \frac{1}{2} \int_{\Omega} (\epsilon(x) \mathbf{grad} u(x)) \cdot \mathbf{grad} u(x) dx \quad (1.2.2.6)$$

C:

What is a suitable configuration space for describing electrostatics in the following configuration:



◁ A metallic (conducting) body is located inside a (grounded) metal box, but it is not connected to any wire, which means that its scalar potential is not known a priori. This situation is often referred to as **floating potential**.

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