

ETH Lecture 401-0674-00L Numerical Methods for Partial Differential Equations

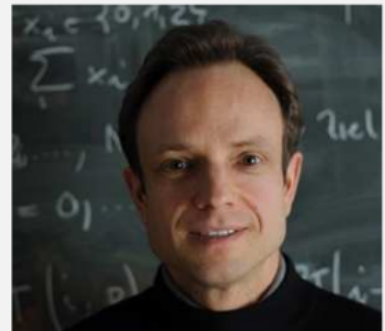
Course Video

Section 11.6.3: Finite Volume Methods (FVMs) for Systems of Conservation Laws

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(C) Seminar für Angewandte Mathematik, ETH Zürich



Prerequisites.

- Diagonalization of matrices
- Multi-dimensional differential calculus
- Curvature of functions

Dependencies. [Lecture → Section 11.3.2], [Lecture → Section 11.3.4], [Lecture → Section 11.6.1]

Duration: 65 minutes



Video and accompanying tablet notes may not match completely!

[Corrections and updates may have been made in tablet notes.]

XI. Numerical Methods for Conservation Laws

11.6. A Glimpse of Systems of Conservation Laws

11.6.3. Finite Volume Methods for Systems of Conservation Laws in 1D

1D C.P. :

$$\text{PDE (11.6.1.2):} \quad \frac{\partial \mathbf{u}}{\partial t}(x, t) + \frac{\partial \mathbf{F}(\mathbf{u})}{\partial x}(x, t) = 0 \quad \text{on } \mathbb{R} \times]0, T[, \quad (11.6.2.4a)$$

$$\text{Initial conditions:} \quad \mathbf{u}(x, 0) = \mathbf{u}_0(x) \quad \text{for } x \in \mathbb{R}, \quad (11.6.2.4b)$$

11.6.3.1 FV Spatial Semi-Discretization

$\mathcal{M} \triangleq$ equidistant spatial mesh

$$\mathcal{M} := \{]x_{j-1}, x_j[: x_j := jh, j \in \mathbb{Z} \} . \quad (11.3.2.1)$$

mesh cells and dual cells $\triangleright$

Unknowns : Dual cell averages

$$\mu_j(t) \approx \frac{1}{h} \int_{x_{j-1/2}}^{x_{j+1/2}} \mathbf{u}(x, t) dx \in \mathbb{R}^m, \quad t \in]0, T[. \quad (11.6.3.3)$$

$$\frac{d\mu_j}{dt}(t) = -\frac{1}{h} \left(\hat{\mathbf{F}}(\mu_{j-m_l+1}(t), \dots, \mu_{j+m_r}(t)) - \hat{\mathbf{F}}(\mu_{j-m_l}(t), \dots, \mu_{j+m_r-1}(t)) \right), \quad j \in \mathbb{Z}, \quad (11.6.3.4)$$

+ truncation to bounded interval $\Rightarrow$ ODE

numerical flux $\hat{\mathbf{F}}: (\mathbb{R}^m)^{m_e+m_r} \rightarrow \mathbb{R}^m$

Most common: $m_r = m_e = 1$

11.6.3.2. Numerical fluxes for linear systems of CL

PDE: $\frac{\partial \mathbf{u}}{\partial t}(x, t) + \mathbf{A} \frac{\partial \mathbf{u}}{\partial x}(x, t) = 0$ on $\mathbb{R} \times]0, T[$, (11.6.1.15a)

+ initial conditions: $\mathbf{u}(x, 0) = \mathbf{u}_0(x)$ for $x \in \mathbb{R}$. (11.6.1.15b)

FVM with 2-point numerical flux

$$\frac{d\mu_j}{dt}(t) = -\frac{1}{h} \left(\hat{\mathbf{F}}(\mu_j(t), \mu_{j+1}(t)) - \hat{\mathbf{F}}(\mu_{j-1}(t), \mu_j(t)) \right) \quad (11.6.3.13)$$

FVM by diagonalization

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{A} \frac{\partial \mathbf{u}}{\partial x} = 0 \quad \xrightarrow{\tilde{\mathbf{u}} := \mathbf{R}^{-1} \mathbf{u}} \quad \underbrace{\frac{\partial \tilde{\mathbf{u}}}{\partial t} + \mathbf{D} \frac{\partial \tilde{\mathbf{u}}}{\partial x}} = 0, \quad \mathbf{D} := \text{diag}(\lambda_1, \dots, \lambda_m). \quad (11.6.3.9)$$

decoupled: $m \times$ scalar FVM applied *componentwise*
 $\hookrightarrow$ numerical flux $\tilde{\mathbf{F}}$

$$\frac{d\tilde{\mu}_j}{dt}(t) = -\frac{1}{h} \left(\tilde{\mathbf{F}}(\tilde{\mu}_j(t), \tilde{\mu}_{j+1}(t)) - \tilde{\mathbf{F}}(\tilde{\mu}_{j-1}(t), \tilde{\mu}_j(t)) \right). \quad (11.6.3.11)$$

$\hookrightarrow$ componentwise $\tilde{\mathbf{F}}$ for $\frac{\partial u}{\partial t} + \lambda_i \frac{\partial u}{\partial x} = 0$

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Trf. back :

$$\hat{\mathbf{F}}(\mathbf{v}, \mathbf{w}) := \mathbf{R} \tilde{\mathbf{F}}(\mathbf{R}^{-1} \mathbf{v}, \mathbf{R}^{-1} \mathbf{w}), \quad \mathbf{v}, \mathbf{w} \in \mathbb{R}^m, \quad (11.6.3.12)$$

• Upwind FVM ($\rightarrow$ Sect. 11.3.4.3)

$$\text{For } \frac{\partial u}{\partial t} + \lambda_i \frac{\partial u}{\partial x} = 0$$

$$F_{uw}(v, w) = \begin{cases} \lambda_i v & , \text{ if } \lambda_i \geq 0, \\ \lambda_i w & , \text{ if } \lambda_i < 0, \end{cases} = \max\{\lambda_i, 0\}v + \min\{\lambda_i, 0\}w, \quad v, w \in \mathbb{R}. \quad (11.6.3.15)$$

$$\tilde{\mathbf{F}}_{uw}(\mathbf{v}, \mathbf{w}) = \mathbf{D}^+ \mathbf{v} + \mathbf{D}^- \mathbf{w}, \quad \begin{aligned} \mathbf{D}^+ &:= \text{diag}(\max\{\lambda_i, 0\}, i = 1, \dots, m), \\ \mathbf{D}^- &:= \text{diag}(\min\{\lambda_i, 0\}, i = 1, \dots, m). \end{aligned} \quad (11.6.3.16)$$

w/ (3.12) :

$$\hat{\mathbf{F}}_{uw}(\mathbf{v}, \mathbf{w}) = \mathbf{A}^+ \mathbf{v} + \mathbf{A}^- \mathbf{w}, \quad \mathbf{v}, \mathbf{w} \in \mathbb{R}^m, \quad \mathbf{A}^\pm := \mathbf{R} \mathbf{D}^\pm \mathbf{R}^{-1}. \quad (11.6.3.17)$$

• Lax-Friedrich / Rusanov FVM ($\rightarrow$ Sect 11.3.4.2)

$$\frac{\partial \tilde{u}_i}{\partial t} + \lambda_i \frac{\partial \tilde{u}_i}{\partial x} = 0 \quad \blacktriangleright \quad F_{LF}(v, w) = \frac{1}{2} \lambda_i (v + w) - \frac{1}{2} |\lambda_i| (w - v), \quad v, w \in \mathbb{R}. \quad (11.6.3.19)$$

$$|\lambda| = \max(0, \lambda) - \min(0, \lambda)$$

$$\tilde{\mathbf{F}}_{LF}(\mathbf{v}, \mathbf{w}) = \frac{1}{2} \mathbf{D}(\mathbf{v} + \mathbf{w}) - \frac{1}{2} |\mathbf{D}|(\mathbf{w} - \mathbf{v}), \quad \mathbf{v}, \mathbf{w} \in \mathbb{R}^m, \quad |\mathbf{D}| := \mathbf{D}^+ - \mathbf{D}^-. \quad (11.6.3.20)$$

w/ (3.12)

$$\hat{\mathbf{F}}_{LF}(\mathbf{v}, \mathbf{w}) = \frac{1}{2} \mathbf{A}(\mathbf{v} + \mathbf{w}) - \frac{1}{2} |\mathbf{A}|(\mathbf{w} - \mathbf{v}), \quad \mathbf{v}, \mathbf{w} \in \mathbb{R}^m, \quad |\mathbf{A}| := \mathbf{A}^+ - \mathbf{A}^-. \quad (11.6.3.21)$$

• Godunov numerical flux

$$\hat{F}_{GD}(\underline{v}, \underline{w}) = \underline{F}(\underline{u}(0, t))$$

↑
similarly solution of R.P.

$$\frac{\partial \underline{u}}{\partial t}(x, t) + \underline{A} \frac{\partial \underline{u}}{\partial x}(x, t) = 0 \quad \text{on } \mathbb{R} \times \mathbb{R}^+, \quad \underline{u}(x, 0) = \begin{cases} \underline{v} & \text{for } x < 0, \\ \underline{w} & \text{for } x \geq 0. \end{cases} \quad (11.6.1.29)$$

(1.31) : shock fan solution :

$$\bar{\underline{u}}(x, t) = \begin{cases} \underline{v} & \text{for } x < \lambda_1 t, \\ \underline{w}_1^r \underline{r}_1 + \sum_{i=2}^m \underline{w}_1^l \underline{r}_i & \text{for } \lambda_1 t < x < \lambda_2 t, \\ \vdots & \vdots \\ \sum_{i=1}^k \underline{w}_i^r \underline{r}_i + \sum_{i=k+1}^m \underline{w}_1^l \underline{r}_i & \text{for } \lambda_k t < x < \lambda_{k+1} t, \\ & k = 1, \dots, m-1, \\ \vdots & \vdots \\ \underline{w} & \text{for } x > \lambda_m t, \end{cases} \quad (11.6.1.31)$$

$$\underline{u}_l = \sum_{k=1}^m \underline{w}_k^l \underline{r}_k, \quad \underline{w}_k^l := (\underline{R}^{-1} \underline{v})_k, \quad \underline{u}_r = \sum_{k=1}^m \underline{w}_k^r \underline{r}_k, \quad \underline{w}_k^r := (\underline{R}^{-1} \underline{w})_k. \quad (11.6.1.30)$$

$$\bar{\underline{u}}^\downarrow(\underline{v}, \underline{w}) := \bar{\underline{u}}(0, t) = \sum_{i=1}^p \underline{w}_i^r \underline{r}_i + \sum_{i=p+1}^m \underline{w}_i^l \underline{r}_i, \quad p := \max\{i \in \{1, \dots, m\} : \lambda_i < 0\}. \quad (11.6.3.23)$$

$$\begin{aligned} \underline{A} \bar{\underline{u}}^\downarrow(\underline{v}, \underline{w}) &= \sum_{i=1}^p \underline{w}_i^r \lambda_i \underline{r}_i + \sum_{i=p+1}^m \underline{w}_i^l \lambda_i \underline{r}_i = \sum_{i=1}^p (\underline{R}^{-1} \underline{w})_i \lambda_i \underline{r}_i + \sum_{i=p+1}^m (\underline{R}^{-1} \underline{v})_i \lambda_i \underline{r}_i \\ &= \underline{R} \underline{D}^- \underline{R}^{-1} \underline{w} + \underline{R} \underline{D}^+ \underline{R}^{-1} \underline{v} = \underline{A}^+ \underline{v} + \underline{A}^- \underline{w}, \end{aligned}$$

$$\triangleright \quad \overset{1}{F}_{GD} = \overset{1}{F}_{\underline{v}\underline{w}}$$

11.6.3.3. Numerical Fluxes for NLSCLs

Recall: Godunov numerical flux

$$\hat{F}_{GD}(\mathbf{v}, \mathbf{w}) := \mathbf{F}(\mathbf{u}^\downarrow(\mathbf{v}, \mathbf{w})) , \quad \mathbf{v}, \mathbf{w} \in U \subset \mathbb{R}^m, \quad \mathbf{u}^\downarrow(\mathbf{v}, \mathbf{w}) := \Psi(0), \quad (11.6.3.26)$$

$$\frac{\partial \bar{\mathbf{u}}}{\partial t} + \frac{\partial \mathbf{F}(\bar{\mathbf{u}})}{\partial x} = 0, \quad \bar{\mathbf{u}}(x, t) = \begin{cases} \mathbf{v} & \text{for } x < 0, \\ \mathbf{w} & \text{for } x \geq 0. \end{cases} \quad (11.6.3.27)$$

$\bar{\mathbf{u}}(x, t) = \Psi(x/t)$

Rem.: Exact Godunov FVM not feasible / not affordable in general

▷ Use *approximate* Godunov N.F.

Replace (3.27) w/ a "linear R.P."

$$\frac{\partial \tilde{\mathbf{u}}}{\partial t}(x, t) + \mathbf{A}_R \frac{\partial \tilde{\mathbf{u}}}{\partial x}(x, t) = 0 \quad \text{on } \mathbb{R} \times \mathbb{R}^+, \quad \tilde{\mathbf{u}}(x, t) = \begin{cases} \mathbf{v} & \text{for } x < 0, \\ \mathbf{w} & \text{for } x \geq 0, \end{cases} \quad (11.6.3.29)$$

$\hookrightarrow \mathbf{A}_R = \mathbf{A}_R(\underline{v}, \underline{w})$

"Correct shock speed preservation"

If $\underline{w} \in \mathcal{R}(\underline{v}) \Rightarrow$ single shock solution, moving with velocity $\dot{s}$

$$(\text{RHJC}) \Rightarrow \dot{s}(\underline{w} - \underline{v}) = \mathbf{F}(\underline{w}) - \mathbf{F}(\underline{v})$$

We demand: Same single-shock solution for (3.29)

$$\mathbf{A}_R(\underline{w} - \underline{v}) = \dot{s}(\underline{w} - \underline{v}) = \mathbf{F}(\underline{w}) - \mathbf{F}(\underline{v})$$

Definition 11.6.3.30. Roe matrix

A matrix-value function $\mathbf{A}_R : U \times U \rightarrow \mathbb{R}^{m,m}$ is a **Roe matrix** for the flux function $\mathbf{F} : U \rightarrow \mathbb{R}^m$, if

- (i) $\mathbf{A}_R(\mathbf{u}, \mathbf{u}) = D\mathbf{F}(\mathbf{u})$ for all $\mathbf{u} \in U$ (consistency),
- (ii) $\mathbf{A}_R(\mathbf{v}, \mathbf{w})(\mathbf{w} - \mathbf{v}) = \mathbf{F}(\mathbf{w}) - \mathbf{F}(\mathbf{v})$ for all $\mathbf{v}, \mathbf{w} \in U$ (correct shock speed),
- (iii) $\mathbf{A}_R(\mathbf{v}, \mathbf{w})$ has m distinct real eigenvalues (strict hyperbolicity).

Finding A_R is an art!

Ex. 11.6.3.31 (Roe matrix SWE)

$$\text{conservation of mass} \quad \blacktriangleright \quad \frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(vh) = 0, \quad (11.6.2.14a)$$

$$\text{conservation of momentum} \quad \blacktriangleright \quad \frac{\partial}{\partial t}(hv) + \frac{\partial}{\partial x}(hv^2 + \tfrac{1}{2}gh^2) = 0. \quad (11.6.2.14b)$$

$$\mathbf{u} = \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} := \begin{bmatrix} h \\ hv \end{bmatrix}, \quad \mathbf{F}(\mathbf{u}) := \begin{bmatrix} u_2 \\ u_2^2/u_1 + \tfrac{1}{2}gu_1^2 \end{bmatrix} \quad \blacktriangleright \quad D\mathbf{F}(\mathbf{u}) = \begin{bmatrix} 0 & 1 \\ -(u_2/u_1)^2 + gu_1 & 2u_2/u_1 \end{bmatrix}. \quad (11.6.2.15)$$

$$\mathbf{A}_R(\mathbf{v}, \mathbf{w}) = D\mathbf{F}\left(\begin{bmatrix} \bar{h} \\ \bar{h}\hat{v} \end{bmatrix}\right), \quad \begin{aligned} \bar{h} &:= \tfrac{1}{2}(h_v + h_w), \\ \hat{v} &:= \frac{\sqrt{h_v}v_v + \sqrt{h_w}v_w}{\sqrt{h_v} + \sqrt{h_w}}, \end{aligned} \quad \mathbf{v} = \begin{bmatrix} h_v \\ h_v v_v \end{bmatrix}, \quad \mathbf{w} = \begin{bmatrix} h_w \\ h_w v_w \end{bmatrix}. \quad (11.6.3.32)$$

↑ Roe average

Problem : $m=1$, Ex. 11.3.4.22

→ Non-physical shock solution of (3.29)

§ 11.6.3.33 (HLL-E "entropy fix")

Idea: Insert an auxiliary "rarefaction" state.

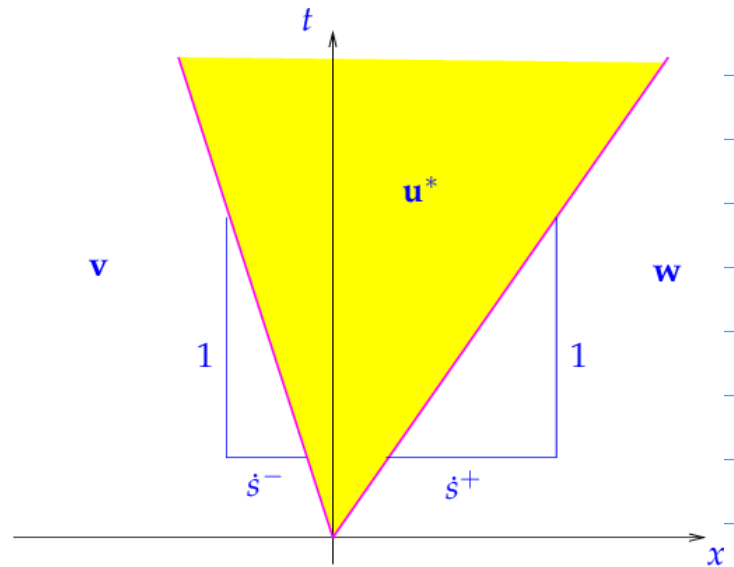
shock fan w/ 3 waves

$$\tilde{\mathbf{u}}(x, t) = \begin{cases} \mathbf{v} & , \text{if } x < \dot{s}^- t, \\ \mathbf{u}^* & , \text{if } \dot{s}^- t \leq x < \dot{s}^+ t, \\ \mathbf{w} & , \text{if } \dot{s}^+ t \leq x, \end{cases} \quad (11.6.3.34)$$

with edge speeds

$$\begin{aligned} \dot{s}^- &:= \min\{\hat{\lambda}_1, \lambda_1(\mathbf{v})\}, \\ \dot{s}^+ &:= \max\{\hat{\lambda}_m, \lambda_m(\mathbf{w})\}, \end{aligned} \quad (11.6.3.35)$$

where the $\hat{\lambda}_i, i = 1, \dots, m$, are the (sorted) eigenvalues of a **Roe matrix** $\mathbf{A}_R(\mathbf{v}, \mathbf{w})$.



Given $\dot{s}^\pm$ the "bridge state" $\mathbf{u}^*$ is obtained from global conservation

$$-\frac{d}{dt} \int_{\mathbb{R}} \tilde{\mathbf{u}}(x, t) dx = \mathbf{F}(\mathbf{w}) - \mathbf{F}(\mathbf{v}). \quad (11.6.3.36)$$

$$\dot{s}^-(\mathbf{u}^* - \mathbf{v}) + \dot{s}^+(\mathbf{w} - \mathbf{u}^*) = \mathbf{F}(\mathbf{w}) - \mathbf{F}(\mathbf{v}) \Rightarrow \mathbf{u}^* = \frac{\mathbf{F}(\mathbf{w}) - \mathbf{F}(\mathbf{v}) - \dot{s}^+ \mathbf{w} + \dot{s}^- \mathbf{v}}{\dot{s}^- - \dot{s}^+}, \quad (11.6.3.37)$$

$$\hat{\mathbf{F}}_{\text{HLL}}(\mathbf{v}, \mathbf{w}) = \begin{cases} \mathbf{F}(\mathbf{v}) & , \text{if } \dot{s}^- > 0 \\ \mathbf{F}(\mathbf{u}^*) & , \text{if } \dot{s}^- < 0 < \dot{s}^+ \\ \mathbf{F}(\mathbf{w}) & , \text{if } \dot{s}^+ < 0 \end{cases} \quad \begin{aligned} & , \dot{s}^- \text{ from (11.6.3.35)}, \\ & , \mathbf{u}^* \text{ as in (11.6.3.37)}, \\ & , \dot{s}^+ \text{ according to (11.6.3.35)}, \end{aligned} \quad (11.6.3.38)$$

$\mathbf{v}, \mathbf{w} \in \mathcal{U}$

Ex. 11.6.3.39 (HLL numerical flux, $m=1$)

$$\frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} = 0, \quad f: \mathbb{R} \rightarrow \mathbb{R} \quad C^1 \text{ \& strictly convex}$$

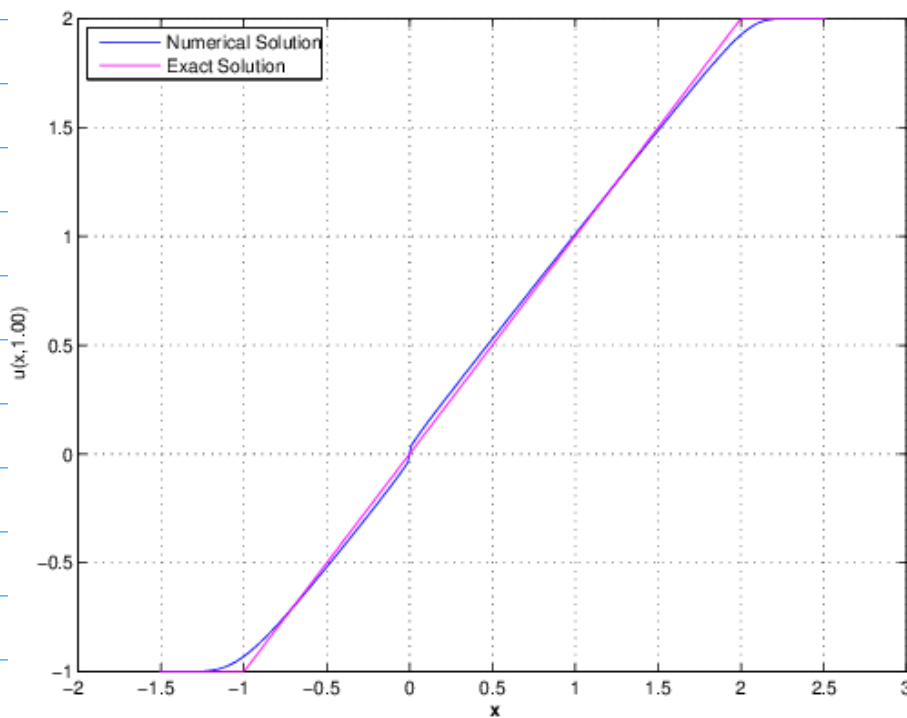
$$\triangleright A_R(v, w) = [\dot{s}], \quad \dot{s} = \begin{cases} f'(v), & v = w \\ \frac{f(w) - f(v)}{w - v}, & v \neq w \end{cases}$$

$$\dot{s}^- = \begin{cases} \dot{s}, & v > w \\ f'(v), & v \leq w \end{cases}, \quad \dot{s}^+ = \begin{cases} \dot{s}, & v > w \\ f'(w), & v \leq w \end{cases}$$

$$F_{\text{HLL}}(v, w) = f(\tilde{u}^\downarrow), \quad \tilde{u}^\downarrow := \begin{cases} v, & \text{if } \dot{s}^- > 0, \\ \frac{f(w) - f(v) - f'(w)w + f'(v)v}{f'(v) - f'(w)}, & \text{if } f'(v) < 0 < f'(w), \\ w, & \text{if } \dot{s}^+ < 0. \end{cases} \quad (11.6.3.42)$$

Case of transonic rarefaction

Exp. 11.6.3.43 (HLL-E-FVM for Burgers eqn.)



△ Transonic rarefaction resolved well

Exp. 11.6.3.46 (HLL-E-FVM for SWE)

Dam break problem, explicit Euler t.s.

§ 11.6.3.46 (Lax-Friedrich / Rusanov N.F.)

Linear case $F(u) = A \cdot u$

$$\hat{F}_{LF}(v, w) = \frac{1}{2}A(v + w) - \frac{1}{2}|A|(w - v), \quad v, w \in \mathbb{R}^m, \quad |A| := A^+ - A^-. \quad (11.6.3.21)$$

central flux

diffusive flux

⇓

$$\rightarrow \frac{1}{2}(F(v) + F(w))$$

Idea: Replace $|A| \rightarrow |A_R|$

$$\hat{F}_{LF}(v, w) = \frac{1}{2}(F(v) + F(w)) - \frac{1}{2}|A_R(v, w)|(w - v), \quad v, w \in U, \quad (11.6.3.48)$$

$$|A_R(v, w)| := R \operatorname{diag}(|\hat{\lambda}_1|, \dots, |\hat{\lambda}_m|) R^{-1}, \quad A_R(v, w) = R \operatorname{diag}(\hat{\lambda}_1, \dots, \hat{\lambda}_m) R^{-1}.$$

Exp. 11.6.3.49 (L.F. - FVM for SWE)

→ dam break problem

Rem. 11.6.3.6: Timestepping for FVM semi-discretization as in Sect 11.4.

▷ Review questions 11.6.3-50





