

ETH Lecture 401-0674-00L Numerical Methods for Partial Differential Equations

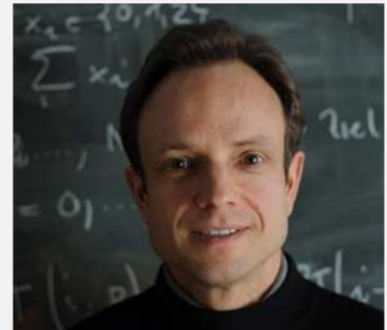
Course Video

Section 12.3.1: Pressure Instability

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(C) Seminar für Angewandte Mathematik, ETH Zürich



Dependencies. [Lecture $\rightarrow$ § 12.3.0.1], [Lecture $\rightarrow$ Section 2.6], [Lecture $\rightarrow$ Section 12.2.2]

Duration: 40 minutes



Video and accompanying tablet notes may not match completely!

[Corrections and updates may have been made in tablet notes.]

XII. Finite Elements for the Stokes Equations

12.3. Galerkin Discretization of the Stokes Saddle Point Problem

$$\underline{v} \in (H_0^1(\Omega))^d, p \in L_*^2(\Omega) \text{ s.t.}$$

$$\begin{aligned} \int_{\Omega} \mu D \underline{v} : D \underline{w} \, dx + \int_{\Omega} \operatorname{div} \underline{w} p \, dx &= \int_{\Omega} \underline{f} \cdot \underline{w} \, dx \quad \forall \underline{w} \in (H_0^1(\Omega))^d, \\ \int_{\Omega} \operatorname{div} \underline{v} q \, dx &= 0 \quad \forall q \in L_*^2(\Omega). \end{aligned} \quad (12.2.2.19)$$



$$\begin{aligned} \underline{v} \in U := (H_0^1(\Omega))^d \\ p \in Q := L_*^2(\Omega) \end{aligned} : \begin{aligned} a(\underline{v}, \underline{w}) + b(\underline{w}, p) &= \ell(\underline{w}) \quad \forall \underline{w} \in U, \\ b(\underline{v}, q) &= 0 \quad \forall q \in Q. \end{aligned} \quad (12.3.0.2)$$

§ 12.3.0.1 Galerkin discretization of saddle point problems

$$\text{Step ①} \quad \begin{array}{ccc} \mathcal{U} & \longrightarrow & \mathcal{U}_h \subset \mathcal{U} \\ \mathcal{Q} & \longrightarrow & \mathcal{Q}_h \subset \mathcal{Q} \end{array} \quad \text{in (2.19)}$$

$$N := \dim \mathcal{U}_h, \quad M := \dim \mathcal{Q}_h < \infty$$

DVP:

$$\begin{array}{l} \mathbf{v}_h \in U_h \\ p_h \in Q_h \end{array} : \quad \begin{array}{l} \mathbf{a}(\mathbf{v}_h, \mathbf{w}_h) + \mathbf{b}(\mathbf{w}_h, p_h) = \ell(\mathbf{w}_h) \quad \forall \mathbf{w}_h \in U_h, \\ \mathbf{b}(\mathbf{v}_h, q_h) = 0 \quad \forall q_h \in Q_h. \end{array} \quad (12.3.0.4)$$

Step ② : Choose bases $\{\mathbf{b}_h^i\}_{i=1}^N$ of U_h
 $\{\beta_h^j\}_{j=1}^M$ of Q_h

+ expansion : $\underline{v}_h \Rightarrow \sum_{i=1}^N v_i \mathbf{b}_h^i$
 $p_h \approx \sum_{j=1}^M \pi_j \beta_h^j$

▷ LSE

saddle point matrix
 [symmetric, not s.p.d.]

saddle point LSE:

$$\begin{bmatrix} \mathbf{A} & \mathbf{B}^T \\ \mathbf{B} & 0 \end{bmatrix} \begin{bmatrix} \vec{v} \\ \vec{\pi} \end{bmatrix} = \begin{bmatrix} \vec{\phi} \\ 0 \end{bmatrix}, \quad (12.3.0.7)$$

$$\mathbf{A} := \left[\mathbf{a}(\mathbf{b}_h^j, \mathbf{b}_h^i) \right]_{i,j=1}^N = \left[\int_{\Omega} \mu \mathbf{D} \mathbf{b}_h^j(x) : \mathbf{D} \mathbf{b}_h^i(x) dx \right]_{i,j=1}^N \in \mathbb{R}^{N,N}, \quad (12.3.0.8a)$$

$$\mathbf{B} := \left[\mathbf{b}(\mathbf{b}_h^j, \beta_h^i) \right]_{\substack{1 \leq i \leq M \\ 1 \leq j \leq N}} = \left[\int_{\Omega} \operatorname{div} \mathbf{b}_h^j(x) \beta_h^i(x) dx \right]_{\substack{1 \leq i \leq M \\ 1 \leq j \leq N}} \in \mathbb{R}^{M,N}, \quad (12.3.0.8b)$$

$$\vec{\phi} := \left[\ell(\mathbf{b}_h^j) \right]_{j=1}^N = \left[\int_{\Omega} \mathbf{f}(x) \cdot \mathbf{b}_h^j(x) dx \right]_{j=1}^N \in \mathbb{R}^N. \quad (12.3.0.8c)$$

12.3.1. Pressure Instability

Simple setting of Sect. 3.1

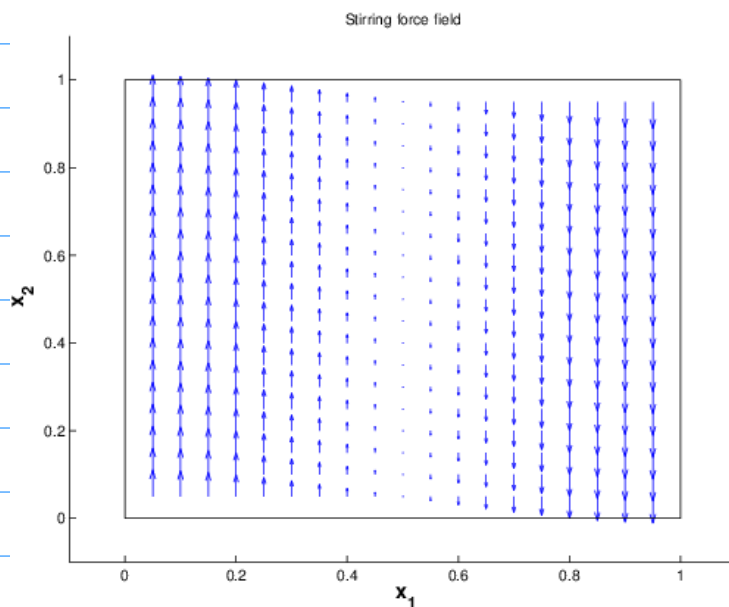
$$\text{V.P. : } u \in V_0 : a(u, v) = l(v) \quad \forall v \in V_0$$

$\hookrightarrow$ *s.p.d.* $\Rightarrow$ induces $\|\cdot\|_a$

Gal. disc: $V_0 \rightarrow V_{0,h} \subset V_0$

Cea's lemma: $\|u - u_h\|_a \leq \inf_{v_h \in V_{0,h}} \|u - v_h\|_a$

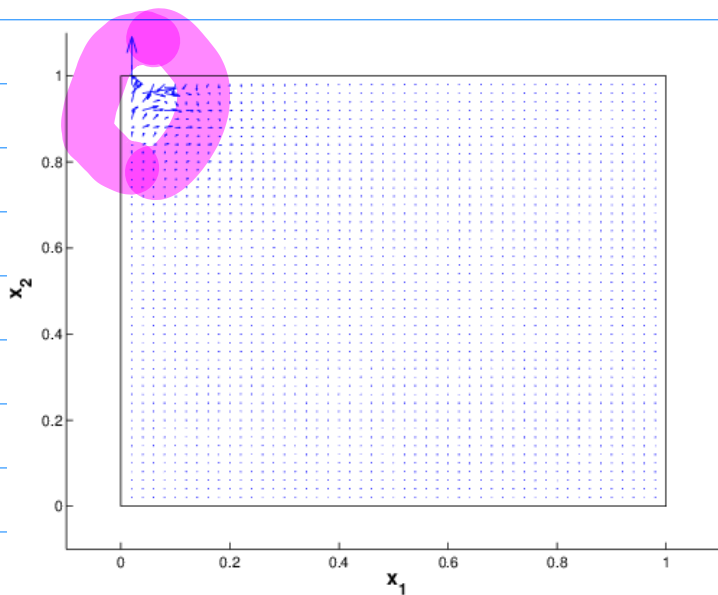
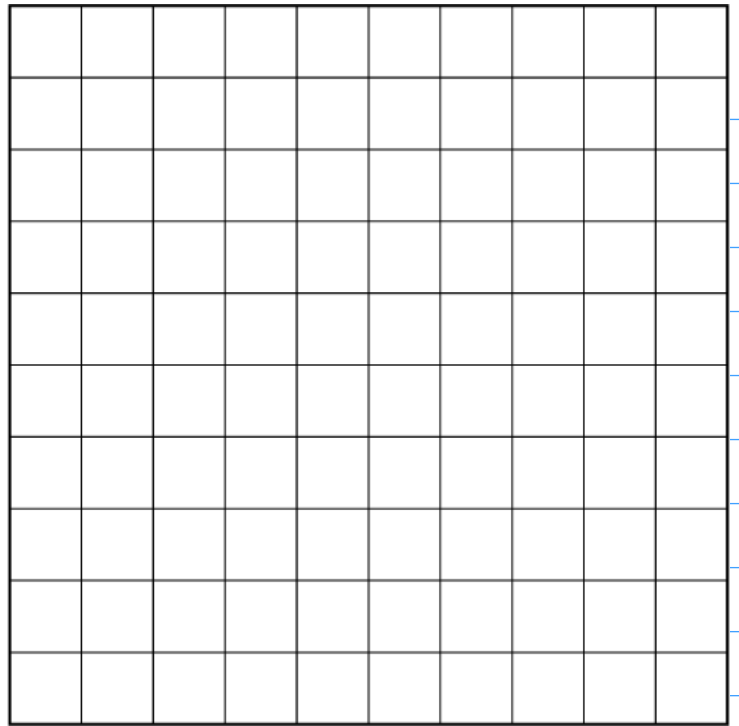
Exp. 12.3.1.1 : Q1-Q1 FEM on tensor-product mesh



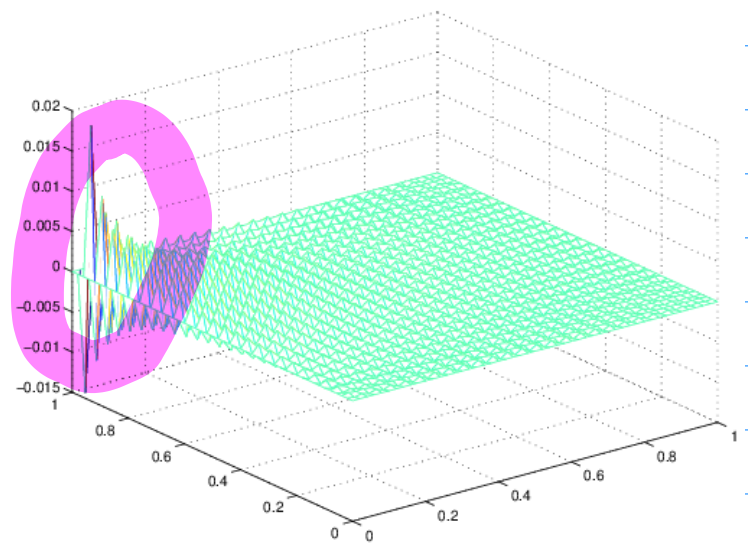
- ◆ on the unit square $\Omega =]0, 1[^2$,
 - ◆ with $\mu \equiv 1$, and
 - ◆ with driving force field $\mathbf{f} = \cos(\pi x_1) \begin{bmatrix} 0 \\ 1 \end{bmatrix}$,
- ◁ Right-hand side force field $\mathbf{f}$

- ◆ $n \times n$ tensor-product quadrilateral mesh $\mathcal{M}$ ▷
- ◆ with $U_h := (\mathcal{S}_{1,0}^0(\mathcal{M}))^2 \subset (H_0^1(\Omega))^2$,
- ◆ and $Q_h := \mathcal{S}_1^0(\mathcal{M}) \subset L^2(\Omega)$,

p.w. - bilinear



velocity solution $\mathbf{v}_h$



pressure solution p_h

→ spurious oscillations

Back $(M+N) \times (M+N)$

saddle point LSE:
$$\begin{matrix} N \updownarrow \\ \begin{bmatrix} \mathbf{A} & \mathbf{B}^T \\ \mathbf{B} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \vec{v} \\ \vec{\pi} \end{bmatrix} = \begin{bmatrix} \vec{\varphi} \\ \mathbf{0} \end{bmatrix} \end{matrix},$$

$$\begin{matrix} \leftarrow M \end{matrix}$$

(12.3.0.7)

$\Rightarrow$ If $\mathcal{N}(\mathbf{B}^T) \neq \{0\} \Rightarrow \vec{\pi}$ not unique

$$M > N \Rightarrow N(B^T) = \{0\}$$

$\dim U_h \geq \dim Q_h$ is a *necessary* condition for the uniqueness of the pressure solution p_h of the discrete saddle point variational problem (12.3.0.4).

FEM with discontinuous pressure:

$$\mathcal{S}_p^{-1}(\mathcal{M}) := \{q_h \in L^2(\Omega) : q_h|_K \in \mathcal{P}_p(\mathbb{R}^d)\}, \quad \mathcal{M} \text{ a mesh of } \Omega \subset \mathbb{R}^d. \quad (12.3.1.2)$$

$$\text{Need } Q_h \subset L_*^2(\Omega) \longrightarrow Q_h := \mathcal{S}_{p,*}^{-1}(\mathcal{M})$$

Ex. 12.3.1.14 (P1-P0 FEM)

"Tensor-product" triangular mesh of $]0,1[^2$ $\triangleright$

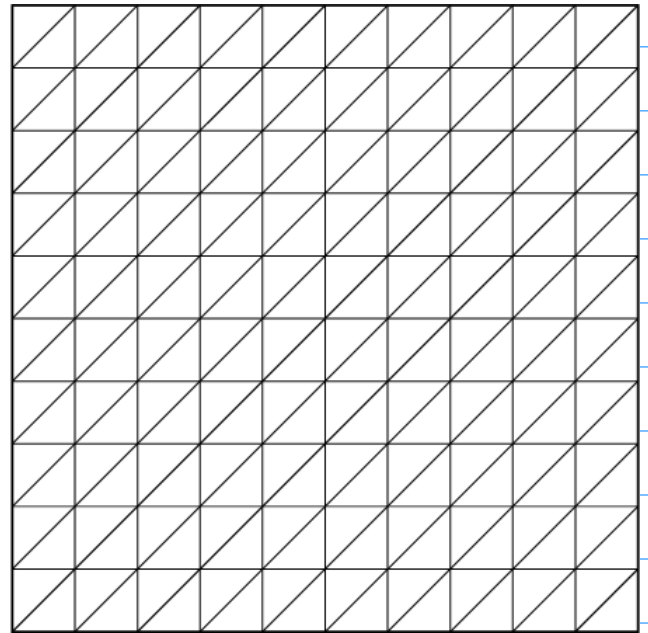
Finite element spaces for (12.2.2.19),

$$U_h := (\mathcal{S}_{1,0}^0(\mathcal{M}))^2,$$

$$Q_h := \mathcal{S}_{0,*}^{-1}(\mathcal{M}) \quad (\mathcal{M}\text{-piecewise constants}).$$

$$\dim U_h = 2(n-1)^2, \quad \dim Q_h = 2n^2 - 1.$$

$$\begin{array}{ccc} \uparrow & & \nwarrow \\ 2\#\{\text{interior nodes}\} & & \#\mathcal{M} - 1 \end{array}$$



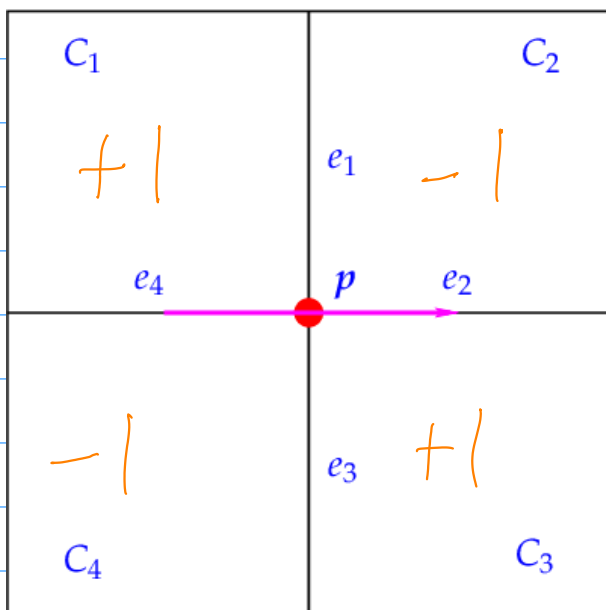
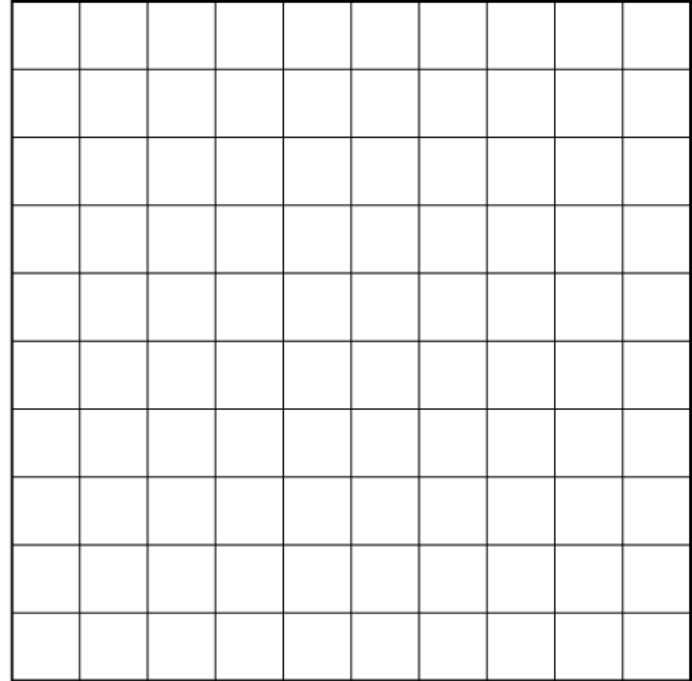
$$\dim U_h < \dim Q_h \Rightarrow \text{Singular LSE}$$

Ex. 12.3.1.5 (Q1-P0 FEM)

- ♦ $\mathcal{M}$ is a uniform tensor product mesh of the unit square domain $\Omega :=]0, 1[^2$ $\triangleright$
- ♦ velocity space $U_h = (\mathcal{S}_{1,0}^0(\mathcal{M}))^2$
- ♦ pressure space $Q_h = \mathcal{S}_0^{-1}(\mathcal{M}) \cap L_*^2(\Omega)$

$$\dim U_h = 2(n-1)^2, \quad \dim Q_h = n^2 - 1.$$

$$\blacktriangleright \dim Q_h < \dim U_h \quad \text{for } n \geq 4.$$



$\leftarrow p_h$ - cell values

$$\int_{\Omega} \operatorname{div} \underline{b}_h^P \cdot p_h \, dx = 0!$$

+1	-1	+1	-1	+1	-1	+1	-1	+1	-1
-1	+1	-1	+1	-1	+1	-1	+1	-1	+1
+1	-1	+1	-1	+1	-1	+1	-1	+1	-1
-1	+1	-1	+1	-1	+1	-1	+1	-1	+1
+1	-1	+1	-1	+1	-1	+1	-1	+1	-1
-1	+1	-1	+1	-1	+1	-1	+1	-1	+1
+1	-1	+1	-1	+1	-1	+1	-1	+1	-1
-1	+1	-1	+1	-1	+1	-1	+1	-1	+1
+1	-1	+1	-1	+1	-1	+1	-1	+1	-1
-1	+1	-1	+1	-1	+1	-1	+1	-1	+1

Checker board mode $\triangleright$

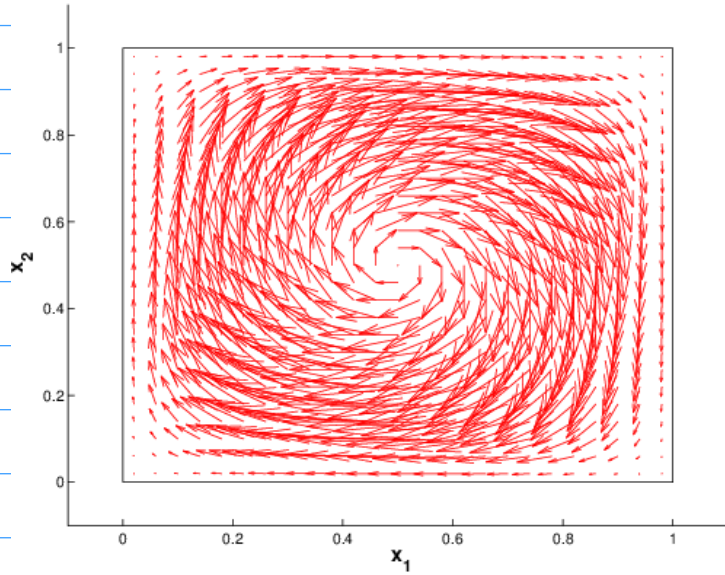
$$\triangleright \in \mathcal{N}(B^T)$$

$\rightarrow$ spurious oscillations!

Exp. 12.3.17 (P1-P0 quadrilateral FEM)

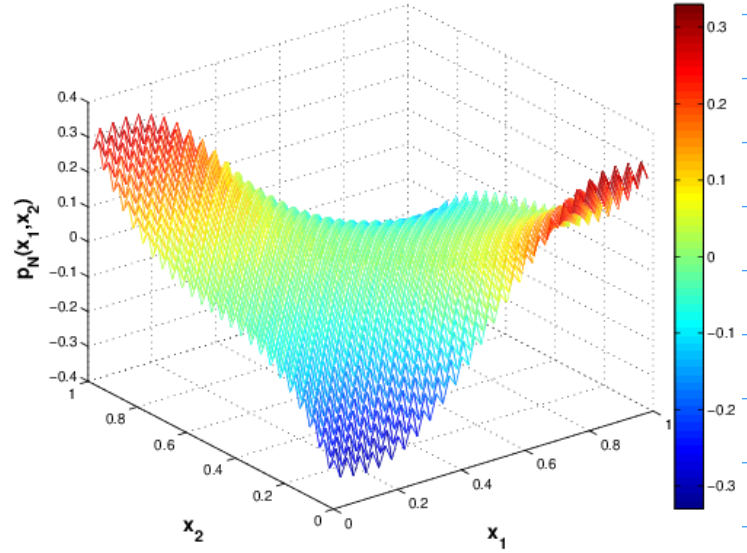
$h = 50$:

P1-P0 velocity vector field, N=50



ok !

P1-P0 pressure field, N=50



bad !







