

ETH Lecture 401-0674-00L Numerical Methods for Partial Differential Equations

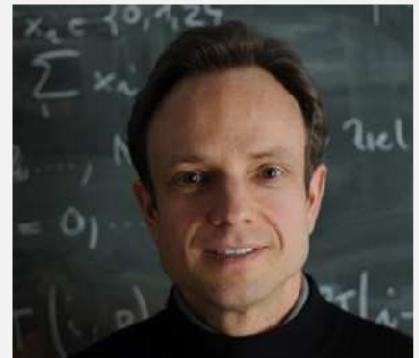
# Course Video

## Section 13.3.3.2: Magnetostatics: Vector-Potential-Based Variational Formulations

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Date: May 9, 2025

(C) Seminar für Angewandte Mathematik, ETH Zürich



**Dependencies.** [Lecture → Section 1.8], [Lecture → Section 13.3.1.2], [Lecture → Section 13.1.4], [Lecture → Section 13.1.5].

Duration: 40 minutes



Video and accompanying tablet notes may not match completely!

[Corrections and updates may have been made in tablet notes.]

# XIII Finite - Element Exterior Calculus (FEEC)

## 13.3. Whitney FEM for Computational Electromagnetism

### 13.3.3. Magnetostatics (MS)

#### 13.3.3.2. MS: Vector-Potential-Based Formulation

BVP:

$$d_1 \mathbf{h} = \mathbf{j}, \quad d_2 \mathbf{b} = 0 \quad \text{in } \Omega \subset \mathbb{R}^3, \quad \mathbf{h}|_{\partial\Omega} = 0 \quad \text{on } \partial\Omega, \quad (13.3.3.9)$$

$$\int_{\Omega} \mathbf{b} \wedge \mathbf{h}' = m_{\mu}(\mathbf{h}, \mathbf{h}') \quad \forall \mathbf{h}' \in L^2 \Lambda^1(\Omega), \quad \int_{\Omega} \mathbf{h} \wedge \mathbf{b}' = m_{\mu}^*(\mathbf{b}, \mathbf{b}') \quad \forall \mathbf{b}' \in L^2 \Lambda^2(\Omega).$$

$$\mathbf{b} = d_1 \mathbf{a}, \quad \mathbf{a} \triangleq \text{magnetic vector potential}$$

$$\underbrace{d_1 \mathbf{h} = \mathbf{j}, \quad \mathbf{b} = d_1 \mathbf{a}}_{(AL)} \quad \text{in } \Omega \subset \mathbb{R}^3, \quad \mathbf{h}|_{\partial\Omega} = 0 \quad \text{on } \partial\Omega. \quad (13.3.3.19)$$

(AL)

$$\begin{aligned}
 \int_{\Omega} g \wedge a' &\stackrel{AL}{\downarrow} = \int_{\Omega} d_1 h \wedge a' \stackrel{i.b.p.}{=} \int_{\Omega} h \wedge d_1 a' + \int_{\partial\Omega} h \wedge a' \\
 &\stackrel{ML}{\downarrow} = m_p^*(b, da') \stackrel{pot.}{=} m_p^*(d_1 a, da') \quad \forall a'
 \end{aligned}$$

$\chi = 0$   
 $[h|_{\partial\Omega} = 0]$

$$a \in H\Lambda^1(d, \Omega): m_{\mu}^*(d_1 a, d_1 a') = \int_{\Omega} j \wedge a' \quad \forall a' \in H\Lambda^1(d, \Omega). \quad (13.3.3.21)$$

LLML & v.p.:

$$\begin{aligned}
 \check{a} \in H(\text{curl}, \Omega): \int_{\Omega} \mu(x)^{-1} \text{curl } \check{a}(x) \cdot \text{curl } \check{a}'(x) dx \\
 = \int_{\Omega} j(x) \cdot \check{a}'(x) dx \quad \forall \check{a}' \in H(\text{curl}, \Omega). \quad (13.3.3.22)
 \end{aligned}$$

PMC b.c. ?  $\stackrel{!}{=} \text{natural b.c.}$  ("hidden in (21)")

### § 13.3.3.23 Gauge Freedom

If  $a$  solves (21)  $\Rightarrow a + d_0 v$  solves (21)  $\forall v \in H\Lambda^0(d, \Omega)$

$$N(d_1) = d_0 H\Lambda^0(d, \Omega)$$

#### Assumption 13.3.3.7. Simple topology of $\Omega$

We assume that the assumptions of Thm. 13.1.4.81 on  $\Omega$  for  $\ell = 1, 2$  are satisfied:

- Every oriented closed 2-surface  $\Omega$  is the boundary of a sub-domain  $\subset \Omega$ .
- Every closed directed curve in  $\Omega$  is the boundary of an oriented surface  $\subset \Omega$ .

(21) : Non-uniqueness

## Existence

$$\begin{aligned}
 \text{r.h.s.}(d_0 v) &= \int_{\Omega} f \, 1 \, d_0 v \\
 &\stackrel{\text{i.p.b.}}{=} - \int_{\Omega} \underbrace{d_2 f \, 1 \, v}_{} + \int_{\partial \Omega} \underbrace{f \, 1 \, v}_{} = 0 \\
 &\quad \quad \quad = 0 \quad \quad \quad \partial \Omega \cap \Omega = \emptyset
 \end{aligned}$$

→ **Consistent** r.h.s.  $\Rightarrow$  Existence of sols

Restoring uniqueness = "suppressing kernel"

Simple setting (Ex. 13.3.3.25)

$$\begin{aligned}
 Z^T \vec{\varphi} &= 0 \\
 \Downarrow
 \end{aligned}$$

$$A \vec{\xi} = \vec{\varphi}, \quad A = A^T \in \mathbb{R}^{N,N}, \quad \vec{\varphi} \in \mathbb{R}^N, \quad \mathcal{N}(A) \neq \{0\}, \quad \vec{\varphi} \perp \mathcal{N}(A).$$

$$Z \in \mathbb{R}^{N,M} : \mathcal{R}(A) = \mathcal{N}(A), \quad M = \dim \mathcal{N}(A)$$



Restore uniqueness of solutions by enforcing orthogonality to  $\mathcal{N}(A)$ !

$$Z^T \vec{\xi} = 0$$

$$A \vec{\xi} + Z \vec{\pi} = \vec{\varphi}$$

$$Z^T \vec{\xi} = 0$$

$$\cancel{Z^T} A \vec{\xi} + \underbrace{\cancel{Z^T} Z}_{\text{invertible}} \vec{\pi} = \cancel{Z^T} \vec{\varphi} = 0 \Rightarrow \vec{\pi} = 0$$

Augmented LSE:

$$\begin{bmatrix} \mathbf{A} & \mathbf{Z} \\ \mathbf{Z}^\top & \mathbf{O} \end{bmatrix} \begin{bmatrix} \vec{\xi} \\ \vec{\pi} \end{bmatrix} = \begin{bmatrix} \vec{\phi} \\ \mathbf{0} \end{bmatrix}. \quad (13.3.3.26)$$

↑ regular square matrix

Apply this idea to

$$\mathbf{a} \in H\Lambda^1(d, \Omega): m_\mu^*(d_1 \mathbf{a}, d_1 \mathbf{a}') = \int_\Omega \mathbf{j} \wedge \mathbf{a}' \quad \forall \mathbf{a}' \in H\Lambda^1(d, \Omega), \quad (13.3.3.21)$$

Orthogonality requires inner product  $m(\cdot, \cdot)$  on  $H\Lambda^1(d, \Omega)$

Use any  $L^2$ -type i.p.!

Important: "kernel representation"

$$\mathcal{N}(d_1) = d_0 H\Lambda^0(d, \Omega)$$

Seek  $\mathbf{a} \in H\Lambda^1(d, \Omega), p \in H_*\Lambda^0(d, \Omega) := \{q \in H\Lambda^0(d, \Omega) : \int_\Omega v.p.(q)(x) dx = 0\}$

$$\begin{aligned} m_\mu^*(d_1 \mathbf{a}, d_1 \mathbf{a}') + m(\mathbf{a}', d_0 p) &= \int_\Omega \mathbf{j} \wedge \mathbf{a}' \quad \forall \mathbf{a}' \in H\Lambda^1(d, \Omega), \\ m(\mathbf{a}, d_0 p') &= 0 \quad \forall p' \in H_*\Lambda^0(d, \Omega). \end{aligned} \quad (13.3.3.28)$$

↑ enforces  $m(\cdot, \cdot)$ -orthogonality to  $\mathcal{N}(d_1)$   
 $p \neq 0!$

$$\text{Note: } \mathcal{N}(d_0) = \mathcal{P}_0 \Lambda^0(\Omega)$$

▷ Enforce uniqueness of  $p$  by orthogonality to constants!

LLML, v.p. notation:

Seek  $\check{\mathbf{a}} \in H(\mathbf{curl}, \Omega)$ ,  $\check{p} \in H_*^1(\Omega) := \{v \in H^1(\Omega) : \int_{\Omega} v(x) dx = 0\}$  such that

$$\int_{\Omega} \mu(x)^{-1} \mathbf{curl} \check{\mathbf{a}}(x) \cdot \mathbf{curl} \check{\mathbf{a}}'(x) dx + \int_{\Omega} \check{\mathbf{a}}'(x) \cdot \mathbf{grad} \check{p}(x) dx = \int_{\Omega} \check{\mathbf{j}}(x) \cdot \check{\mathbf{a}}'(x) dx,$$

$$\int_{\Omega} \check{\mathbf{a}}(x) \cdot \mathbf{grad} \check{p}'(x) dx = 0.$$

(13.3.3.29)

for all  $\check{\mathbf{a}}' \in H(\mathbf{curl}, \Omega)$ ,  $\check{p}' \in H_*^1(\Omega)$ .

### § 13.3.3.30 Galerkin FEM

▷ Discrete kernel representation

$$\{v_h \in W^1(\mathcal{M}) : d_1 v_h = 0\} = d_0 W^0(\mathcal{M})$$

Step I: f.d. subspaces  $\rightarrow$  Whitney FE spaces

Seek  $\mathbf{a}_h \in \mathcal{W}^1(\mathcal{M})$ ,  $p_h \in \mathcal{W}^0(\mathcal{M}) \cap H_* \Lambda^0(d, \Omega)$  such that

$$m_{\mu}^*(d_1 \mathbf{a}_h, d_1 \mathbf{a}'_h) + m(\mathbf{a}'_h, d_0 p_h) = \int_{\Omega} \mathbf{j} \wedge \mathbf{a}'_h \quad \forall \mathbf{a}'_h \in \mathcal{W}^1(\mathcal{M}),$$

$$m(\mathbf{a}_h, d_0 p'_h) = 0 \quad \forall p'_h \in \mathcal{W}^0(\mathcal{M}) \cap H_* \Lambda^0(d, \Omega).$$

(13.3.3.32)

$\hat{=}$  orthogonality to (discrete) kernel

Vanishing mean constraint:

Seek  $\mathbf{a}_h \in \mathcal{W}^1(\mathcal{M})$ ,  $p_h \in \mathcal{W}^0(\mathcal{M})$ , and  $\lambda \in \mathbb{R}$  such that

$$m_\mu^*(d_1 \mathbf{a}_h, d_1 \mathbf{a}'_h) + m(\mathbf{a}'_h, d_0 p_h) = \int_{\Omega} \mathbf{j} \wedge \mathbf{a}'_h,$$

$$m(\mathbf{a}_h, d_0 p'_h) + \lambda \int_{\Omega} p_h(x)' dx = 0, \quad (13.3.33)$$

$$\int_{\Omega} p_h(x) dx = 0.$$

for all  $\mathbf{a}'_h \in \mathcal{W}^1(\mathcal{M})$ ,  $p'_h \in \mathcal{W}^0(\mathcal{M})$ .

$\lambda = 0$

"Orthogonality to const."

Step II: Ordered bases {Whitney 1/O-forms}

▷ LSE:

$$\begin{bmatrix} \mathbf{A} & \mathbf{B}^\top & \mathbf{0} \\ \mathbf{B} & \mathbf{0} & \mathbf{c} \\ \mathbf{0} & \mathbf{c}^\top & \mathbf{0} \end{bmatrix} \begin{bmatrix} \vec{\mathbf{a}}_h \\ \vec{p}_h \\ \lambda \end{bmatrix} = \begin{bmatrix} \vec{\psi} \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}, \quad (13.3.34)$$

↳ dummy unknowns = 0



















