

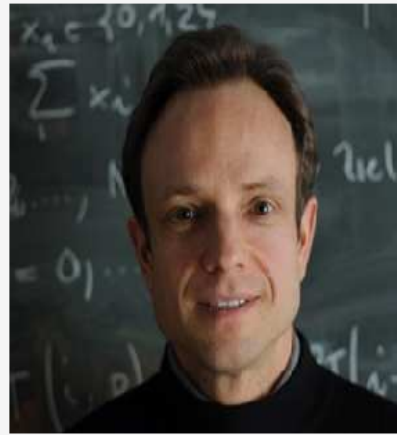
Homework Problem Video Tutorial

Problem 6-4: The One-Dimensional Wave Equation with Absorbing Boundary Condition

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1D wave propagation

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0, \quad \text{on }]0, 1[\times]0, T[. \quad (6.4.1)$$

 $\uparrow c \equiv \text{wave speed}$

$$c \frac{\partial u}{\partial x}(0, t) - \frac{\partial u}{\partial t}(0, t) = 0, \quad \leftarrow \text{absorbing b.c.} \quad (6.4.2a)$$

$$u(1, t) = g(t) := \begin{cases} \sin t & 0 \leq t \leq \pi, \\ 0 & \text{otherwise,} \end{cases} \quad \left. \vphantom{\begin{matrix} u(1, t) = g(t) \\ \sin t \\ 0 \end{matrix}} \right\} \begin{array}{l} \text{spatial} \\ \text{Dirichlet b.c.} \end{array} \quad (6.4.2b)$$

$$u(x, 0) = 0, \quad \frac{\partial u}{\partial t}(x, 0) = 0, \quad 0 < x < 1. \quad (6.4.2c)$$

initial conditions

a, Spatial variational formulation

Test space $V_0 := \{v \in H^1(\mathbb{R}) : v(1) = 0\}$ Test with $v \in V_0$ & integrate over $]0, 1[$ & i.b.p.

$$\int_0^1 \frac{\partial^2 u}{\partial t^2}(x, t) v(x) dx + c \int_0^1 \frac{\partial u}{\partial x}(x, t) \frac{\partial v}{\partial x}(x) dx + c \frac{\partial u}{\partial x}(0, t) v(0) = 0 \quad \forall v \in V_0$$

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$$u :]0, T[\rightarrow \{v \in H^1(\mathbb{R}_+^1) : v(1) = g(t)\}$$

② $\triangleright$ t -dependent affine trial space

$$m(\ddot{u}, v) + b(\dot{u}, v) + a(u, v) = 0 \quad \forall v \in V_0, \quad (6.4.3)$$

where

$$m(u, v) := \int_0^1 uv \, dx,$$

$$b(u, v) := u(0)v(0),$$

$$a(u, v) := c \int_0^1 \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} \, dx.$$

(6.4.4)

$$\mathbf{M}_{\text{full}} = h \begin{pmatrix} \frac{1}{2} & & & \\ & 1 & & \\ & & \ddots & \\ & & & 1 \\ & & & & \frac{1}{2} \end{pmatrix}, \quad \mathbf{A}_{\text{full}} = \frac{c}{h} \begin{pmatrix} 1 & -1 & & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & -1 & 2 & -1 \\ & & & -1 & 1 \end{pmatrix}.$$

b , Spatial Galerkin discretization based $S_1^0(\mathcal{M})$, $\mathcal{M} \cong$ equidistant mesh

By the method-of-lines approach as outlined above we arrive at an ordinary differential equation of the form

$$\mathbf{M}\ddot{\vec{\mu}} + \mathbf{B}\dot{\vec{\mu}} + \mathbf{A}\vec{\mu} = \vec{\varphi}(t), \quad (6.4.5)$$

with matrices $\mathbf{M}, \mathbf{A}, \mathbf{B} \in \mathbb{R}^{N,N}$ and a time-dependent vector $\vec{\varphi}$. Here N stands for the dimension of the finite element test space.

Give concrete formulas for the entries of $\mathbf{M}, \mathbf{A}, \mathbf{B} \in \mathbb{R}^{N,N}$ and the components of $\vec{\varphi}$.

$\triangleright$ All integrations based on trapezoidal rule

$\Rightarrow$ Mass $\mathbf{M}$ matrix diagonal

$$(\mathbf{B})_{ij} = \begin{cases} 1 & \text{for } i=j=1 \\ 0 & \text{else} \end{cases}$$

Galerkin matrices for full FE space $S_1^0(\mathcal{M})$

$$\mathbf{M}_{\text{full}} = h \begin{pmatrix} \frac{1}{2} & & & \\ & 1 & & \\ & & \ddots & \\ & & & 1 \\ & & & & \frac{1}{2} \end{pmatrix}, \quad \mathbf{A}_{\text{full}} = \frac{c}{h} \begin{pmatrix} 1 & -1 & & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & -1 & 2 & -1 \\ & & & -1 & 1 \end{pmatrix}.$$

Non-zero Dirichlet b.c. are taken into account by **offset function**

$$u_h(t) = \sum_{e=1}^N p_e(t) b_n^e + g(t) b_n^{N+1}$$

$N := \dim(S_1^0(\mathcal{M}) \cap V_0)$ $\uparrow$
tent function at $x=1$

We start from:

$$m(\ddot{u}_h(t), b_n^e) + b(\dot{u}_h(t), b_n^e) + a(u_h(t), b_n^e) = 0$$

$$\begin{aligned} (\vec{\varphi})_e &= -m(\ddot{g}(t) b_n^{N+1}, b_n^e) - b(\dot{g}(t) b_n^{N+1}, b_n^e) \\ &\quad - a(g(t) b_n^{N+1}, b_n^e), \quad e=1, \dots, N \\ &= \begin{cases} \frac{c}{h} g(t) & \text{for } e=N \\ 0 & \text{elsewhere} \end{cases} \end{aligned}$$

③ Leapfrog timestepping: $\vec{v}^{(j)} \approx \dot{\vec{u}}(t_{j+1/2})$ d,

$$\mathbf{M} \frac{\vec{v}^{(j+1/2)} - \vec{v}^{(j-1/2)}}{\tau} = -\mathbf{A}\vec{\mu}^{(j)} - \mathbf{B} \frac{\vec{v}^{(j+1/2)} + \vec{v}^{(j-1/2)}}{2} + \vec{\varphi}(t_j), \quad (6.4.6)$$

$$\frac{\vec{\mu}^{(j+1)} - \vec{\mu}^{(j)}}{\tau} = \vec{v}^{(j+1/2)} \approx \dot{\vec{u}}(t_j)$$

Initial step: $\vec{v}^{(-1/2)} = 2\vec{v}_0 - \vec{v}^{(1/2)}$. (6.4.7)

c, Implementation of leapfrog timestepping II

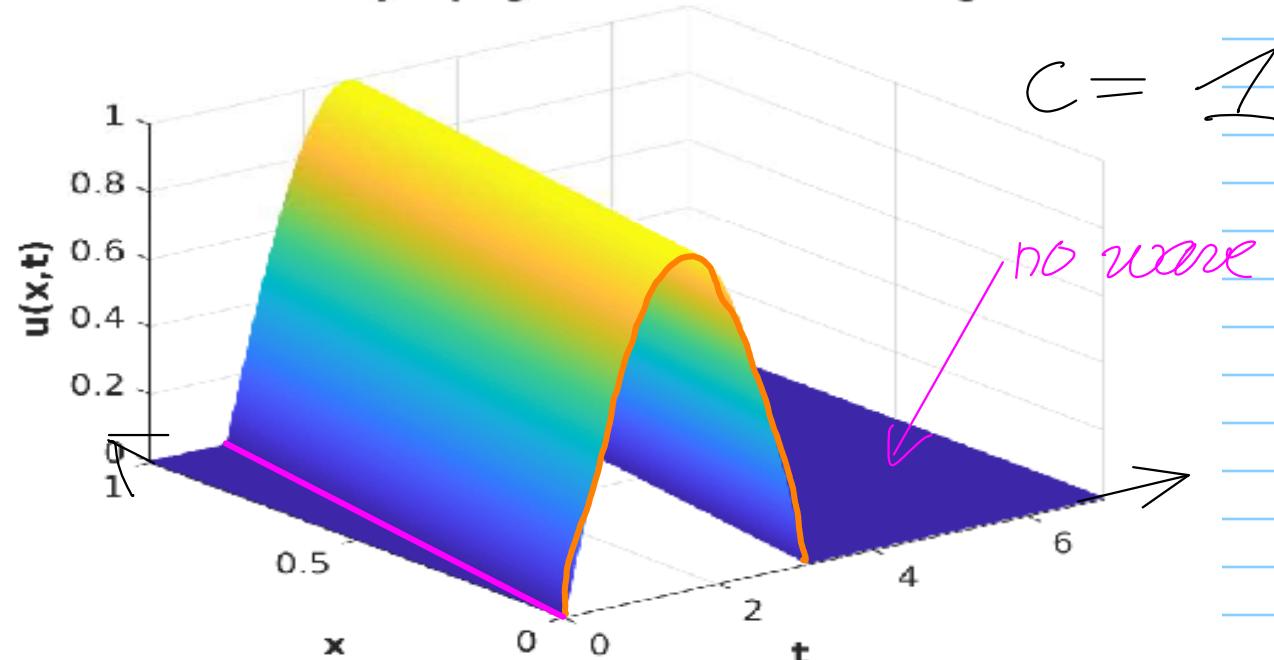
Initial step: $\vec{v}^{(1/2)} = 0$
 (6.4.6) for $j=0$: $2\mathbf{M}\vec{v}^{(1/2)} = -\tau\mathbf{A}\vec{\mu}_0 - \tau\mathbf{B}\vec{v}_0 + 2\mathbf{M}\vec{v}_0 + \tau\vec{\varphi}(0) = 0$
 $\uparrow$ initial values $= 0$

$$\vec{\mu}^{(j+1)} := \vec{\mu}^{(j)} + \tau\vec{v}^{(j+1/2)}, \quad j \in \mathbb{N}_0, \quad (6.4.10)$$

$$(\mathbf{M} + \frac{1}{2}\tau\mathbf{B})\vec{v}^{(j+1/2)} := -\tau\mathbf{A}\vec{\mu}^{(j)} - \frac{1}{2}\mathbf{B}\vec{v}^{(j-1/2)} + \mathbf{M}\vec{v}^{(j-1/2)} + \tau\vec{\varphi}(t_j), \quad j \in \mathbb{N}. \quad (6.4.11)$$

II diagonal matrix

1D wave propagation with absorbing b.c.



Excitation just propagates, cf. d'Alembert solution

e, In the file 1dwaveabsorbingbc.cc write a C++ function

```
std::pair<Eigen::VectorXd, Eigen::VectorXd> computeEnergies(
    const Eigen::MatrixXd &full_solution,
    double c, double tau);
```

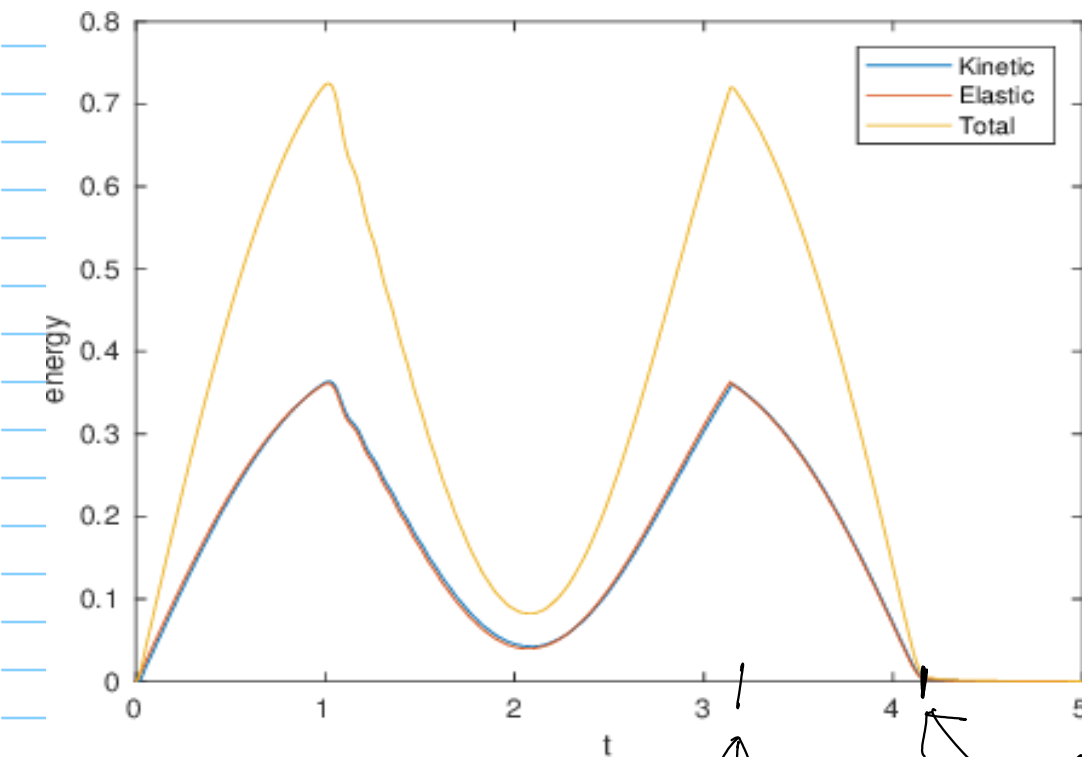
that accepts as an argument the output matrix of `waveLeapfrogABC()` from Sub-problem (6-4.c) plus the constant c and the timestep size $\tau > 0$, and computes the discrete energies (using the notations of (6.4.6))

potential energy: $E_{\text{pot}}^{(j)} := \frac{1}{2}\vec{\mu}^{(j)}\mathbf{A}\vec{\mu}^{(j)}, \quad j = 0, \dots, m,$

kinetic energy: $E_{\text{kin}}^{(j)} := \frac{1}{2}\vec{v}^{(j+1/2)}\mathbf{M}\vec{v}^{(j+1/2)}, \quad j = 0, \dots, m-1.$

These energies are returned as a pair of vectors, the potential energy in the first component, the kinetic energy in the second.

$$\vec{v}(j+1/2) := \frac{1}{\tau} (\vec{p}(j) - \vec{p}(j+1))$$



$$C = 1$$

excitation stops

All the waves have left domain at time $\pi + 1$

h, Now we consider the method-of-lines ode (6.4.5) without excitation, that is, in the case $\vec{\phi} \equiv 0$. Show that even in this case the total energy

$$E(t) = \frac{1}{2} \dot{\vec{\mu}}(t)^\top \mathbf{M} \dot{\vec{\mu}}(t) + \frac{1}{2} \vec{\mu}(t)^\top \mathbf{A} \vec{\mu}(t) \quad (6.4.15)$$

will not be conserved, but decay, that is, $t \mapsto E(t)$ is decreasing, if $\vec{\mu}(t)$ solves (6.4.5).

Hint: Generalized product rule

$$S = S^\top : \frac{d}{dt} (\vec{p}(t)^\top S \vec{p}(t)) = 2 \dot{\vec{p}}(t)^\top S \vec{p}(t)$$



$$\triangleright \frac{dE}{dt}(t) = \dot{\vec{\mu}}(t)^\top \mathbf{M} \ddot{\vec{\mu}}(t) + \dot{\vec{\mu}}(t)^\top \mathbf{A} \vec{\mu}(t)$$

$$\text{MOL-QDE : } \mathbf{M} \ddot{\vec{p}} + \mathbf{B} \dot{\vec{p}} + \mathbf{A} \vec{p} = 0$$

$$\cdot \vec{p}^\top : \dot{\vec{\mu}}(t)^\top \mathbf{M} \ddot{\vec{\mu}}(t) + \dot{\vec{\mu}}(t)^\top \mathbf{A} \vec{\mu}(t) + \dot{\vec{\mu}}(t)^\top \mathbf{B} \dot{\vec{\mu}}(t) = 0$$

$$\triangleright \frac{dE}{dt}(t) = -\dot{\vec{\mu}}(t)^\top \mathbf{B} \dot{\vec{\mu}}(t) \leq 0$$

$$\dot{\vec{p}}(0) \neq 0 \Rightarrow \frac{dE}{dt}(t) < 0$$



