

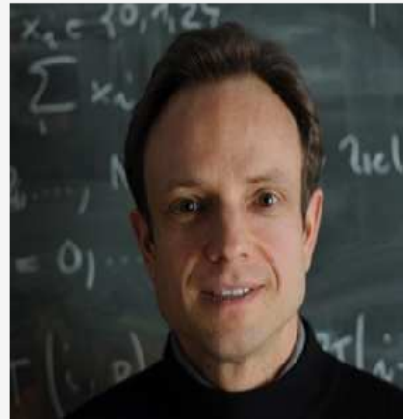
Homework Problem Video Tutorial

Problem 6-2: Implicit Timestepping for Parabolic IBVP

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Non-dimensional heat equation

$$\frac{\partial u}{\partial t} - \Delta u = 0 \quad \text{in } \Omega \times [0, T], \quad (\text{evolution equation})$$

$$-\text{grad } u \cdot n = cu \quad \text{on } \partial\Omega \times [0, T], \quad (\text{boundary conditions}) \quad (*) \quad (6.2.1)$$

$$u(x, 0) = u_0(x) \quad \text{in } \Omega, \quad (\text{initial conditions}) \quad c > 0$$

→ convective cooling b.c.

a, Spatial variational formulation

Test with $v \in H^1(\Omega)$ & integrate & i.b.p.

$$u: [0, T] \rightarrow H^1(\Omega);$$

$$\int_{\Omega} \dot{u} v \, dx + \int_{\Omega} \text{grad } u \cdot \text{grad } v \, dx - \int_{\partial\Omega} v \text{grad } u \cdot n \, dx = 0$$

$$\Rightarrow \underbrace{\int_{\Omega} \dot{u} v \, dx}_{m(\dot{u}, v)} + \underbrace{\int_{\Omega} \text{grad } u \cdot \text{grad } v \, dx}_{a(u, v)} + \underbrace{\int_{\partial\Omega} cu v \, dx}_{\ell(v)} = \underbrace{0}_{\ell(v)}, \quad \forall v \in H^1(\Omega) \quad (6.2.3)$$

$$\begin{array}{ccc} m(\dot{u}, v) & + & a(u, v) = \ell(t)(v) \\ \uparrow \text{s.p.d.} & & \uparrow \text{s.p.d.} \end{array}$$

b, Argue why the total thermal energy [heat capacity $S \equiv 1$]

$$\text{"Heat content": } E(t) := \int_{\Omega} u(x, t) \, dx, \quad (6.2.4)$$

decreases with time, if $u_0(x) > 0$ for all $x \in \Omega$.

② "Physical intuition": $u(x,t) \geq 0 \quad \forall (t,x)$
 $\Leftrightarrow$ **Maximum principle** for parabolic evolutions]

$$\begin{aligned} \frac{d}{dt} E(t) &= \frac{d}{dt} \int_{\Omega} u(x,t) dx \\ &= \int_{\Omega} \dot{u}(x,t) dx \end{aligned}$$

[Use variational formulation with $v \equiv 1$]

$$= -c \int_{\partial\Omega} \underbrace{u(x,t)}_{\geq 0} dS(x) \leq 0$$

(f $u|_{\partial\Omega} > 0 \Rightarrow t \rightarrow E(t)$ decrease.

c, Relevant bilinear forms

$$a(u,v) := \int_{\Omega} \text{grad } u \cdot \text{grad } v \, dx + \int_{\partial\Omega} cuv \, dx,$$

$$m(u,v) := \int_{\Omega} u(x)v(x) \, dx.$$

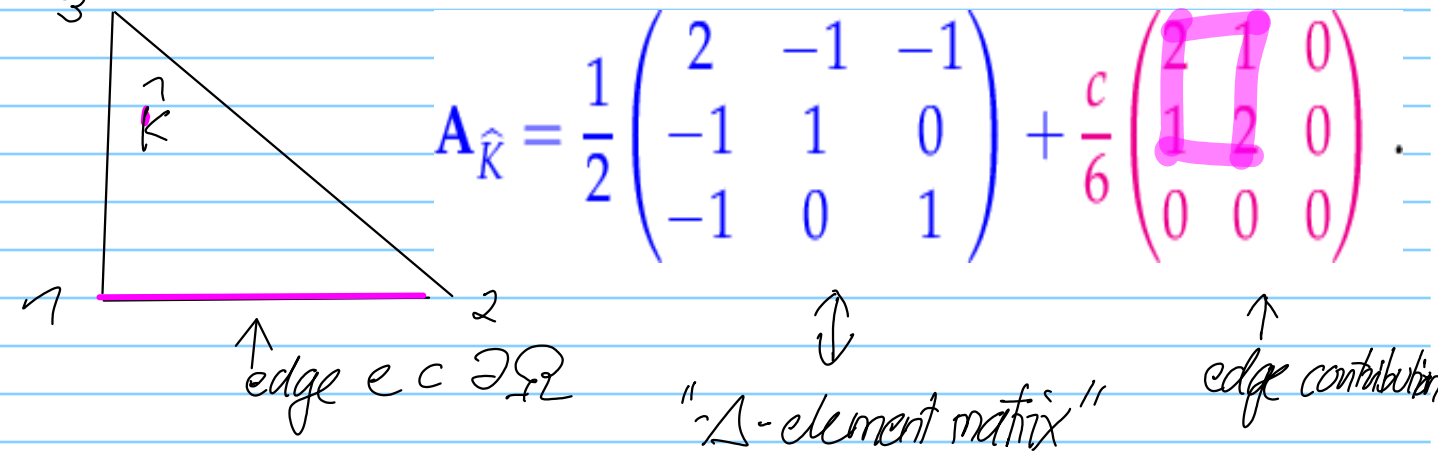
$S_1^0(\mathcal{M})$ Galerkin FEM

Here: Element mass matrix ($\Leftrightarrow m(\cdot, \cdot)$) on unit triangle $\hat{K}$,

LSF $\hat{=}$ barycentric coordinate functions $\lambda_1, \lambda_2, \lambda_3$
 $M_{\hat{K}} = \left[\int_{\hat{K}} \lambda_i \lambda_j \, dx \right]_{i,j=1}^3$

$$M_{\hat{K}} = \frac{1}{24} \begin{pmatrix} 2 & 1 & 1 \\ 1 & 2 & 1 \\ 1 & 1 & 2 \end{pmatrix}. \quad (6.2.6)$$

d, $a(\cdot, \cdot) - \hat{K}$ element matrix



$$A_{\hat{K}} = \frac{1}{2} \begin{pmatrix} 2 & -1 & -1 \\ -1 & 1 & 0 \\ -1 & 0 & 1 \end{pmatrix} + \frac{c}{6} \begin{pmatrix} 2 & 1 & 0 \\ 1 & 2 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

"Δ-element matrix"

f, Implement a LEHRFEM++-based C++ function (in the file `sdirkmethodoflines.cc`)

```
std::pair<Eigen::SparseMatrix<double>, Eigen::SparseMatrix<double>>
assembleGalerkinMatrices(const lf::assemble::DofHandler &dofh,
double c);
```

that assembles the $S_1^0(\mathcal{M})$ -Galerkin matrices for the bilinear forms $a(\cdot, \cdot)$ and $m(\cdot, \cdot)$ found in Sub-problem (6-2.a). The argument `dofh` is to provide a hybrid mesh $\mathcal{M}$ and information about the numbering of local and global basis functions. The coefficient $c > 0$ is given through `c`.

Your implementation can rely on the LEHRFEM++ library class `lf::uscalfe::LinearFELaplaceElementMatrix` and, of course, on `lf::assemble::assembleMatrix` see [Lecture $\rightarrow$ Code 2.7.4.23], which even accepts a co-dimension argument! You will also have to implement your own slightly extended version of the class `LinFEMassEdgeMatProvider` from [Lecture $\rightarrow$ Code 2.7.4.39].

An alternative is the use of `lf::uscalfe::ReactionDiffusionElementMatrixProvider` with suitable coefficient functions, see [Lecture $\rightarrow$ Ex. 2.8.3.28].

③ Implicit 2-stage RK-SSM : SDIRK-2

$$\begin{array}{c|c} c & A \\ \hline 1 & \begin{array}{cc} \lambda & 0 \\ 1-\lambda & \lambda \end{array} \\ \hline 1-\lambda & \lambda \end{array} \triangleq \begin{array}{c|c} \lambda & 0 \\ \hline 1 & \begin{array}{cc} \lambda & 0 \\ 1-\lambda & \lambda \end{array} \\ \hline 1-\lambda & \lambda \end{array}, \quad \lambda := 1 - \frac{1}{2}\sqrt{2}. \quad (6.2.8)$$

Definition 6.2.6.31. General Runge-Kutta method →

For coefficients $b_i, a_{ij} \in \mathbb{R}$, $c_i := \sum_{j=1}^s a_{ij}$, $i, j = 1, \dots, s$, $s \in \mathbb{N}$, the discrete evolution $\Psi^{s,t}$ of an s -stage Runge-Kutta single step method (RK-SSM) for the ODE $\dot{u} = f(t, u)$, is defined by

$$k_i := f(t + c_i \tau, u + \tau \sum_{j=1}^s a_{ij} k_j), \quad i = 1, \dots, s, \quad \Psi^{t, t+\tau} u := u + \tau \sum_{i=1}^s b_i k_i.$$

The $k_i \in V_0$ are called increments.

↑ discrete evolution operator

g, SDIRK-2 for $\dot{y} = -\gamma y =: f(y)$
[scalar linear ODE]



Increment eqn.:

$$\begin{aligned} k_1 &= -\gamma y - \gamma \tau \lambda k_1, \\ k_2 &= -\gamma y - \gamma \tau (1-\lambda) k_1 - \gamma \tau \lambda k_2 \end{aligned} \quad \approx 2 \times 2 \text{ LSE}$$

$$\triangleright \quad k_1 = \frac{-\gamma y}{1 + \gamma \tau \lambda}, \quad k_2 = \frac{-\gamma y (1 + 2\gamma \tau \lambda - \gamma \tau)}{(1 + \gamma \tau \lambda)^2} \quad (6.2.9)$$

$$\begin{aligned} y^{(j+1)} &= y^{(j)} + \tau (1-\lambda) k_1 + \tau \lambda k_2 \\ &= y^{(j)} \left(1 - \frac{\gamma \tau (1 + \gamma \tau \lambda^2)}{(1 + \gamma \tau \lambda)^2} \right). \end{aligned} \quad (6.2.10) \quad (6.2.11)$$

h, "Measuring" order of SDIRK-2 RK-SSM
= rate of alg. cvg.

Model problem: $\dot{y} = -\gamma y$ on $[0, 2]$, $y_0 = 1$

ℓ	m_2 #timesteps	error ε_ℓ	estimated rate
0	10	0.000446558	
1	20	0.000110501	2.01478
2	40	2.74922e-05	2.00697
3	80	6.85696e-06	2.00338
4	160	1.71226e-06	2.00167
5	320	4.2782e-07	2.00083
6	640	1.06924e-07	2.00041
7	1280	2.67273e-08	2.00021
8	2560	6.68136e-09	2.0001
9	5120	1.67028e-09	2.00005

reduction by
factor ≈ 4

Alg. cvg.

$$\varepsilon_\ell \approx C m_\ell^{-\alpha}$$

$$\log \varepsilon_\ell \approx \log C - \alpha \log m_\ell$$

$$\log \frac{\varepsilon_\ell}{\varepsilon_{\ell-1}} \approx -\alpha \log \frac{m_\ell}{m_{\ell-1}}$$

$$\alpha \approx - \frac{\log \frac{\varepsilon_\ell}{\varepsilon_{\ell-1}}}{\log 2}$$

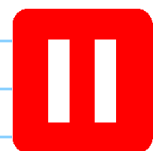
f, Show: SDIRK-2 is $L(\pi)$ -stable

$$S(-\tau\gamma) := \left(1 - \frac{\gamma \tau (1 + \gamma \tau \lambda^2)}{(1 + \gamma \tau \lambda)^2} \right) \Leftrightarrow S(z) = 1 - \frac{-z(1 - z\lambda^2)}{(1 - z\lambda)^2}. \quad (6.2.16)$$

→ stability function

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$L(\pi)$ -stable $\Leftrightarrow |S(z)| < 1, \forall z < \infty$ ✓
 $\lim_{z \rightarrow -\infty} S(z) = 0$ ✓



Compute, $v_n \in S_1^0(\mathcal{M})$
 $\int_{\Omega} v_n dx = \sum_K \sum_{e=1}^3 v_n(a_k^e) \int \lambda_e dx$
 entries of coefficient vector $\vec{v}$ of v_n = $\frac{1}{3|K|}$

7, SDIRK-2 for MOL-ODE

$$\mathbf{M} \left\{ \frac{d}{dt} \vec{\mu}(t) \right\} + \mathbf{A} \vec{\mu}(t) = \vec{\varphi}(t) \Leftrightarrow \underbrace{\dot{\vec{\mu}} = \mathbf{M}^{-1}(\vec{\varphi}(t) - \mathbf{A} \vec{\mu}(t))}_{=f(t, \vec{\mu})} \quad (7.1.4.4)$$

apply RK-SSM formulas

Increment equations LSE

$$(7.1.7.7) \Leftrightarrow (\mathbf{I}_s \otimes \mathbf{M} + \tau \mathbf{A} \otimes \mathbf{A}) \begin{bmatrix} \vec{\kappa}_1 \\ \vdots \\ \vec{\kappa}_s \end{bmatrix} = \begin{bmatrix} \vec{\varphi}(t_j + c_1 \tau) - \mathbf{A} \vec{\mu}^{(j)} \\ \vdots \\ \vec{\varphi}(t_j + c_s \tau) - \mathbf{A} \vec{\mu}^{(j)} \end{bmatrix} \quad (6.2.7.9)$$

↑ ↑
Kronecker product of matrices



$$\mathbf{M} \vec{\kappa}_1 + \tau \lambda \mathbf{A} \vec{\kappa}_1 = -\mathbf{A} \vec{\mu}^{(j)},$$

(6.2.19)

$$\mathbf{M} \vec{\kappa}_2 + \tau(1 - \lambda) \mathbf{A} \vec{\kappa}_1 + \tau \lambda \mathbf{A} \vec{\kappa}_2 = -\mathbf{A} \vec{\mu}^{(j)}.$$

$\triangleq$ A staggered LSE

C++11 code 6.2.22: Sub-problem (6-2.k), function `thermalEnergy()` → GITLAB

```
2 double thermalEnergy(const If::assemble::DofHandler &dofh,
3                       const Eigen::VectorXd &temperature_vec) {
4     double thermal_energy = 0.0;
5     /* SOLUTION_BEGIN */
6     auto mesh_p = dofh.Mesh(); // pointer to mesh
7
8     // The thermal energy of the system is defined as the integral of the
9     // temperature solution over the domain. It is computed for linear
10    // lagrangian
11    // finite elements using the trapezoidal rule by summing up the
12    // contribution
13    // of that quadrature rule over each triangle
14    double thermal_energy_loc;
15    for (const If::mesh::Entity &tria : mesh_p->Entities(0)) {
16        thermal_energy_loc = 0.0;
17        // Compute the area of the triangle
18        const double area = If::geometry::Volume(*(tria.Geometry()));
19        // Obtain the global indices of the nodal degrees of freedom
20        for (const If::assemble::gdof_idx_t &g_idx :
21             dofh.GlobalDofIndices(tria)) {
22            thermal_energy_loc += temperature_vec[g_idx];
23        }
24        // Sum local contribution of quadrature rule
25        thermal_energy += thermal_energy_loc * area / 3.0;
26    }
27    /* SOLUTION_END */
28    return thermal_energy;
29 } // thermalEnergy
```


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In the file `sdirkmethodoflines.cc` implement a C++ function

```
std::pair<Eigen::VectorXd, Eigen::VectorXd>
solveTemperatureEvolution(const lf::assemble::DofHandler &dofh,
    unsigned int m, double cool_coeff, [→  $c \in \mathbb{R}$ ]
    Eigen::VectorXd initial_temperature_vec)
```

that carries out m equal timesteps of the SDIRK-2 in order to solve (6.2.1) over the time interval $[0, 1]$. The argument `initial_temperature_vec` provides the initial conditions for the method-of-lines ODE, cf. [Lecture → Eq. (6.2.4.4)], the parameter `cool_coeff` the coefficient $c \in \mathbb{R}$. The discretization in space relies on $S_1^0(\mathcal{M})$ -FEM.

The function should return the basis expansion coefficient vector of the solution at final time $T = 1$ and the “energies” (6.2.25) for times $\tau j, j = 0, \dots, m$, where $\tau := 1/m$ is the constant timestep size.

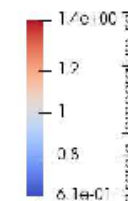
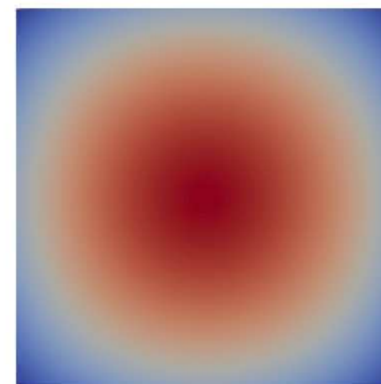
In every time step we have to solve

$$\mathbf{M}\vec{\kappa}_1 + \tau\lambda\mathbf{A}\vec{\kappa}_1 = -\mathbf{A}\vec{\mu}^{(j)}, \quad (6.2.19)$$

$$\mathbf{M}\vec{\kappa}_2 + \tau(1-\lambda)\mathbf{A}\vec{\kappa}_1 + \tau\lambda\mathbf{A}\vec{\kappa}_2 = -\mathbf{A}\vec{\mu}^{(j)}.$$

▷ Two $N \times N$ LSEs, $N := \dim S_1^0(\mathcal{M})$ with the same system matrix $\mathbf{M} + \tau\lambda\mathbf{A}$:
[the same for all timesteps] $\uparrow$ fixed

▷ Assemble $\mathbf{A}, \mathbf{M}$ & LU-factorize $\mathbf{M} + \tau\lambda\mathbf{A}$ before entering the timestepping loop.



Δu at $T = 1$
 $M_b \equiv \nearrow$

▷ Material has cooled close to $\partial\Omega$.

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