

Lecture 10: Viscoelastic modulus reconstruction and full-field OCT elastography

Habib Ammari

Department of Mathematics, ETH Zürich

Elastography

- **Elastography**: quantitative visualization of the mechanical properties of human tissues by using the relation between the wave propagation velocity and the mechanical properties of the tissues.
- Mechanical properties of tissue include the shear modulus, shear viscosity, and compression modulus.
- Quantification of the tissue shear modulus in vivo can provide evidence of the manifestation of tissue diseases.
- Image reconstruction methods for tissue viscoelasticity imaging.
- Recover the distribution of the **shear modulus** (μ) and shear viscosity (η) from the internal measurement of the time-harmonic mechanical displacement field u produced by the application of an external time harmonic excitation at frequency $\omega/2\pi$ in the range $50 \sim 200\text{Hz}$ through the surface of the subject.
- Modeling soft tissue as being linearly viscoelastic and **nearly incompressible**.

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- **Displacement:** $\Re(u(x)e^{i\omega t})$,

$$\nabla \cdot \left((\mu + i\omega\eta)(\nabla u + \nabla u^T) \right) + \nabla((\lambda + i\omega\eta_\lambda)\nabla \cdot u) + \rho\omega^2 u = 0;$$

ρ : density of the medium, ∇u^T : transpose of ∇u , λ : compression modulus and η_λ : compression viscosity.

- **Algebraic inversion** method: For any non-zero constant vector a ,

$$\mu + i\omega\eta = -\frac{\rho\omega^2(a \cdot u)}{\nabla \cdot \nabla(a \cdot u)}.$$

- Strong assumptions of $\nabla(\mu + i\omega\eta) \approx 0$ (**local homogeneity**) and $(\lambda + i\omega\eta_\lambda)\nabla \cdot u \approx 0$ (**negligible pressure**).
- Algebraic formula: **ignores reflection effects** of the propagating wave due to abrupt changes of $\mu + i\omega\eta$, so that the method cannot measure any change of $\mu + i\omega\eta$ in the direction of a .

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- **Boundary conditions:** Γ_D and Γ_N s.t. $\overline{\Gamma_D \cup \Gamma_N} = \partial\Omega$ and $\Gamma_D \cap \Gamma_N = \emptyset$,

$$u = g \quad \text{on } \Gamma_D,$$

$$2(\mu + i\omega\eta)\nabla^s u \nu + (\lambda + i\omega\eta\lambda)(\nabla \cdot u)\nu = 0 \quad \text{on } \Gamma_N.$$

- Soft tissues: **nearly incompressible**.
- $\lambda \approx \infty \Rightarrow \nabla \cdot u \approx 0$.
- Internal pressure $p = \lim_{\lambda \rightarrow +\infty} \lambda \nabla \cdot u$.

- **Quasi-incompressible viscoelasticity model:**

$$\begin{cases} 2\nabla \cdot ((\mu + i\omega\eta)\nabla^s u) + \nabla p + \rho\omega^2 u = 0 & \text{in } \Omega, \\ \nabla \cdot u = 0 & \text{in } \Omega, \\ u = g & \text{on } \Gamma_D, \\ 2(\mu + i\omega\eta)\nabla^s u \nu + p\nu = 0 & \text{on } \Gamma_N. \end{cases}$$

- If $\Gamma_D = \partial\Omega$ ($\Gamma_N = \emptyset$), then g should satisfy the compatibility condition $\int_{\partial\Omega} g \cdot \nu \, ds = 0$.
- $u^{(m)} = u^{(m)}[\mu_*, \eta_*]$: displacement measured in Ω ; μ_* and η_* : true distributions of shear elasticity and viscosity.
- Inverse problem: reconstruct the distribution of μ and η from $u^{(m)}$

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- For a fixed $\epsilon > 0$; $\Omega' := \{x \in \Omega \mid \text{dist}(x, \partial\Omega) > \epsilon\}$ and $\mathcal{E} := \Omega \setminus \overline{\Omega'}$.
- μ and $\eta \in \tilde{\mathcal{S}} := \{(\mu_0, \eta_0) + (\phi_1, \phi_2) \mid (\phi_1, \phi_2) \in \mathcal{S}\}$; μ_0 and η_0 : positive constants.
- $\tilde{\mathcal{S}} = (\mu_0, \eta_0) + \mathcal{S}$:

$$\mathcal{S} := \{(\phi_1, \phi_2) \in H_0^2(\Omega) \times H_0^2(\Omega) : c_1 < \phi_1 + \mu_0 < c_2,$$

$$c_1 < \phi_2 + \eta_0 < c_2, \|\phi_j\|_{W^{2,2}(\Omega)} \leq c_3, \text{supp } \phi_j \subset \Omega' \text{ for } j = 1, 2\};$$

c_1, c_2, c_3 : positive constants.

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- **Optimal control algorithm.**
- Discrepancy functional:

$$J[\mu, \eta] = \frac{1}{2} \int_{\Omega} |u[\mu, \eta] - u^{(m)}|^2 dx.$$

- **Fréchet derivatives** of $J[\mu, \eta]$ with respect to μ and η .
- Assume that δ_{μ} and δ_{η} : small perturbations of μ and η , respectively, by regarding $\frac{\delta\mu + i\omega\delta\eta}{\mu + i\omega\eta} \approx 0$.
- $u_0 := u[\mu, \eta]$, $p_0 :=$ the pressure corresponding to u_0 and $p_0 + p_1 :=$ the pressure corresponding to $u[\mu + \delta_{\mu}, \eta + \delta_{\eta}]$;

$$\delta u := u[\mu + \delta_{\mu}, \eta + \delta_{\eta}] - u_0.$$

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$$2\nabla \cdot ((\mu + i\omega\eta)\nabla^s \delta u) + \nabla p_1 + \rho\omega^2 \delta u = -2\nabla \cdot ((\delta_{\mu} + i\omega\delta_{\eta})\nabla^s u_0) - 2\nabla \cdot ((\delta_{\mu} + i\omega\delta_{\eta})\nabla^s \delta u) \quad \text{in } \Omega.$$

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- u_1 : solution of

$$\begin{cases} 2\nabla \cdot ((\mu + i\omega\eta)\nabla^s u_1) + \nabla p_1 + \rho\omega^2 u_1 = & \text{in } \Omega, \\ -2\nabla \cdot ((\delta_\mu + i\omega\delta_\eta)\nabla^s u_0) & \\ \nabla \cdot u_1 = 0 & \text{in } \Omega, \\ u_1 = 0 & \text{on } \Gamma_D, \\ 2(\mu + i\omega\eta)\nabla^s u_1 \nu + p_1 \nu = 0 & \text{on } \Gamma_N. \end{cases}$$

- For $(\delta\mu + \mu, \delta\eta + \eta) \in \tilde{\mathcal{S}}$,

$$\Re \int_{\Omega} u_1 \overline{(u_0 - u^{(m)})} dx = \Re \int_{\Omega} 2(\delta_\mu + i\omega\delta_\eta)\nabla^s u_0 : \nabla^s \bar{v} dx.$$

- **Fréchet derivatives** of $J[\mu, \eta]$ with respect to μ and η :

$$\frac{\partial}{\partial \mu} J[\mu, \eta] = \Re [2\nabla^s u_0 : \nabla^s \bar{v}], \quad \frac{\partial}{\partial \eta} J[\mu, \eta] = \Re [2(i\omega\nabla^s u_0) : \nabla^s \bar{v}];$$

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- v solution of the **adjoint problem**:

$$\begin{cases} 2\nabla \cdot ((\mu - i\omega\eta)\nabla^s v) + \nabla q + \rho\omega^2 v = (u_0 - u^{(m)}) & \text{in } \Omega, \\ \nabla \cdot v = 0 & \text{in } \Omega, \\ v = 0 & \text{on } \Gamma_D, \\ 2(\mu - i\omega\eta)\nabla^s v \nu + q\nu = 0 & \text{on } \Gamma_N. \end{cases}$$

- $J[\mu, \eta]$: **Fréchet differentiable** for $(\mu, \eta) \in \tilde{S}$.
- As $\delta_\mu, \delta_\eta \rightarrow 0$,

$$\begin{aligned} & \left| J[\mu + \delta_\mu, \eta + \delta_\eta] - J[\mu, \eta] - \Re \int_{\Omega} u_1(\overline{u_0 - u^{(m)}}) dx \right| \\ & = O\left(\left(\|\delta_\mu\|_{W^{2,2}(\Omega)} + \|\delta_\eta\|_{W^{2,2}(\Omega)}\right)^2\right). \end{aligned}$$

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- **Gradient descent scheme:**

- Let $n = 0$. Start with an initial guess of shear modulus $\mu^{(0)}$ and shear viscosity $\eta^{(0)}$.
- For $n = 0, 1, \dots$, compute $u_0^{(n)}$ by solving the **forward problem** with μ and η replaced by $\mu^{(n)}$ and $\eta^{(n)}$, respectively. Compute $v^{(n)}$ by solving the **adjoint problem** with μ, η, u_0 replaced by $\mu^{(n)}, \eta^{(n)}, u_0^{(n)}$, respectively.
- For $n = 0, 1, \dots$, compute the **Fréchet derivatives** $\frac{\partial J}{\partial \mu}[\mu^{(n)}, \eta^{(n)}]$ and $\frac{\partial J}{\partial \eta}[\mu^{(n)}, \eta^{(n)}]$.
- Update μ and η as follows:

$$\begin{cases} \mu^{(n+1)} &= \mu^{(n)} - \delta \frac{\partial J}{\partial \mu}[\mu^{(n)}, \eta^{(n)}], \\ \eta^{(n+1)} &= \eta^{(n)} - \delta \frac{\partial J}{\partial \eta}[\mu^{(n)}, \eta^{(n)}], \end{cases}$$

δ : **step size**.

- Repeat Steps 2, 3, and 4 until $\|\mu^{(n+1)} - \mu^{(n)}\| \leq \epsilon$ and $\|\eta^{(n+1)} - \eta^{(n)}\| \leq \epsilon$ for a given $\epsilon > 0$.

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- **Initial guess:**

- Ignore the pressure term:

$$2\nabla \cdot (\mu + i\omega\eta)\nabla^s u^\diamond + \rho\omega^2 u^\diamond = 0 \quad \text{in } \Omega,$$

- **Helmholtz decomposition:**

$$(\mu + i\omega\eta)\nabla^s u^\diamond = \nabla f + \nabla \times W \quad \text{with } \nabla \cdot W = 0,$$

- f and W : vector and matrix, respectively.

$$\mu + i\omega\eta = \frac{\nabla f : \nabla^s \bar{u}^\diamond}{|\nabla^s u^\diamond|^2} + \frac{\nabla \times W : \nabla^s \bar{u}^\diamond}{|\nabla^s u^\diamond|^2}.$$

$$\Delta f = -\frac{1}{2}\rho\omega^2 u^\diamond \quad \text{in } \Omega.$$

$$\Delta W = \nabla \times ((\mu + i\omega\eta)\nabla^s u^\diamond) \quad \text{in } \Omega.$$

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- \tilde{f} :

$$\begin{cases} \Delta \tilde{f} = -\frac{1}{2} \rho \omega^2 \mathbf{u}^{(m)} & \text{in } \Omega, \\ \nabla \tilde{f} \nu = (\mu_0 + i\omega\eta_0) \nabla^s \mathbf{u}^{(m)} \nu & \text{on } \partial\Omega. \end{cases}$$

- W_1 :

$$\begin{cases} \Delta W_1 = \nabla \times \left(\frac{\nabla \tilde{f} \cdot \nabla^s \mathbf{u}^{(m)}}{|\nabla^s \mathbf{u}^{(m)}|^2} \nabla^s \mathbf{u}^{(m)} \right) & \text{in } \Omega, \\ W_1 = 0 & \text{on } \partial\Omega. \end{cases}$$

- W_2 :

$$\begin{cases} \Delta W_2 = \nabla \times \left(-\frac{\rho \omega^2 (\mathbf{a} \cdot \mathbf{u}^{(m)})}{\nabla \cdot \nabla (\mathbf{a} \cdot \mathbf{u}^{(m)})} \nabla^s \mathbf{u}^{(m)} \right) & \text{in } \Omega, \\ W_2 = 0 & \text{on } \partial\Omega, \end{cases}$$

\mathbf{a} : any nonzero vector.

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- Initial guess formula:

$$\mu^{(0)} + i\omega\eta^{(0)} = \frac{\nabla \tilde{f} : \nabla^s \bar{u}^{(m)}}{|\nabla^s u^{(m)}|^2} + \frac{\nabla \times (W_1 + W_2) : \nabla^s \bar{u}^{(m)}}{2|\nabla^s u^{(m)}|^2} .$$

Elastography

- **Local reconstruction:**

- $\Omega_{\text{loc}} \Subset \Omega$; Localized minimization problem:

$$J_{\text{loc}}[\mu, \eta] = \frac{1}{2} \int_{\Omega_{\text{loc}}} |u_{\text{loc}}[\mu, \eta] - u^{(m)}|^2 dx;$$

- $u_{\text{loc}}[\mu, \eta]$:

$$\begin{cases} 2\nabla \cdot ((\mu + i\omega\eta)\nabla^s u) + \nabla p + \rho\omega^2 u = 0 & \text{in } \Omega_{\text{loc}}, \\ \nabla \cdot u = 0 & \text{in } \Omega_{\text{loc}}, \\ u = u^{(m)} & \text{on } \partial\Omega_{\text{loc}}. \end{cases}$$

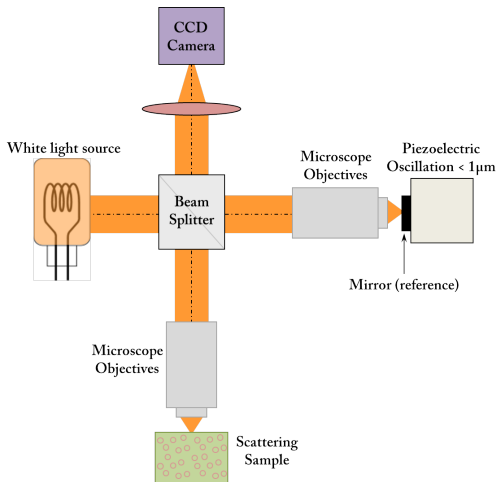
- **Adjoint problem:**

$$\begin{cases} 2\nabla \cdot ((\mu - i\omega\eta)\nabla^s v) + \nabla q + \rho\omega^2 v = u_{\text{loc}} - u^{(m)} & \text{in } \Omega_{\text{loc}}, \\ \nabla \cdot v = 0 & \text{in } \Omega_{\text{loc}}, \\ v = 0 & \text{on } \partial\Omega_{\text{loc}}. \end{cases}$$

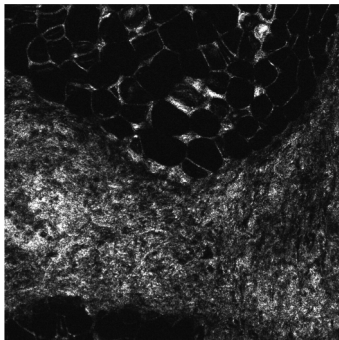
Full-field optical coherence elastography

- Full-field optical coherence tomography (OCT): **optical image** with **sub-cellular resolution**.
- Apply a load on the sample.
- **OCTE**: Use a set of optical images **before** and **after** mechanical solicitation to reconstruct the **shear modulus** distribution inside the sample.
- Map of mechanical properties: added as a **supplementary contrast** mechanism to **enhance specificity**.

Full-field optical coherence elastography



Full-field optical coherence elastography



Full-field optical coherence elastography

- Reconstruct the **shear modulus** μ from ε and ε_u .
- $\varepsilon(x) = \varepsilon_u(x + u(x))$;
- Displacement field u :

$$\begin{cases} \nabla \cdot (\mu(\nabla u + \nabla u^T)) + \nabla p = 0 & \text{in } \Omega, \\ \nabla \cdot u = 0 & \text{in } \Omega, \\ u = f & \text{on } \partial\Omega. \end{cases}$$

Full-field optical coherence elastography

- $BV(\Omega)$: subspace of $L^1(\Omega)$ of all the functions f whose weak derivative Df : a finite Radon measure:
- f s.t.

$$\int_{\Omega} f \nabla \cdot F \leq C \sup_{x \in \Omega} |F|, \quad \forall F \in \mathcal{C}_0^1(\Omega)^d$$

for some positive constant C with $\mathcal{C}_0^1(\Omega)$: set of compactly supported \mathcal{C}^1 functions.

Full-field optical coherence elastography

- Derivative of a function $f \in BV(\Omega)$ can be decomposed as

$$Df = \nabla f \mathcal{H}^d + [f] \nu_S \mathcal{H}_S^{d-1} + D_c f;$$

- \mathcal{H}^d : Lebesgue measure on Ω , \mathcal{H}_S^{d-1} : surface Hausdorff measure on a rectifiable surface S , ν_S : normal vector defined a.e. on S ;
- $\nabla f \in L^1(\Omega)$: smooth derivative of f , $[f] \in L^1(S, \mathcal{H}_S^{d-1})$: jump of f across S and $D_c f$: vector measure supported on a set of Hausdorff dimension less than $(d - 1)$.

Full-field optical coherence elastography

- $SBV(\Omega)$: subspace of $BV(\Omega)$ of all the functions f satisfying $D_c f = 0$.
- For any $1 \leq p \leq +\infty$,

$$SBV^p(\Omega) = \left\{ f \in SBV(\Omega) \cap L^p(\Omega), \nabla f \in L^p(\Omega)^d \right\}.$$

Full-field optical coherence elastography

- u^* : true displacement; $\tilde{\varepsilon}$: measured deformed optical:

$$\tilde{\varepsilon} = \varepsilon \circ (\mathbb{I} + u^*)^{-1}.$$

- Optical image: **discontinuous**.
- **Optimal control algorithm**:

$$I(u) = \frac{1}{2} \int_{\Omega} |\tilde{\varepsilon} \circ (\mathbb{I} + u) - \varepsilon|^2 dx.$$

- I has a **nonempty subgradient**.
- $\xi \in \partial I$:

$$\xi : h \mapsto \int_{\Omega} [\tilde{\varepsilon}(x + u) - \varepsilon(x)] h(x) \cdot D\tilde{\varepsilon} \circ (\mathbb{I} + u)(x) dx.$$

Full-field optical coherence elastography

- Nondifferentiable functional $u \mapsto I(u)$ has a **nonempty subgradient** if there exists ξ s.t.

$$I(u + h) - I(u) \geq (\xi, h)$$

holds for $\|h\|$ small enough;

- $\xi \in \partial I$ with ∂I : subgradient of I .
- In order to **minimize** I , it is sufficient to **find one** $\xi \in \partial I$.

Full-field optical coherence elastography

Initial guess:

- Detect the **surface of jumps** of the optical image (edge detection algorithm).
- **Local recovery by linearization:** $\text{data} = \varepsilon - \varepsilon_u (\approx \nabla \varepsilon \cdot u)$

$$J_x(u) = \int_{\Omega} |\nabla \varepsilon(y) \cdot u - \text{data}(y)|^2 w_{\delta}(|x - y|) dy.$$

- $w_{\delta} = \frac{1}{\delta^d} w\left(\frac{\cdot}{\delta}\right)$; w : a mollifier supported on $[-1, 1]$.
- **Least-squares solution:**

$$u^T = \left(\int_{\Omega} w_{\delta}(|x - y|) \nabla \varepsilon(y) \nabla \varepsilon^T(y) dy \right)^{-1} \int_{x+\delta B} \text{data} w_{\delta}(|x - y|) \nabla \varepsilon(y) dy.$$

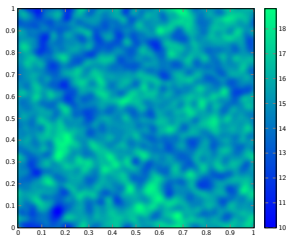
- If all vectors $\nabla \varepsilon$ in $\{y : w_{\delta}(|y - x|) \neq 0\}$ **not collinear**, then

$$\int_{\Omega} w_{\delta}(|x - y|) \nabla \varepsilon(y) \nabla \varepsilon^T(y) dy \quad \text{invertible.}$$

- **Resolution = variation of ε .**

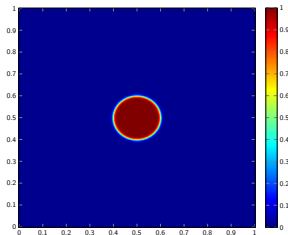
Full-field optical coherence elastography

Optical image ε of the medium:

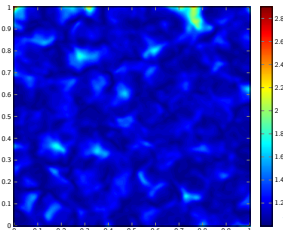


Full-field optical coherence elastography

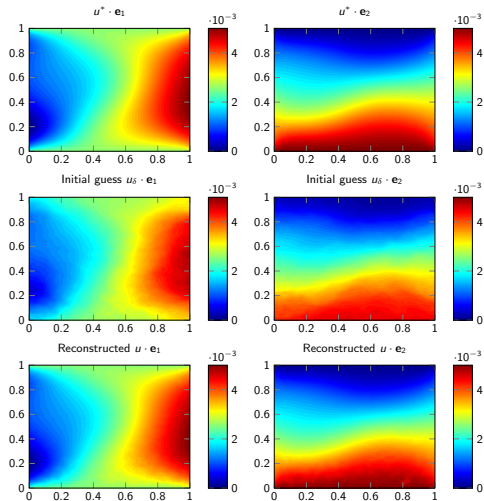
Averaging kernel w_δ :



Conditioning of the matrix $w_\delta \star \nabla \epsilon \nabla \epsilon^T$:

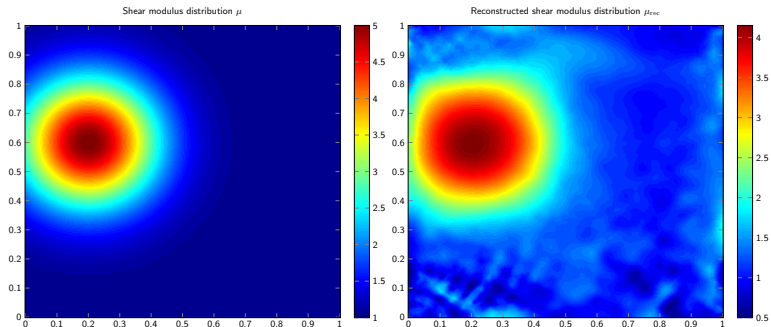


Full-field optical coherence elastography



Displacement field and its reconstruction.

Full-field optical coherence elastography



Shear modulus reconstruction.