Lecture 12: Plasmonic nanoparticle imaging

Habib Ammari

Department of Mathematics, ETH Zürich

Habib Ammari

Mathematics of super-resolution biomedical imaging

- Gold nano-particles: accumulate selectively in tumor cells; bio-compatible; reduced toxicity.
- Detection: localized enhancement in radiation dose (strong scattering).
- Ablation: localized damage (strong absorption).
- Functionalization: targeted drugs.





M.A. El-Sayed et al.

- D: nanoparticle in \mathbb{R}^d , d = 2, 3; $\mathcal{C}^{1,\alpha}$ boundary ∂D , $\alpha > 0$.
- ε_c(ω): complex permittivity of D; ε_m > 0: permittivity of the background medium;
- Permittivity contrast: λ(ω) = (ε_c(ω) + ε_m)/(2(ε_c(ω) − ε_m)).
- Causality \Rightarrow Kramer-Krönig relations (Hilbert transform); Drude model for the dielectric permittivity $\varepsilon_c(\omega)$.
- G: Fundamental solution to the Laplacian; S_D : Single-layer potential.
- Neumann-Poincaré operator K^{*}_D:

$$\mathcal{K}^*_D[\varphi](x) := \int_{\partial D} rac{\partial \mathcal{G}}{\partial
u(x)} (x-y) \varphi(y) \, ds(y) \ , \quad x \in \partial D.$$

 ν : normal to ∂D .

K^{*}_D: compact operator on *L*²(∂*D*); Spectrum of *K*^{*}_D lies in] − ¹/₂, ¹/₂] (Kellog).

- \mathcal{K}_D^* self-adjoint on $L^2(\partial D)$ iff D is a disk or a ball.
- Symmetrization technique for Neumann-Poincaré operator \mathcal{K}_D^* :
 - Calderón's identity: $\mathcal{K}_D \mathcal{S}_D = \mathcal{S}_D \mathcal{K}_D^*$;
 - In three dimensions, \mathcal{K}_D^* : self-adjoint in the Hilbert space $\mathcal{H}^*(\partial D) = H^{-\frac{1}{2}}(\partial D)$ equipped with

$$(u, v)_{\mathcal{H}^*} = -(u, \mathcal{S}_D[v])_{-\frac{1}{2}, \frac{1}{2}}$$

 $(\cdot, \cdot)_{-\frac{1}{2}, \frac{1}{2}}$: duality pairing between $H^{-\frac{1}{2}}(\partial D)$ and $H^{\frac{1}{2}}(\partial D)$.

• In two dimensions: $\exists ! \widetilde{\varphi}_0 \text{ s.t. } \mathcal{S}_D[\widetilde{\varphi}_0] = \text{constant on } \partial D \text{ and}$ $(\widetilde{\varphi_0}, 1)_{-\frac{1}{2}, \frac{1}{2}} = 1. \ \mathcal{S}_D \to \widetilde{\mathcal{S}}_D:$

$$\widetilde{\mathcal{S}}_{D}[\varphi] = \left\{ egin{array}{c} \mathcal{S}_{D}[\varphi] & ext{if } (arphi,1)_{-rac{1}{2},rac{1}{2}} = 0, \ -1 & ext{if } arphi = \widetilde{arphi}_{0}. \end{array}
ight.$$

- Symmetrization technique for Neumann-Poincaré operator *K*^{*}_D:
 - Spectrum $\sigma(\mathcal{K}_D^*)$ discrete in] -1/2, 1/2[;
 - Ellipse: $\pm \frac{1}{2} \left(\frac{a-b}{a+b}\right)^{j}$, elliptic harmonics (a, b): long and short axis).
 - Ball: $\frac{1}{2(2i+1)}$, spherical harmonics.
 - Twin property in two dimensions;
 - $(\lambda_j, \varphi_j), j = 0, 1, 2, \ldots$: eigenvalue and normalized eigenfunction pair of \mathcal{K}_D^* in $\mathcal{H}^*(\partial D); \lambda_j \in]-\frac{1}{2}, \frac{1}{2}]$ and $\lambda_j \to 0$ as $j \to \infty$;
 - φ₀: eigenfunction associated to 1/2 (φ̃₀ multiple of φ₀);
 - Spectral decomposition formula in $H^{-1/2}(\partial D)$,

$$\mathcal{K}_D^*[\psi] = \sum_{j=0}^{\infty} \lambda_j (\psi, \varphi_j)_{\mathcal{H}^*} \varphi_j.$$

• u^i : incident plane wave; Helmholtz equation:

$$\left\{\begin{array}{l} \nabla \cdot \left(\varepsilon_m \chi(\mathbb{R}^d \setminus \bar{D}) + \varepsilon_c(\omega) \chi(\overline{D})\right) \nabla u + \omega^2 u = 0,\\ u^s := u - u^i \text{ satisfies the outgoing radiation condition} \end{array}\right.$$

• Uniform small volume expansion with respect to the contrast: $D = z + \delta B$, $\delta \to 0$, $|x - z| \gg 2\pi/k_m$,

$$u^{s} = -M(\lambda(\omega), D) \nabla_{z} G_{k_{m}}(x-z) \cdot \nabla u^{i}(z) + O(\frac{\delta^{d+1}}{\operatorname{dist}(\lambda(\omega), \sigma(\mathcal{K}_{D}^{*}))}).$$

- G_{k_m} : outgoing fundamental solution to $\Delta + k_m^2$; $k_m := \omega/\sqrt{\varepsilon_m}$;
- Polarization tensor:

$$M(\lambda(\omega), D) := \int_{\partial D} x(\lambda(\omega)I - \mathcal{K}_D^*)^{-1}[\nu](x) \, ds(x).$$

• Spectral decomposition: (1, m)-entry

$$M_{l,m}(\lambda(\omega),D) = \sum_{j=1}^{\infty} \frac{(\nu_m,\varphi_j)_{\mathcal{H}^*}(\nu_l,\varphi_j)_{\mathcal{H}^*}}{(1/2-\lambda_j)(\lambda(\omega)-\lambda_j)}$$

- (ν_m, φ₀)_{H*} = 0; φ₀: eigenfunction of K^{*}_D associated to 1/2.
- Quasi-static plasmonic resonance: dist(λ(ω), σ(K^{*}_D)) minimal (ℜe ε_c(ω) < 0).

•
$$M(\lambda(\omega), B) = \left(\frac{\varepsilon_c(\omega)}{\varepsilon_m} - 1\right) \int_B \nabla v(y) dy$$
:

$$\begin{cases} \nabla \cdot \left(\varepsilon_m \chi(\mathbb{R}^d \setminus \overline{B}) + \varepsilon_c(\omega) \chi(\overline{B})\right) \nabla v = 0, \\ v(y) - y \to 0, \quad |y| \to +\infty. \end{cases}$$

• Corrector *v*:

$$\mathbf{v}(\mathbf{y}) = \mathbf{y} + \mathcal{S}_B(\lambda(\omega)\mathbf{I} - \mathcal{K}_B^*)^{-1}[\mathbf{\nu}](\mathbf{y}), \quad \mathbf{y} \in \mathbb{R}^d.$$

• Inner expansion: $\delta \rightarrow 0$, $|x - z| = O(\delta)$,

$$u(x) = u^{i}(z) + \delta v(\frac{x-z}{\delta}) \cdot \nabla u^{i}(z) + O(\frac{\delta^{2}}{\operatorname{dist}(\lambda(\omega), \sigma(\mathcal{K}_{D}^{*}))}).$$

• Monitoring of temperature elevation due to nanoparticle heating:

$$\begin{cases} \rho C \frac{\partial T}{\partial t} - \nabla \cdot \tau \nabla T = \frac{\omega}{2\pi} (\Im(\varepsilon_c(\omega)) |u|^2 \chi(D), \\ T|_{t=0} = 0. \end{cases}$$

 ρ : mass density; C: thermal capacity; τ : thermal conductivity.

• Scattering amplitude:

$$u^{s}(x) = -ie^{-\frac{\pi i}{4}} \frac{e^{ik_{m}|x|}}{\sqrt{8\pi k_{m}|x|}} A_{\infty}[D, \varepsilon_{c}, \varepsilon_{m}, \omega](\theta, \theta') + o(|x|^{-\frac{1}{2}}),$$

 $|x| \rightarrow \infty; \, \theta, \, \theta'$: incident and scattered directions.

• Scattering cross-section:

$$Q^{s}[D,\varepsilon_{c},\varepsilon_{m},\omega](\theta'):=\int_{0}^{2\pi}\left|A_{\infty}[D,\varepsilon_{c},\varepsilon_{m},\omega](\theta,\theta')\right|^{2}d\theta.$$

• Enhancement of Q^s at plasmonic resonances:

 $Q^s \propto \left| \operatorname{tr}(M(\lambda(\omega), D)) \right|^2.$



Norm of the polarization tensor for two disks for various separating distances.

- *K*^{*}_D: scale invariant ⇒ Quasi-static plasmonic resonances: size independent.
- Analytic formula for the first-order correction to quasi-static plasmonic resonances in terms of the particle's characteristic size δ:



M.A. El-Sayed et al.

• Helmholtz equation:

$$\begin{cases} \nabla \cdot \Big(\varepsilon_m \chi(\mathbb{R}^d \setminus \overline{D}) + \varepsilon_c(\omega) \chi(\overline{D}) \Big) \nabla u + \omega^2 u = 0, \\ u^s := u - u^i \text{ satisfies the outgoing radiation condition.} \end{cases}$$

- u^i : incident plane wave; $k_m := \omega \sqrt{\varepsilon_m}, k_c := \omega \sqrt{\varepsilon_c(\omega)}$.
- Integral formulation on ∂D :

$$\begin{cases} \mathcal{S}_D^{k_c}[\phi] - \mathcal{S}_D^{k_m}[\psi] = u^i, \\ \varepsilon_c (\frac{l}{2} - (\mathcal{K}_D^{k_c})^*)[\phi] - \varepsilon_m (\frac{l}{2} + (\mathcal{K}_D^{k_m})^*)[\psi] = \varepsilon_m \partial u^i / \partial \nu. \end{cases}$$

• Operator-valued function $\delta \mapsto \mathcal{A}_{\delta}(\omega) \in \mathcal{L}(\mathcal{H}^*(\partial B), \mathcal{H}^*(\partial B))$:

$$\mathcal{A}_{\delta}(\omega) = \overbrace{(\lambda(\omega)I - \mathcal{K}_{B}^{*})}^{\mathcal{A}_{0}(\omega)} + (\omega\delta)^{2}\mathcal{A}_{1}(\omega) + O((\omega\delta)^{3}).$$

Quasi-static limit:

$$\mathcal{A}_0(\omega)[\psi] = \sum_{j=0}^{\infty} \tau_j(\omega)(\psi,\varphi_j)_{\mathcal{H}^*}\varphi_j, \quad \tau_j(\omega) := \frac{1}{2} \big(\varepsilon_m + \varepsilon_c(\omega)\big) - \big(\varepsilon_c(\omega) - \varepsilon_m\big)\lambda_j.$$

• Shift in the plasmonic resonances:

$$\arg\min_{\omega} \left| \frac{1}{2} (\varepsilon_m + \varepsilon_c(\omega)) - (\varepsilon_c(\omega) - \varepsilon_m) \lambda_j + (\omega \delta)^2 \tau_{j,1} \right|$$

•
$$\tau_{j,1} := (\mathcal{A}_1(\omega)[\varphi_j], \varphi_j)_{\mathcal{H}^*}$$

Nanoparticle imaging

• Single nanoparticle imaging:

$$\max_{z^{S}} I(z^{S}, \omega)$$

- $I(z^{s}, \omega)$: imaging functional; z^{s} : search point.
- Resolution: limited only by the signal-to-noise-ratio.
- Cross-correlation techniques: robustness with respect to medium noise.





Nanoparticle imaging

• Blow-up of the electric field in the gap at the plasmonic resonances:

$$abla u = O(rac{1}{(\delta/R)^{3/2}\ln(R/\delta)}).$$



Nanoparticle imaging

• Reconstruction from plasmonic spectroscopic data:



Habib Ammari