

Lecture 6: Quantitative thermo-acoustic imaging

Habib Ammari

Department of Mathematics, ETH Zürich

Quantitative thermo-acoustic imaging

- Quantitative thermoacoustic imaging:
 - Reconstruct the **absorption coefficient** from thermal **energy measurements**;
 - Given several data sets \Rightarrow **analytical reconstruction formula**;
 - Formula involves **derivatives** of the given data \Rightarrow **unstable**: small measurement noises may cause large errors;
 - Regularization technique \Rightarrow **good initial guess**;
 - **Optimal control approach**.

Quantitative thermo-acoustic imaging

- u^k : solution of

$$\begin{cases} \Delta u^k + (k^2 + ikq)u^k = 0 & \text{in } \Omega, \\ u^k = g & \text{on } \partial\Omega. \end{cases}$$

- Thermoacoustic imaging problem: inverse problem of reconstructing the absorption coefficient q from thermoacoustic measurements $q|u^k|^2$ in Ω for $k \in (\underline{k}, \bar{k})$.
- $q|u^k|^2$ in Ω : **heat energy** due to the absorption distribution q .
- Heat energy generates an acoustic wave propagating inside the medium. Finding the initial data in the acoustic wave from boundary measurements \rightarrow the heat energy distribution.
- **Quantitative** imaging: **separate q from u^k** .

Quantitative thermo-acoustic imaging

- **Polarization procedure** to obtain more data:
 - $g_1, g_2 \in L^2(\partial\Omega)$. For $j = 1, 2$, and $k \in (\underline{k}, \bar{k})$, u_j^k : solution of

$$\begin{cases} \Delta u_j^k + (k^2 + ikq)u_j^k = 0 & \text{in } \Omega, \\ u_j^k = g_j & \text{on } \partial\Omega. \end{cases}$$

- Function:

$$E_2^k(x) = q(x)u_2^k(x)\overline{u_1^k(x)}, x \in \Omega$$

can be evaluated from the knowledge of the data

$$q|u_1^k + u_2^k|^2 \text{ and } q|iu_1^k + u_2^k|^2 \leftarrow g = g_1 + g_2 \text{ and } g = ig_1 + g_2.$$

- Proof:

$$E_2^k = \frac{1}{2}(q|u_1^k + u_2^k|^2 - q|u_1^k|^2 - q|u_2^k|^2) + \frac{i}{2}(q|iu_1^k + u_2^k|^2 - q|u_1^k|^2 - q|u_2^k|^2).$$

Quantitative thermo-acoustic imaging

- u_j^k : solution corresponding to $(g_j)_{j=1}^{d+1} = (1, x_1, \dots, x_d)$.
- **Polarized data:** $\mathcal{E}^k = (E_j^k)_{j=1}^{d+1} = (qu_j^k \bar{u}_1^k)_{j=1}^{d+1} \quad \forall k \in (\underline{k}, \bar{k})$.
- **Proper set of measurements:**
 - Assume that q : compactly supported in Ω and denote by Ω' its support.
 - \Rightarrow There exist $N > 1$ pairwise disjoint open subsets B_1, B_2, \dots, B_N of Ω , and N frequencies $k_1, \dots, k_N \in (\underline{k}, \bar{k})$ s.t. $\bar{\Omega}' \subset \cup_{j=1}^N \bar{B}_j \subset \Omega$ and, for any $n = 1, \dots, N$,
 - (i) $|u_1^{k_n}| > 0$ in B_n ;
 - (ii) The matrix $[u_j^{k_n}, \nabla u_j^{k_n}]_{1 \leq j \leq d+1}$ is invertible for all $x \in B_n$.

Quantitative thermo-acoustic imaging

- Proof:
 - Fix an arbitrary point $z \in \Omega$, assume that $u_1^k(z) = 0$ for all $k \in (\underline{k}, \bar{k})$.
 - $k \mapsto u_1^k(z)$: analytic $\Rightarrow u_1^k(z) = 0$ for all $k \in \mathbb{R} \Rightarrow u_1^0(z) = 0$.
 - $k = 0$ and $g = 1 \Rightarrow u_1^0(z) = 1$.
 - \Rightarrow For all $z \in \Omega$, there is $k_z \in (\underline{k}, \bar{k})$ s.t. $u_1^{k_z}(z)$ does not vanish.
 - By the continuity of $u_1^{k_z}$, $|u_1^{k_z}| > 0$ in B_z , a small neighborhood of z in Ω .
 - $\bar{\Omega}$: compact \Rightarrow one can extract B_1, \dots, B_N from $\{B_z : z \in \Omega\}$ s.t. $\bar{B}_1, \dots, \bar{B}_N$ cover $\bar{\Omega}$, and hence (i) holds.
 - (ii) can be proved similarly using the differentiability of the determinant.
- $(g_j)_{j=1}^{d+1} \subset L^2(\partial\Omega)$: **proper set of measurements** on (\underline{k}, \bar{k}) iff the corresponding (u_j^k) , $j = 1, \dots, d+1$, $k \in (\underline{k}, \bar{k})$ satisfy (i) and (ii).

Quantitative thermo-acoustic imaging

- **Exact formula:**
- **Proper set of measurements** $(g_j)_{j=1}^{d+1}$.
- Fix $n \in \{1, \dots, N\}$. In B_n , for $2 \leq j \leq d+1$ and $1 \leq l \leq N$,

$$u_j^{k_l} = \alpha_j^{k_l} u_1^{k_l}, \quad \alpha_j^{k_l} = \frac{E_j^{k_l}}{E_1^{k_l}}.$$

- Let $\beta^{k_l} = \Im(\bar{u}_1^{k_l} \nabla u_1^{k_l})$. Then

$$-\nabla \cdot \beta^{k_l} = k_l E_1^{k_l} \quad \text{in } B_n.$$

Quantitative thermo-acoustic imaging

- Proof:
 - $\varphi \in C_c^\infty(B_n, \mathbb{R})$: arbitrary function; Use $\varphi u_1 \in H_0^1(B_n)$ as a test function in

$$-\Delta u_1^{k_l} = (k_l^2 + ik_l q)u_1^{k_l}$$

- \Rightarrow

$$\int_{\Omega} \varphi |\nabla u_1^{k_l}|^2 dx + \int_{\Omega} \overline{u_1^{k_l}} \nabla u_1^{k_l} \cdot \nabla \varphi dx = \int_{\Omega} (k_l^2 + ik_l q) |u_1^{k_l}|^2 \varphi dx.$$

- Taking the imaginary part \Rightarrow

$$-\int_{\Omega} \nabla \cdot (\Im \overline{u_1^{k_l}} \nabla u_1^{k_l}) \varphi dx = \int_{\Omega} k_l q |u_1^{k_l}|^2 \varphi dx = \int_{\Omega} k_l E_1^{k_l} \varphi dx.$$

Quantitative thermo-acoustic imaging

- For all $2 \leq j \leq d + 1$ and $1 \leq l \leq N$,

$$\nabla \alpha_j^{k_l} \cdot \left(\nabla \log \frac{q}{E_1^{k_l}} - \frac{2iq\beta^{k_l}}{E_1^{k_l}} \right) = \Delta \alpha_j^{k_l} \quad \text{in } B_n.$$

- Proof: Fix $j \in \{2, \dots, d + 1\}$. $u_j^{k_l}$: solution of the Helmholtz equation \Rightarrow

$$\begin{aligned} (k_l^2 + ik_l q) \alpha_j^{k_l} u_1^{k_l} &= -\Delta (\alpha_j^{k_l} u_1^{k_l}) \\ &= (k_l^2 + ik_l q) \alpha_j^{k_l} u_1^{k_l} - u_1^{k_l} \Delta \alpha_j^{k_l} - 2 \nabla u_1^{k_l} \cdot \nabla \alpha_j^{k_l}. \end{aligned}$$

- \Rightarrow

$$\begin{aligned} -E_1^{k_l} \Delta \alpha_j^{k_l} &= 2q \overline{u_1^{k_l}} \nabla u_1^{k_l} \cdot \nabla \alpha_j^{k_l} \\ &= q \left(\nabla |u_1^{k_l}|^2 + 2i \Im \overline{u_1^{k_l}} \nabla u_1^{k_l} \right) \cdot \nabla \alpha_j^{k_l}. \end{aligned}$$

- $\Rightarrow q \nabla |u_1^{k_l}|^2 \cdot \nabla \alpha_j^{k_l} = -E_1^{k_l} \Delta \alpha_j^{k_l} - 2iq\beta^{k_l} \cdot \nabla \alpha_j^{k_l}$.

- Differentiating $E_1^{k_l} = q |u_1^{k_l}|^2 \Rightarrow$

$$\nabla E_1^{k_l} = q \nabla |u_1^{k_l}|^2 + E_1^{k_l} \nabla \log q.$$

- $\Rightarrow (\nabla E_1^{k_l} - E_1^{k_l} \nabla \log q) \cdot \nabla \alpha_j^{k_l} = -E_1^{k_l} \Delta \alpha_j^{k_l} - 2iq\beta^{k_l} \cdot \nabla \alpha_j^{k_l}$.

Quantitative thermo-acoustic imaging

- The set $\{\nabla\alpha_j^{k_l}\}_{2\leq j\leq d+1}$: **linearly independent** for all $x \in \bar{\Omega}$.
- $d = 2$:

$$\begin{aligned}\det \begin{bmatrix} \nabla\alpha_2^{k_l} \\ \nabla\alpha_3^{k_l} \end{bmatrix} &= \frac{1}{(u_1^{k_l})^4} \det \begin{bmatrix} u_1^{k_l} \nabla u_2^{k_l} - u_2^{k_l} \nabla u_1^{k_l} \\ u_1^{k_l} \nabla u_3^{k_l} - u_3^{k_l} \nabla u_1^{k_l} \end{bmatrix} \\ &= \frac{1}{(u_1^{k_l})^4} \left(\det \begin{bmatrix} u_1^{k_l} \nabla u_2^{k_l} \\ u_1^{k_l} \nabla u_3^{k_l} - u_3^{k_l} \nabla u_1^{k_l} \end{bmatrix} \right. \\ &\quad \left. - u_2^{k_l} \det \begin{bmatrix} \nabla u_1^{k_l} \\ u_1^{k_l} \nabla u_3^{k_l} - u_3^{k_l} \nabla u_1^{k_l} \end{bmatrix} \right) \\ &= \frac{1}{(u_1^{k_l})^3} \left(u_1^{k_l} \det \begin{bmatrix} \nabla u_2^{k_l} \\ \nabla u_3^{k_l} \end{bmatrix} + u_3^{k_l} \det \begin{bmatrix} \nabla u_1^{k_l} \\ \nabla u_2^{k_l} \end{bmatrix} \right. \\ &\quad \left. - u_2^{k_l} \det \begin{bmatrix} \nabla u_1^{k_l} \\ \nabla u_3^{k_l} \end{bmatrix} \right) \\ &= \frac{1}{(u_1^{k_l})^3} \det \begin{bmatrix} u_1^{k_l} & \nabla u_1^{k_l} \\ u_2^{k_l} & \nabla u_2^{k_l} \\ u_3^{k_l} & \nabla u_3^{k_l} \end{bmatrix} \neq 0.\end{aligned}$$

Quantitative thermo-acoustic imaging

- $A^{k_l} = [\nabla \alpha_{j+1}^{k_l}]_{1 \leq j \leq d}$: invertible in $B_n \Rightarrow$

$$\nabla \log \frac{q}{E_1^{k_l}} - \frac{2iq\beta^{k_l}}{E_1^{k_l}} = \mathbf{a}^{k_l},$$

\mathbf{a}^{k_l} : vector $(A^{k_l*} A^{k_l})^{-1} [A^{k_l*} (\nabla \cdot A^{k_l})]$.

- Evaluate q : split \Re and \Im parts \Rightarrow

$$\nabla \log \frac{q}{E_1^{k_l}} = \frac{\nabla q}{q} - \nabla \log E_1^{k_l} = \Re(\mathbf{a}^{k_l}), \quad \beta^{k_l} = -\frac{E_1^{k_l} \Im(\mathbf{a}^{k_l})}{2q}.$$

- Differentiation \Rightarrow

$$\nabla \cdot \beta^{k_l} = \frac{E_1^{k_l} \Im(\mathbf{a}^{k_l}) \cdot \nabla q}{2q^2} - \frac{\nabla \cdot (E_1^{k_l} \Im(\mathbf{a}^{k_l}))}{2q}.$$

- \Rightarrow

$$\begin{aligned} q &= -\frac{E_1^{k_l} (\Re(\mathbf{a}^{k_l}) + \nabla \log E_1^{k_l}) \cdot \Im(\mathbf{a}^{k_l}) - \nabla \cdot (E_1^{k_l} \Im(\mathbf{a}^{k_l}))}{2k_l E_1^{k_l}} \\ &= -\frac{\Re(\mathbf{a}^{k_l}) \cdot \Im(\mathbf{a}^{k_l}) - \nabla \cdot \Im(\mathbf{a}^{k_l})}{2k_l}. \end{aligned}$$

Quantitative thermo-acoustic imaging

- A^k :

$$A^k = [\nabla \alpha_{j+1}^k]_{j=1}^d,$$

$$\alpha_j^k = \frac{E_j^k}{E_1^k}$$

and the data $(E_j)_{j=1}^{d+1}$: given by **proper set of measurements** $\{1, x_1, \dots, x_d\}$.

- **Exact formula:**

$$q(x) = \frac{-\Re(a^k) \cdot \Im(a^k) + \nabla \cdot \Im(a^k)}{2k};$$

- $a^k = ((A^k)^* A^k)^{-1} [(A^k)^* \nabla \cdot A^k]$.

Quantitative thermo-acoustic imaging

- Optimal control approach:
 - Minimization of the discrepancy functional:

$$J[q] = \frac{1}{2} \int_{\underline{k}}^{\bar{k}} \int_{\Omega} |q|u^k|^2 - E^k|^2 dx dk$$

- Initial guess:

$$q_I = \frac{1}{\bar{k} - \underline{k}} \int_{\underline{k}}^{\bar{k}} \frac{-\Re(a^k) \cdot \Im(a^k) + \nabla \cdot \Im(a^k)}{2k} dk.$$

Quantitative thermo-acoustic imaging

- Differentiability of the data map and its inverse:
 - Let $0 < \underline{q} < \bar{q}$. Let $L_+^\infty(\Omega)$: open set in $L^\infty(\Omega)$:

$$L_+^\infty(\Omega) = \left\{ p \in L^\infty(\Omega) : \underline{q} < p < \bar{q} \text{ in } \Omega \right\}.$$

- Fix $k \in (\underline{k}, \bar{k})$. Define the solution and the data map as

$$\begin{aligned} u^k : L_+^\infty(\Omega) &\rightarrow H^1(\Omega) \\ q &\mapsto u^k[q] \end{aligned}$$

$$\begin{aligned} F^k : L_+^\infty(\Omega) &\rightarrow L^2(\Omega) \\ q &\mapsto F^k[q] = q|u^k[q]|^2, \end{aligned}$$

$u^k[q]$: solution of

$$\begin{cases} \Delta u^k + (k^2 + ikq)u^k = 0 & \text{in } \Omega, \\ u^k = 1 & \text{on } \partial\Omega. \end{cases}$$

- F^k : **well-defined** $\leftarrow u^k \in C^1(\bar{\Omega})$.

Quantitative thermo-acoustic imaging

- u^k : **Fréchet differentiable** in $L_+^\infty(\Omega)$. Derivative at q given by

$$du^k[q](\rho) = v^k(\rho), \quad \forall \rho \in B_q,$$

$B_q \subset L_+^\infty(\Omega)$: open neighborhood of 0 (that depends on q) in $L^\infty(\Omega)$ and $v^k(\rho)$: solution of

$$\begin{cases} \Delta v^k + (k^2 + ikq)v^k = -ik\rho u^k[q] & \text{in } \Omega, \\ v^k = 0 & \text{on } \partial\Omega. \end{cases}$$

- **Differentiability of F^k :**

(i) F^k : **Fréchet differentiable**;

$$dF^k[q](\rho) = \rho |u^k[q]|^2 + 2q \Re(u^k[q] \bar{v}^k(\rho)), \quad \forall q \in L_+^\infty(\Omega), \rho \in B_q.$$

(ii) Dual of dF^k , dF^{k*} : given by

$$(\rho, dF^{k*}[q](h)) = \Re \int_{\Omega} (q |u^k[q]|^2 h + iku^k[q] \bar{p}^k(h)) \rho dx,$$

Quantitative thermo-acoustic imaging

- $p^k(h)$: solution of the **dual problem**

$$\begin{cases} \Delta p^k + (k^2 + ikq)p^k = 2qu^k[q]h & \text{in } \Omega, \\ p^k = 0 & \text{on } \partial\Omega. \end{cases}$$

- (iii) There exists $c > 0$ s.t. for all $\rho \in L^2(\Omega)$

$$\int_{\underline{k}}^{\bar{k}} \|dF^k[q](\rho)\|_{L^2(\Omega)} dk \geq c\|\rho\|_{L^2(\Omega)}.$$

Quantitative thermo-acoustic imaging

- Proof: It is sufficient to show that

$$\lim_{\|\rho\|_{L^\infty(\Omega)} \rightarrow 0} h(\rho) = 0; \quad h(\rho) = \frac{\|u^k[q + \rho] - u^k[q] - v^k(\rho)\|_{L^2(\Omega)}}{\|\rho\|_{L^\infty(\Omega)}}.$$

- $u^k[q + \rho] - u^k[q] - v^k(\rho)$ solution to

$$\begin{cases} (\Delta + k^2 + ikq)(u^k[q + \rho] - u^k[q] - v^k(\rho)) = -ik\rho(u^k[q + \rho] - u^k[q]) & \text{in } \Omega, \\ u^k[q + \rho] - u^k[q] - v^k(\rho) = 0 & \text{on } \partial\Omega. \end{cases}$$

- \Rightarrow

$$\|u^k[q + \rho] - u^k[q] - v^k(\rho)\|_{L^2(\Omega)} \leq \frac{\|\rho\|_{L^\infty(\Omega)} \|(u^k[q + \rho] - u^k[q])\|_{L^2(\Omega)}}{\inf q}.$$

- $u^k[q + \rho] - u^k[q]$ satisfies

$$\begin{cases} \Delta(u^k[q + \rho] - u^k[q]) + (k^2 + ik(q + \rho))(u^k[q + \rho] - u^k[q]) & = -ik\rho u^k[q] \\ u^k[q + \rho] - u^k[q] & = 0 \end{cases}$$

Quantitative thermo-acoustic imaging

- \Rightarrow

$$\|u^k[q + \rho] - u^k[q]\|_{L^2(\Omega)} \leq \frac{\|\rho\|_{L^\infty(\Omega)} \|u^k[q]\|_{L^2(\Omega)}}{\inf(q + \rho)}.$$

- Since $dF^k[q] = |u^k[q]|^2(I + \text{compact})$, it is sufficient to prove that

$$\bigcap_{k \in (\underline{k}, \bar{k})} \text{Ker} dF^k[q] = \{0\}$$

and then apply the Fredholm alternative.

- Assume there exists $\rho \in L^2(\Omega) \setminus \{0\}$ s.t.

$$dF^k[q](\rho) = 0$$

for all $k \in (\underline{k}, \bar{k})$. By analyticity, it follows that $dF^0[q](\rho) = 0$.

Contradiction: $v^0(\rho) = 0$ and $u^0[q] = 1$.

- Regularity theory $\Rightarrow u^k[q] \in L^\infty(\Omega) \Rightarrow dF^k[q]$ can be extended so that its domain is $L^2(\Omega)$.

Quantitative thermo-acoustic imaging

- J : **Fréchet differentiable** in q . For all $q \in L_+^\infty(\Omega)$,

$$dJ[q](\rho) = \int_{\underline{k}}^{\bar{k}} \int_{\Omega} \rho \left(|u^k[q]|^2 (q|u^k|^2 - E^k) + \Re ikp^k \bar{u}^k[q] \right) dx dk;$$

p^k : solution of the **dual problem** with $h = (q|u^k|^2 - E^k)$.

- **Gradient descent method** to minimize J :

$$q^{(n+1)} = Tq^{(n)} - \eta dJ[Tq^{(n)}],$$

$\eta > 0$: step size,

$$Tf = \max\{\underline{q}, \min\{\bar{q}, f\}\}.$$

- **Landweber scheme**:

$$q^{(n+1)} = Tq^{(n)} - \eta \int_{\underline{\omega}}^{\bar{\omega}} \left[dF^*[Tq^{(n)}](F^k(Tq^{(n)}) - E^k) \right] dk.$$

- Landweber scheme or **equivalently** the optimal control approach **converges** to q_* .