

We shall now put the theory of Sobolev spaces $W^{1,p}(I)$, $I \subset \mathbb{R}$, to use to study simple 2nd order ODEs on I using the 2 approaches outlined earlier.

Solutions via the Riesz representation theorem ("APPROACH 2").

Let $I = (a, b) \subset \mathbb{R}$, $-\infty < a < b < \infty$.

Let $f \in C^0(\bar{I})$ and consider

$$\begin{cases} u'' = f & \text{in } I \\ u(a) = u(b) = 0. \end{cases} \quad \otimes$$

Using Riesz, $\exists!$ $u \in H_0^1(I)$ solving \otimes weakly, in that

$$\int_I u' v' dx = - \int_I f v dx \quad \forall v \in H_0^1(I). \quad \oplus$$

Proposition (P. 8) The weak solution u satisfies $u \in C^2(I)$ and is a classical solution of \otimes .

Proof. By \oplus , $u' \in L^2(I)$ has weak derivative

$$\begin{aligned} (u')' &= f \in C^0(\bar{I}) \subset L^2(I) && \text{continuous} \\ \Rightarrow u' \in W^{1,2}(I) &\Rightarrow u'(x) = u'(x_0) + \int_{x_0}^x \widetilde{f(t)} dt \text{ lies} \\ \text{in } C^1(\bar{I}) &\Rightarrow u \in C^2(\bar{I}) \text{ since } u(x) = u(x_0) + \int_{x_0}^x u'(t) dt. \end{aligned}$$

• Boundary conditions: since $C_c^\infty(I) \subset H_0^1(I)$ is dense, and $H_0^1(I) \subset C^0(I)$ by Theorem (T.20), every $u \in H_0^1(I)$ vanishes at a, b (since this is true for $u \in C_c^\infty(I)$). □

Solutions via self-adjointness arguments ("APPROACH 1").

Let again $I = (a, b)$, $-\infty < a < b < \infty$.

We define $A = \frac{d^2}{dx^2}$ on $\mathcal{D}(A) = \{u \in C^2(\bar{I}) : u(a) = u(b) = 0\}$.

We aim to prove Theorem (T.15) (\bar{A} is self-adjoint). First, we identify the domain of \bar{A} . We need:

Lemma (L.11) $H_0^1(I) = \{u \in H^1(I) : u(a) = u(b) = 0\}$.

(This was our original definition of $H_0^1(I)$, so really this lemma is about showing that the original definition is consistent with Definition (D.13).

Proof " \subseteq " Easy.

" \supseteq ". Given $u \in H^1(I)$, $u(a) = u(b) = 0$, let $v(x) = \begin{cases} u(x), & x \in I \\ 0, & x \notin I. \end{cases}$

For $\varphi \in C_c^\infty(\mathbb{R})$,

$$\begin{aligned} \int \varphi' v \, dx &= \int_I \varphi' u \, dx \\ &= \int_a^b (\varphi u)' - \varphi u' \, dx \\ &= \underbrace{(\varphi(b)u(b) - \varphi(a)u(a))}_{=0} - \int_a^b \varphi u' \, dx; \end{aligned}$$

so v has a weak derivative $v' \in L^2(\mathbb{R})$ defined by

$$v'(x) = \begin{cases} u'(x), & x \in I \\ 0, & x \notin I. \end{cases} \Rightarrow v \in H^1(\mathbb{R}), \text{ supp } v \subset \bar{I}.$$

• WLOG, $I = (-1, 1)$. Let $v_k(x) = v\left(\frac{x}{1 - \frac{1}{k}}\right)$, then

$$v_k \in H^1(\mathbb{R}), \quad v_k \xrightarrow{k \rightarrow \infty} v \text{ in } H^1(\mathbb{R}), \quad \text{supp } v_k \subset \left(-1 + \frac{1}{k}, 1 - \frac{1}{k}\right).$$

We can then approximate v_k in $H^1(\mathbb{R})$ by

$$v_{k,\ell} := \rho_{\ell^{-1}} * v_k \in C_c^\infty\left(\left(-1 + \frac{1}{k} - \frac{1}{\ell}, -1 + \frac{1}{k} + \frac{1}{\ell}\right)\right) \subset C_c^\infty(\mathbb{I}) \quad (\ell \geq k)$$

as in Corollary (C.3). So v is the limit of a sequence

$$v_{k,\ell_k} \in C_c^\infty(\mathbb{I}) \text{ in } H^1(\mathbb{R}). \text{ This proves } u = v|_{\mathbb{I}} \in H_0^1(\mathbb{I}), \square$$

Lemma (L.12) $D(\bar{A}) = H^2(\mathbb{I}) \cap H_0^1(\mathbb{I})$.

Proof " \supseteq " Given $u \in H^2(\mathbb{I}) \cap H_0^1(\mathbb{I})$, pick (using Theorem (T.19')(iii))

$u_k \in C^\infty(\bar{\mathbb{I}})$ with $u_k \rightarrow u$ in $H^2(\mathbb{I})$. Since $H^2(\mathbb{I}) \subset C^0(\bar{\mathbb{I}})$,

Lemma (L.11) gives $u_k(a), u_k(b) \xrightarrow{k \rightarrow \infty} 0$, so

$$v_k(x) := u_k(a) + \frac{x-a}{b-a} (u_k(b) - u_k(a)) \xrightarrow{k \rightarrow \infty} 0 \text{ in } H^2(\mathbb{I})$$

$$\Rightarrow \tilde{u}_k := u_k - v_k \rightarrow u \text{ in } H^2(\mathbb{I}),$$

$$\tilde{u}_k(a) = \tilde{u}_k(b) = 0,$$

so $\tilde{u}_k \in D(\bar{A})$. Since $\tilde{u}_k \rightarrow u$ in $L^2(\mathbb{I})$,

$$\tilde{u}_k'' \rightarrow u'' \text{ in } L^2(\mathbb{I}),$$

we get $u \in D(\bar{A})$.

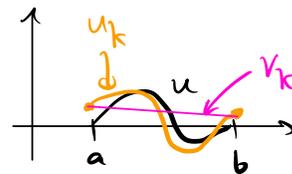
" \subseteq " Let $u \in D(\bar{A})$, so $\exists \{u_k\} \subset D(\bar{A}) \stackrel{\text{Lemma (L.11)}}{=} C^2(\bar{\mathbb{I}}) \cap H_0^1(\mathbb{I})$

$$\text{s.t. } \begin{cases} u = \lim_{k \rightarrow \infty} u_k \text{ in } L^2(\mathbb{I}), \\ \end{cases}$$

$$\begin{cases} \exists \lim_{k \rightarrow \infty} Au_k = \lim_{k \rightarrow \infty} u_k'' =: v \text{ in } L^2(\mathbb{I}). \end{cases}$$

To conclude, we only need to show that u_k' converges in $L^2(\mathbb{I})$.

But this follows from



$$\begin{aligned}
\|u'_k - u'_l\|_{L^2(I)}^2 &= \int_I |u'_k - u'_l|^2 dx \\
&= - \int_I (u''_k - u''_l)(u_k - u_l) dx \\
&\leq \|u''_k - u''_l\|_{L^2(I)} \|u_k - u_l\|_{L^2(I)} \\
&\xrightarrow{k, l \rightarrow \infty} 0.
\end{aligned}$$

□

Proof of Theorem (T.15) in $n=1$ dimension: \bar{A} is self-adjoint.

If $v \in L^2(I)$ satisfies $v \in D(\bar{A}^*)$, i.e. $\exists C$ s.t.

$$|(v, \bar{A}u)_{L^2(I)}| \leq C \|u\|_{L^2(I)} \quad \forall u \in D(\bar{A}), \quad \textcircled{\#}$$

then we must show that $v \in D(\bar{A})$.

Well, $\textcircled{\#}$ implies the existence of $g \in L^2(I)$ s.t.

$$(v, \bar{A}u) = (g, u) \quad \forall u \in D(\bar{A}).$$

In particular, $(v, \varphi'') = (g, \varphi) \quad \forall \varphi \in C_c^\infty(I)$.

Let $G(x) = \int_a^x g(t) dt$, so $G \in H^1(I)$ and

$$(v, \varphi'') = (G', \varphi) = -(G, \varphi') \quad \forall \varphi \in C_c^\infty(I). \quad \textcircled{\otimes}$$

Formally, $(v, \varphi'') = -(v', \varphi')$, which suggests that $v' = G - c$, $c \in \mathbb{R}$,

(similarly to Lemma (L.9)).

Claim: $v \in H^2(I)$. To prove this, let $\chi \in C_c^\infty(I)$, $\int \chi dx = 1$; given

$\psi \in C_c^\infty(I)$, we then have $\varphi(x) := \int_a^x (\psi(t) - \chi(t) \int_I \psi ds) dt \in C_c^\infty(I)$,

$$\varphi' = \psi - \chi \int_I \psi dt,$$

$$\varphi'' = \psi' - \chi' \int_I \psi dt,$$

and therefore from $\textcircled{\otimes}$

$$\begin{aligned}
0 &= (v, \psi'') + (G, \psi') \\
&= \int_{\mathbb{I}} v(x) \left(\psi'(x) - \chi'(x) \int_{\mathbb{I}} \psi dt \right) + G(x) \left(\psi(x) - \chi(x) \int_{\mathbb{I}} \psi dt \right) dx \\
&= \int_{\mathbb{I}} v \psi' + G \psi - \psi \underbrace{\left(\int v \chi' + G \chi dt \right)}_{=: c \in \mathbb{R}} dx \\
&= \int_{\mathbb{I}} v \psi' + (G - c) \psi dx.
\end{aligned}$$

Since $\psi \in C_0^\infty(\mathbb{I})$ is arbitrary, this implies that $v \in H^1(\mathbb{I})$ with $v' = G - c \in H^1(\mathbb{I})$, so in fact even $v \in H^2(\mathbb{I})$ with $v'' = G' = g$; so $\bar{A}^* v = g = v''$.

Claim: $v(a) = v(b) = 0$. Indeed, for $u \in \mathcal{D}(A)$ we have

$$\begin{aligned}
0 &= \int_{\mathbb{I}} v u'' - g u dx = v u' \Big|_a^b - \int_{\mathbb{I}} v' u' + g u dx \\
&= (v(b) u'(b) - v(a) u'(a)) - \underbrace{v' u \Big|_a^b}_{=0} + \underbrace{\int (v'' - g) u dx}_{=0} \\
&= v(b) u'(b) - v(a) u'(a).
\end{aligned}$$

Plug in $u \in \mathcal{D}(A)$ with $u'(a) = 1, u'(b) = 0 \Rightarrow v(a) = 0$;
 $u'(a) = 0, u'(b) = 1 \Rightarrow v(b) = 0$.

By Lemma (L.11), this gives $v \in H_0^1(\mathbb{I})$, finishing the proof. \square