

We return to Corollary (C.10) and ask: what can one say about the eigenvalues  $0 < \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$  of  $-\Delta$  as an operator on  $L^2(\Omega)$  with domain  $H^2(\Omega) \cap H_0^1(\Omega)$ , where  $\Omega \subset \mathbb{R}^n$  is bounded and of class  $C^\infty$ ?

Theorem (T.35) (Weyl's Law) Let  $\Omega \subset \mathbb{R}^n$  be  $C^\infty$ , let  $0 < \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$  denote the Dirichlet eigenvalues of  $-\Delta$  on  $\Omega$ , and let

$$N(T) = \# \{ i : \lambda_i \leq T \}$$

denote the *eigenvalue counting function*. Then

$$\lim_{T \rightarrow \infty} \frac{N(T)}{T^{n/2}} = \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(\Omega),$$

where  $\omega_n = \mathcal{L}^n(B_1(0))$ ,  $\mathcal{L}^n = n$ -dimensional Lebesgue measure.

It is important that we can work in less regular domains.

Lemma (L.19) Let  $\Omega \subset \mathbb{R}^n$  be bounded. Let

$$\sigma(-\Delta) = \{ \lambda \in \mathbb{C} : \exists u \in H_0^1(\Omega) \text{ s.t. } -\Delta u = \lambda u \}.$$

(i)  $\sigma(-\Delta) = \{ \lambda_1, \lambda_2, \dots \}$  where  $0 < \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$ , and  $\exists$  complete ONB  $\{u_1, u_2, \dots\} \subset L^2(\Omega)$  with  $u_j \in H_0^1(\Omega)$ ,  $-\Delta u_j = \lambda_j u_j$ .

(ii) We have the min-max characterization

$$\lambda_k = \inf_{\substack{V \subset H_0^1(\Omega) \\ \dim V = k}} \sup_{\substack{u \in V \\ u \neq 0}} \frac{\|\nabla u\|_{L^2(\Omega)}^2}{\|u\|_{L^2(\Omega)}^2}.$$

This largely generalizes Corollary (C.10), except it is not true anymore that the eigenfunctions  $u_k$  lie in  $C^\infty(\bar{\Omega})$  (or even in  $H^2(\Omega)$ ).

Proof Step 1. We define  $A: H_0^1(\Omega) \rightarrow H^{-1}(\Omega) := (H_0^1(\Omega))^*$  by

$$(Au)(v) := (\nabla u, \nabla v)_{L^2(\Omega)}.$$

To understand what this is, note that there is a continuous embedding  $L^2(\Omega) \hookrightarrow H^{-1}(\Omega)$  given by

$$L^2(\Omega) \ni f \mapsto (H_0^1(\Omega) \ni v \mapsto (f, v)_{L^2(\Omega)});$$

and then  $Au = f$  — an equation in  $H^{-1}(\Omega)$  — means

$$(\nabla u, \nabla v)_{L^2(\Omega)} = (f, v)_{L^2(\Omega)} \quad \forall v \in H_0^1(\Omega),$$

i.e.  $-\Delta u = f$  weakly.

1.1.  $A$  is an isomorphism.

Indeed, since  $(u, v)_{H_0^1} = (Au)(v)$ , this is precisely the content of the Riesz representation theorem for  $H_0^1(\Omega)$ .

1.2. We compose  $A^{-1}: H^{-1}(\Omega) \rightarrow H_0^1(\Omega)$  with embeddings from/into  $L^2(\Omega)$  and get the operator

$$\begin{array}{ccc}
 T: L^2(\Omega) & \longrightarrow & L^2(\Omega) \\
 \downarrow & & \uparrow \\
 H^{-1}(\Omega) & \xrightarrow{A^{-1}} & H_0^1(\Omega)
 \end{array}$$

The inclusion map  $H_0^1(\Omega) \rightarrow L^2(\Omega)$  is compact (*exercise* — this does not require any regularity on  $\partial\Omega$ , unlike Corollary (C.8)). Therefore,  $T$  is a compact operator.

One can also easily check (*exercise*) that  $T$  is symmetric. Finally,  $T$  is positive: if  $f \in L^2(\Omega)$ , then

$$\begin{aligned}
 (Tf, f)_{L^2(\Omega)} &= (u, f)_{L^2(\Omega)} && \text{(where } u := Tf \in H_0^1(\Omega) \\
 &= (\nabla u, \nabla u)_{L^2(\Omega)} && \text{solves } -\Delta u = f) \\
 &\geq 0,
 \end{aligned}$$

with equality iff  $u=0$  in  $H_0^1(\Omega)$  and thus  $f=0$ .

Step 2. (i) follows from the spectral theorem for compact self-adjoint operators, with the  $\lambda_j =$  reciprocals of the eigenvalues of  $T$  (which are all  $>0$ ).

(ii) follows either from a similar characterization of the eigenvalues of  $T$  (see FA I), or directly using the complete ONB of eigenfunctions of  $-\Delta$  (*exercise*).  $\square$

Example (E.21)  $\Omega = (0, L)^n$ ,  $L > 0$ .

(i) Claim: for  $k \in \mathbb{N}^n$ , let

$$u_k(x) = \left(\frac{2}{L}\right)^{n/2} \prod_{j=1}^n \sin\left(\frac{\pi}{L} k_j x_j\right).$$

Then  $u_k \in H_0^1(\Omega)$ ,  $-\Delta u_k = \left(\frac{\pi}{L}\right)^2 |k|^2 u_k$ , and

$\{u_k\}_{k \in \mathbb{N}^n}$  is a complete ONB of  $L^2(\Omega)$ .

Indeed, only the final claim is not clear. Extend  $u_k$  to an odd function  $\tilde{u}_k$  on  $\tilde{\Omega} := (-L, L)^n$ , i.e.

$$\tilde{u}_k(x_1, \dots, -x_i, \dots, x_n) = -\tilde{u}_k(x_1, \dots, x_i, \dots, x_n) \quad \forall i$$

— that is,  $\tilde{u}_k(x) = \left(\frac{2}{L}\right)^{n/2} \prod_{j=1}^n \sin\left(\frac{\pi}{L} k_j x_j\right)$  simply.

• Every  $L^2$ -function on  $\tilde{\Omega}$  can be expanded into Fourier series,

$$v = \sum_{k \in \mathbb{Z}^n} v_k \tilde{e}_k, \quad (v_k)_{k \in \mathbb{Z}^n} \in \ell^2, \quad v_k = \int_{\tilde{\Omega}} v(x) \overline{\tilde{e}_k(x)} dx,$$

$$\tilde{e}_k(x) = (2L)^{-n/2} e^{i \frac{\pi}{L} x \cdot k}.$$

If now  $u \in L^2(\Omega)$ , let  $v \in L^2(\tilde{\Omega})$  be its odd extension to  $\tilde{\Omega}$ . Let  $k = (k_1, \dots, k_n)$ .

— If  $k_i = 0$  for some  $i$ , then  $\tilde{e}_k(\dots, x_i, \dots) = \tilde{e}_k(\dots, -x_i, \dots)$

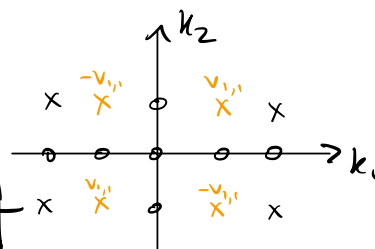
is even in  $x_i$ , but  $v$  is odd in  $x_i$ , so  $v_k = 0$ .

— If  $k_1, \dots, k_n \neq 0$ , and  $k^\# = (\pm k_1, \dots, \pm k_n)$ , then

$$v_{k^\#} = (-1)^{\text{number of } - \text{ signs}} v_k$$

These  $2^n$  terms contribute to

the Fourier series of  $v$  a total of



$$\begin{aligned}
\sum_{k^\# = (\pm k_1, \dots, \pm k_n)} v_{k^\#} \tilde{e}_{k^\#}(x) &= v_k \sum_{k^\# = (\pm k_1, \dots, \pm k_n)} (-1)^{\#(-\text{sign})} \tilde{e}_{k^\#}(x) \\
&= v_k \sum_{k' = (\pm k_1, \dots, \pm k_{n-1})} (-1)^{\#(-\text{sign})} \left( \tilde{e}_{(k', k_n)}(x) - \tilde{e}_{(k', -k_n)}(x) \right) \\
x = (x', x_n) \downarrow \\
&= v_k \sum_{k' = (\pm k_1, \dots, \pm k_{n-1})} (-1)^{\#(-\text{sign})} 2i (2L)^{-\frac{n}{2}} e^{\frac{i\pi}{L} x' \cdot k'} \sin\left(\frac{\pi}{L} x_n k_n\right) \\
&= \dots = (2i)^n (2L)^{-\frac{n}{2}} v_k \prod_{j=1}^n \sin\left(\frac{\pi}{L} x_j k_j\right) \\
&= v_k i^n \tilde{u}_k(x).
\end{aligned}$$

Therefore,  $v = \sum_{k \in \mathbb{N}^n} i^n v_k \tilde{u}_k$  in  $L^2(\tilde{\Omega})$ ,

and thus also  $u = \sum_{k \in \mathbb{N}^n} i^n v_k u_k$  in  $L^2(\Omega)$ , as required.

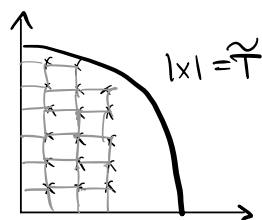
(ii) We conclude that the set of Dirichlet eigenvalues of  $-\Delta$  on  $\Omega = (0, L)^n$  is

$$S = \left\{ \left(\frac{\pi}{L}\right)^2 |k|^2 : k \in \mathbb{N}^n \right\}.$$

$$\begin{aligned}
\left(\frac{\pi}{L}\right)^2 |k|^2 &\leq T \\
\Leftrightarrow |k| &\leq \frac{L}{\pi} \sqrt{T}
\end{aligned}$$

(iii) Claim. Weyl's law holds on  $\Omega = (0, L)^n$ .

Indeed,  $N(T) = \# \left( \mathbb{N}^n \cap \underbrace{\overline{B_{\tilde{T}}(0)}}_{= \{x \in \mathbb{R}^n : |x| \leq \tilde{T}\}} \right)$ ,  $\tilde{T} = \frac{L}{\pi} T^{\frac{1}{2}}$



This is the Lebesgue measure of the union of the gray boxes, so

$$\begin{aligned}
N(T) &\leq \mathcal{L}^n([0, \infty)^n \cap \overline{B_{\tilde{T}}(0)}) \\
&= 2^{-n} \mathcal{L}^n(\overline{B_{\tilde{T}}(0)}) \\
&= 2^{-n} \tilde{T}^n \omega_n \quad (\omega_n = \mathcal{L}^n(\overline{B_1(0)}))
\end{aligned}$$

$$= \left(\frac{L}{2\pi}\right)^n T^{n/2} \omega_n$$

On the other hand,  $[0, \infty)^n \times \overline{B_{\tilde{T}-\sqrt{n}}(0)} \subset \text{union of gray boxes,}$

$$\text{so } N(T) \geq 2^{-n} (\tilde{T} - \sqrt{n})^n \omega_n = \underbrace{\left(\frac{L}{2\pi}\right)^n T^{n/2} \omega_n}_{\text{from } \tilde{T}} \underbrace{\left(1 - \frac{\pi\sqrt{n}}{L\sqrt{T}}\right)^n}_{\left(\frac{\tilde{T}-\sqrt{n}}{\tilde{T}}\right)^n}$$

$$\Rightarrow \frac{N(T)}{T^{n/2}} = \frac{\omega_n}{(2\pi)^n} Z^n(\Omega) \cdot (1 + O(T^{-1/2}))$$

$$\xrightarrow{T \rightarrow \infty} \frac{\omega_n}{(2\pi)^n} Z^n(\Omega).$$

We next have the following monotonicity result for eigenvalues:

Lemma (L.20) Let  $\lambda_k(\Omega)$  denote the  $k^{\text{th}}$  eigenvalue ( $k \geq 1$ , counting with multiplicity) of  $-\Delta$  on  $\Omega \subset \mathbb{R}^n$ .

If  $\Omega \subset \Omega'$ , then  $\lambda_k(\Omega') \leq \lambda_k(\Omega) \quad \forall k \in \mathbb{N}$ .

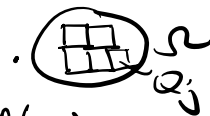
Proof This follows directly from Lemma (L.19) (ii): if  $V \subset H_0^1(\Omega)$ ,  $\dim V = k$ , then extending the elements of  $V$  by 0 to  $\Omega'$  gives a subspace  $V \subset H_0^1(\Omega')$ . So  $\lambda_k(\Omega)$  is an infimum over a smaller set (of linear subspaces of  $H_0^1(\Omega')$  of dimension  $k$ ) than  $\lambda_k(\Omega')$ , so  $\lambda_k(\Omega) \geq \lambda_k(\Omega')$ .  $\square$

We can now already prove one half of Weyl's law.

Proof that  $\liminf_{T \rightarrow \infty} \frac{N(T)}{T^{n/2}} \geq \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(\Omega).$

Given  $\varepsilon > 0$ , we can choose  $L > 0$  so small that there exists a disjoint union  $\bigcup_{j=1}^J Q_j \subset \Omega$  of cubes  $Q_j = q_j + Q$ ,  $Q = (0, L)^n$ , s.t.

$$\mathcal{L}^n(\Omega) - \varepsilon < \mathcal{L}^n\left(\bigcup_{j=1}^J Q_j\right) = \sum \mathcal{L}^n(Q) < \mathcal{L}^n(\Omega).$$



Write  $\lambda_1^L \leq \lambda_2^L \leq \dots$  for the Dirichlet eigenvalues of  $Q_j$  (see Example (E.21)), and  
 $\lambda_1^0 \leq \lambda_2^0 \leq \dots$  for the sequence  $\lambda_1^L, \dots, \lambda_1^L$  ( $J$  times),  
 $\lambda_2^L, \dots, \lambda_2^L$  ( $J$  times),  
 $\dots$

We claim that  $\lambda_k^0 \geq \lambda_k(\Omega)$  for all  $k \in \mathbb{N}$ .

Indeed, by Lemma (L.19),

$$\begin{aligned} \lambda_k^0 &= \inf_{\substack{V \subset H_0^1\left(\bigcup_{j=1}^J Q_j\right) \\ \dim V = k}} \sup_{\substack{u \in V \\ u \neq 0}} \frac{\|\nabla u\|^2}{\|u\|^2} \\ &\geq \inf_{\substack{V \subset H_0^1(\Omega) \\ \dim V = k}} \sup_{\substack{u \in V \\ u \neq 0}} \frac{\|\nabla u\|^2}{\|u\|^2} = \lambda_k(\Omega). \end{aligned}$$

$$\begin{aligned} \Rightarrow N(T) &= \#\{k: \lambda_k(\Omega) \leq T\} \\ &\geq \#\{k: \lambda_k^0 \leq T\} \\ &= J \cdot \#\{k: \lambda_k^L \leq T\} \\ &= J \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(Q) T^{\frac{n}{2}} (1 + o(1)) \end{aligned}$$

$$> \frac{\omega_n}{(2\pi)^n} (Z^n(\Omega) - \varepsilon) T^{\frac{n}{2}} (1 + o(1)), \quad T \rightarrow \infty.$$

$$\Rightarrow \liminf_{T \rightarrow \infty} \frac{N(T)}{T^{\frac{n}{2}}} \geq \frac{\omega_n}{(2\pi)^n} (Z^n(\Omega) - \varepsilon).$$

Since  $\varepsilon > 0$  is arbitrary, we are done.  $\square$

The idea here was to **use only** test functions  $u$  in the min-max characterization of  $\lambda_k(\Omega)$  which are  $H^1_0$ -functions on each cube  $Q_j$  to get an **upper bound on  $\lambda_k(\Omega)$** . That this gives a sharp asymptotic lower bound for  $N(T)$  (equivalently, a sharp asymptotic upper bound on  $\lambda_k(\Omega)$  — **exercise**) is not obvious; we shall prove this (i.e. the other inequality required for a proof of Weyl's law) by getting a **lower bound on  $\lambda_k(\Omega)$**  by using a **larger space** of test functions than  $H^1_0(\Omega)$ .

We have few choices apart from considering (e.g. for  $\Omega = (0, 1)^n$ )

$$\nu_k := \inf_{\substack{V \subset H^1(\Omega) \\ \dim V = k}} \sup_{\substack{u \in V \\ u \neq 0}} \frac{\|\nabla u\|_{L^2(\Omega)}^2}{\|u\|_{L^2(\Omega)}^2}.$$

What is **this**? We have the following analogue of Lemma (L.19):

Lemma (L.21) Let  $\Omega \in \mathbb{R}^n$ . Assume  $H^1(\Omega) \rightarrow L^2(\Omega)$  is compact.



$$\sigma_p(-\Delta) = \{ \lambda \in \mathbb{C} : \exists u \in H^1(\Omega) \text{ s.t.}$$

$$(\nabla u, \nabla v)_{L^2(\Omega)} = \lambda (u, v)_{L^2(\Omega)} \forall v \in H^1(\Omega) \}.$$

(i)  $\sigma_p(-\Delta) = \{ \lambda_1, \lambda_2, \dots \}$  where  $0 < \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$ ,  
and  $\exists$  complete ONB  $\{u_1, u_2, \dots\} \subset L^2(\Omega)$  with  $u_j \in H^1(\Omega)$ ,  
 $(\nabla u_j, \nabla v) = \lambda_j (u_j, v) \quad \forall v \in H^1(\Omega)$ .

(ii) We have the min-max characterization  $\lambda_k = \mu_k$ .

Proof. We set  $B: H^1(\Omega) \rightarrow (H^1(\Omega))^*$ ,

$$(Bu)(v) = (u, v)_{L^2} + (\nabla u, \nabla v)_{L^2} = (u, v)_{H^1(\Omega)}.$$

(So  $Bu = -\Delta u + u$  when  $u \in H^2(\Omega)$ .) This is an isomorphism. By assumption,

$$T: L^2(\Omega) \hookrightarrow (H^1(\Omega))^* \xrightarrow{B^{-1}} H^1(\Omega) \hookrightarrow L^2(\Omega)$$

is compact. Its symmetry is a simple calculation: let  $f, g \in L^2(\Omega)$ , and set  $u = Tf, v = Tg \in H^1(\Omega)$ ; then  $Bv = g$  (regarded as an element of  $(H^1(\Omega))^*$  via  $H^1(\Omega) \ni u \mapsto (g, u)_{L^2(\Omega)}$ ), and therefore

$$\begin{aligned} (Tf, g)_{L^2(\Omega)} &= (u, g)_{L^2(\Omega)} = (Bv)(u) = (u, v)_{H^1(\Omega)} \\ &= (v, u)_{H^1(\Omega)} = \dots = (f, Tg)_{L^2(\Omega)} \end{aligned}$$

(working with real-valued functions here).

•  $T$  is positive:  $(Tf, f) = (u, u)_{H^1(\Omega)} > 0$  unless  $u=0$ .

• Get  $0 < \mu_1 \leq \mu_2 \leq \dots \rightarrow \infty$ , complete ONB  $\{u_k\}_{k \in \mathbb{N}} \subset L^2(\Omega)$ ,

$$Tu_k = \mu_k u_k \Leftrightarrow u_k = \mu_k Bu_k \Leftrightarrow Bu_k = \mu_k^{-1} u_k$$

$$\Leftrightarrow (\nabla u_k, \nabla v)_{L^2(\Omega)} = \left(\frac{1}{\mu_k} - 1\right) (u_k, v)_{L^2(\Omega)} \quad \forall v \in H^1(\Omega).$$

$$\Rightarrow \lambda_k = \frac{1}{\mu_k} - 1.$$

To show that  $\lambda_k \geq 0$ , note that  $(\nabla u, \nabla v) = \lambda(u, v) \quad \forall v \in H^1(\Omega)$   
implies (for  $v = u$ )  $\|\nabla u\|^2 = \lambda \|u\|^2$ , and hence  $\lambda \geq 0$  indeed.  $\square$

Remark (R.29) (i) For  $\Omega \in \mathbb{R}^n$ ,  $H^1(\Omega) \hookrightarrow L^2(\Omega)$  may fail to be compact. Example:  $\Omega = \bigcup_{k=1}^{\infty} B_{\frac{1}{k^3}}(\frac{1}{k}, 0, \dots, 0)$ . (Disjoint union.)

$$\left( \begin{array}{c} \text{...} \textcircled{1} \textcircled{2} \\ u_2 \end{array} \right) \text{ let } u_j = \begin{cases} j^{3n/2} & \text{on } B_{\frac{1}{j^3}}(\frac{1}{j}, 0, \dots, 0), \\ 0 & \text{otherwise.} \end{cases}$$

Then all derivatives of  $u_j$  vanish, so

$$\|u_j\|_{H^1(\Omega)}^2 = \|u_j\|_{L^2(\Omega)}^2 = \int_{B_{\frac{1}{j^3}}} (j^{3n/2})^2 dx = \omega_n$$

is bounded, yet  $\|u_j - u_l\|_{L^2(\Omega)}^2 = 2\omega_n$  for  $j \neq l$ .

So  $\{u_j\}$  does not have a convergent subsequence in  $L^2(\Omega)$ .

(ii) For  $\Omega = (0, L)^n$ , the inclusion  $H^1(\Omega) \hookrightarrow L^2(\Omega)$  is compact.

(Exercise.) For  $\Omega$  of class  $C^1$ , compactness of the inclusion follows from Corollary (C.8).

Now, what are the Neumann eigenvalues  $\nu_1, \nu_2, \dots$  of  $-\Delta$  on  $\Omega \in \mathbb{R}^n$  as in Lemma (L.21)? If  $u \in H^1(\Omega)$ ,  $-\Delta u = f$  (in the sense that  $(\nabla u, \nabla v) = (f, v) \quad \forall v \in H^1(\Omega)$ ), suppose that  $\Omega$  is  $C^\infty$ , and  $u, f \in C^\infty(\bar{\Omega})$ . Then for  $v \in C^\infty(\bar{\Omega})$

$$\int_{\Omega} \nabla u \cdot \nabla v \, dx = \int_{\partial\Omega} \frac{\partial u}{\partial \nu} \cdot v \, d\sigma - \int_{\Omega} \Delta u \cdot v \, dx$$

$$\stackrel{!}{=} \int_{\Omega} f \cdot v \, dx,$$

so  $\int_{\Omega} (\Delta u + f) \cdot v \, dx = \int_{\partial\Omega} \frac{\partial u}{\partial \nu} \cdot v \, d\sigma.$

Plugging in  $v \in C_c^\infty(\Omega)$ , we conclude that  $-\Delta u = f$  in  $\Omega$ .

Thus  $\int_{\partial\Omega} \frac{\partial u}{\partial \nu} v \, d\sigma = 0 \quad \forall v \in C^\infty(\bar{\Omega}) \Rightarrow \frac{\partial u}{\partial \nu}|_{\partial\Omega} = 0.$

Therefore, we are in fact working with a weak formulation of the **Neumann problem**

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial\Omega. \end{cases}$$

Remark (R.30) One can develop a regularity theory analogous to Theorem (T.32) also for this problem.

In particular, the  $u_k$  from Lemma (L.21) are weak solutions of the eigenvalue problem

$$\begin{cases} -\Delta u_k = \lambda_k u_k \\ \frac{\partial u_k}{\partial \nu} = 0 & \text{on } \partial\Omega. \end{cases}$$

Example (E.22) Consider  $\Omega = (0, L)^n$ . Then  $H^1(\Omega) \hookrightarrow L^2(\Omega)$  is compact (**exercise**).

(i) Claim: for  $k \in \mathbb{N}_0^n$ , let

$$u_k(x) = \left(\frac{2}{L}\right)^{\frac{n}{2}} \prod_{j=1}^n \cos\left(\frac{\pi}{L} k_j x_j\right) \in H^1(\Omega).$$

Then  $u_k \in H^1(\Omega)$  solves  $-\Delta u_k = \left(\frac{\pi}{L}\right)^2 |k|^2 u_k$  weakly, and  $\{u_k\}_{k \in \mathbb{N}_0^n}$  is a complete ONB of  $L^2(\Omega)$ .

Indeed, note that  $\frac{\partial u_k}{\partial x_i} = 0$  at  $x_i = 0, L$  for  $i=1, \dots, n$ ;

this implies via an integration by parts that  $u_k$  is a weak solution of the **Neumann eigenvalue equation** indeed.

The completeness of  $\{u_k\}$  follows by a variation of the argument used in Example (E.21).

(ii) We conclude that the set of **Neumann eigenvalues** of  $-\Delta$  on  $\Omega = (0, L)^n$  is

$$S = \left\{ \left(\frac{\pi}{L}\right)^2 |k|^2 : k \in \mathbb{N}_0^n \right\}.$$

(iii) Claim. Weyl's law holds on  $\Omega = (0, L)^n$  with Neumann boundary conditions; i.e.

$$\lim_{T \rightarrow \infty} \frac{\#\{v \in S : |v| \leq T\}}{T^{n/2}} = \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(\Omega).$$

Proof via a minor modification of the argument in Example (E.21).

We can now already prove the other half of Weyl's law.

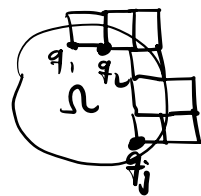
Proof that  $\limsup_{T \rightarrow \infty} \frac{N(T)}{T^{n/2}} \leq \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(\Omega).$

( $\Rightarrow$  Theorem (T.35).)

• Given  $\varepsilon > 0$ , we choose  $L > 0$  so small that there exist distinct points  $q_1, \dots, q_J \in L \cdot \mathbb{Z}^n$  s.t.  $\Omega \subset \bigcup_{j=1}^J (\bar{Q} + q_j)$  where  $\bar{Q} = [0, L]^n$ ,

and  $Z^n(\Omega) \leq J \cdot Z^n(\bar{\Omega}) < Z^n(\Omega) + \varepsilon$ .

- Write  $0 \leq \nu_1 \leq \nu_2 \leq \dots \rightarrow \infty$  for the Neumann eigenvalues of  $Q$ , and  $0 \leq \lambda_1^N \leq \lambda_2^N \leq \dots \rightarrow \infty$



for the sequence  $\underbrace{\nu_1, \dots, \nu_1}_{J \text{ times}}, \underbrace{\nu_2, \dots, \nu_2}_{J \text{ times}}, \dots$ .

- We claim that  $\lambda_k(\Omega) \geq \lambda_k^N$  for all  $k \in \mathbb{N}$ .

Indeed, by Lemma (L.19),

$$\lambda_k(\Omega) = \inf_{\substack{V \subset H_1^0(\Omega) \\ \dim V = k}} \sup_{\substack{u \in V \\ u \neq 0}} \frac{\|\nabla u\|_{L^2(\Omega)}^2}{\|u\|_{L^2(\Omega)}^2}$$

if  $u \in H_1^0(\Omega)$ , then  $u|_{Q+q_j} \in H^1(Q+q_j)$   $\Rightarrow$

$$\geq \inf_{\substack{V \subset L^2(\mathbb{R}^n), \dim V = k \\ \forall j: V|_{Q+q_j} \in H^1(Q+q_j)}} \sup_{\substack{u \in V \\ \|u\|_{L^2(\Omega)} = 1}} \|\nabla u\|_{L^2(\Omega)}^2$$

$$= \lambda_k^N$$

The final equality follows by writing  $u \in V$  as

$$u(x) = \sum_{j=1}^J 1_{Q_j} \sum_{\substack{c_{j\ell} \\ \ell \in \mathbb{N}}} u_\ell(x - q_j), \quad \sum_{j,\ell} |c_{j\ell}|^2 = 1,$$

with

$$\|\nabla u\|_{L^2(\Omega)}^2 = \sum_{j=1}^J \sum_{\ell} |c_{j\ell}|^2 \nu_\ell$$

$\nwarrow$  ONB of  $H^1(Q)$  from Example (E.22)

and following the arguments in the proof of the min-max formula from here.

$$\Rightarrow N(T) = \# \{k: \lambda_k(\Omega) \leq T\}$$

$$\leq \# \{k: \lambda_k^N \leq T\}$$

$$= J \cdot \# \{k: \nu_k \leq T\}$$

$$= \int \frac{\omega_n}{(2\pi)^n} \mathcal{L}^n(Q) T^{\frac{n}{2}} (1 + o(1))$$

$$< \frac{\omega_n}{(2\pi)^n} (\mathcal{L}^n(\Omega) + \varepsilon) T^{\frac{n}{2}} (1 + o(1)), \quad T \rightarrow \infty.$$

$$\Rightarrow \limsup_{T \rightarrow \infty} \frac{N(T)}{T^{\frac{n}{2}}} \leq \frac{\omega_n}{(2\pi)^n} (\mathcal{L}^n(\Omega) - \varepsilon).$$

Since  $\varepsilon > 0$  is arbitrary, we are done. □