# Talk Diableretes

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### 2 Vertex algebra of Joyce and the conformal element

To explain the Lie bracket that appeared in Miguel's talk, I will describe how it is derived from the Vertex algebra constructed by Joyce(17'). For this, we will begin with the moduli stack of complexes on X

$$\mathcal{M}_{X}$$

and consider its shifted homology for each component labeled by a class  $\alpha$ :

$$\ddot{H}_*(\mathcal{M}_\alpha) = H_{*-\chi(\alpha,\alpha)}(\mathcal{M}_X,\mathbb{C}),$$

where  $\chi$  is the Euler pairing<sup>\*</sup>. The vertex algebra on

$$V_* = \bigoplus_{\alpha} \hat{H}_*(\mathcal{M}_{\alpha}) \tag{2.1}$$

consists mainly of two important ingredients:

• the translation operator

 $T: V_* \longrightarrow V_{*+2}$ 

• and the state-to-field correspondence

$$Y(-,z): V_* \longrightarrow \operatorname{End}(V_*) \llbracket z, z^{-1} \rrbracket$$

satisfying additional compatibilities. One of them being the condition of locality

$$(z-w)^N \Big[ Y(v,z), Y(u,z) \Big] = 0$$

for any two  $u, v \in V$ . These are constructed out of their respective ingredients

<sup>\*</sup>We will always work over  $\mathbb C$  for a reason that will become apparent later.

• the action of  $\rho : [*/\mathbb{G}_m] \times \mathcal{M}_X \to \mathcal{M}_X$  which in S-families of complexes on X is given by rescalling by a line bundle L on S pulled back from S:

$$(\mathcal{E} \to X \times S) \mapsto (\mathcal{E} \otimes \pi_S^*(L) \to X \times S).$$

• The morphism induced by mapping sheaves to their direct sums

$$\sigma:\mathcal{M}_X\times\mathcal{M}_X\longrightarrow\mathcal{M}_X$$

and the complex

$$\mathcal{E}\mathrm{xt}\longrightarrow \mathcal{M}_X\times \mathcal{M}_X$$

given by  $\operatorname{Ext}(E, F)$  at each point  $([E], [F]) \in \mathcal{M}_X \times \mathcal{M}_X$  together with the translation operator T itself.

The construction of both of these are then given by

1. letting t be the dual of  $c_1(\mathcal{L})$  for the universal line bundle on  $[*/\mathbb{G}_m]$  and setting

$$T(u) = \phi_*(t \boxtimes u) \,.$$

2. For  $u \in \hat{H}_d(\mathcal{M}_\alpha)$  and  $v \in \hat{H}_e(\mathcal{M}_\beta)$ , we construct

$$Y(u,z)v = (-1)^{\chi(\alpha,\beta)} z^{\chi(\alpha,\beta) + \chi(\beta)} \sigma_* \left( e^{zT} \otimes \operatorname{id} \left( u \boxtimes v \cap c_{z^{-1}}(\mathcal{E}\mathrm{xt}^{\vee} + \sigma^* \mathcal{E}\mathrm{xt}) \right), \quad (2.2)$$

It was the observation of Joyce(17') that these satisfy the necessary axioms of a vertex algebra.

**Example 2.1.** We will mostly focus here on the case of a curve C. Cheating a little bit this will also include acyclic quivers Q and complexes of representations, because both are homological dimension 1. The construction above still makes sense in this case. The homology (2.1) has now an explicit description in terms of  $K^0 := K^0(C)$  (resp.  $K^0(Q)$ ). Fixing a basis B of  $K^0 \otimes \mathbb{C}$  constructed as

$$B = \sqcup_{p,q} B^{p,q},$$

where  $B^{p,q}$  is a basis of  $H^{p,q}(C,\mathbb{C})$ , we define the classes

$$\mu_{v,i} = \operatorname{ch}_i(\mathbb{G}/v^{\vee}).$$

where  $\mathbb{F} \to X \times \mathcal{M}_X$  is the universal sheaf.

To compare with Miguel's notation one sees by a simple computation that

$$\operatorname{ch}_{i}(w) = \sum_{v \in B} \int_{X} w \cdot v \mu_{v,i-[p,q]}$$

where  $[p,q] := \left\lfloor \frac{p-q}{2} \right\rfloor$ . We may write

$$H^*(\mathcal{M}_C) = \mathbb{C}[K^0(C)] \otimes \operatorname{SSym}^{\bullet}_{\mathbb{C}}\llbracket \mu_{v,k}, v \in B, k > 0 \rrbracket, H_*(\mathcal{M}_C) = \mathbb{C}[K^0(C)] \otimes \operatorname{SSym}^{\bullet}_{\mathbb{C}}\llbracket u_{v,k}, v \in B, k > 0 \rrbracket,$$

where  $\mu_{v,k} = \frac{1}{(k-1)!} \frac{\partial}{\partial u_{v,k}}$ . The latter has a natural vertex algebra structure called the *lattice* vertex algebra associated to the pairing

$$\chi_{\text{ssym}}: K^0 \times K^0 \longrightarrow \mathbb{Z}, \qquad \chi_{\text{ssym}}(v, w) = \chi(v, w) + (-1)^{\deg(v)} \chi(w, v).$$

and by Gross(19')/Joyce(17') it is the one described by (2.2).

Miguel constructed in his talk a map  $\xi^{\mathbb{G}} : \mathbb{D}^C \to H^*(\mathcal{M}_C)$  using the universal sheaf  $\mathbb{G}$  on  $C \times \mathcal{M}_C$  such that after imposing the relations

$$\operatorname{ch}_{i}(\gamma) = \begin{cases} 0 & \text{if } p - q \ge 2i \text{ and } p \neq q, \\ \int_{Y} \operatorname{ch}(\alpha) \cdot \gamma & \text{if } p = q, i = 0. \end{cases}$$
(2.3)

it becomes an isomorphisms only because we are working with curves. This has to do with  $B^{0,q}$  being bounded by q = 1.

Continuing to use the notation of the example, recall that Miguel defined the operators

 $T_k^*, R_k^*$ 

on the algebra  $\mathbb{D}^C$  in terms of  $ch_i(v)$  and their derivations. The condition that these can be formulated on the full  $H^*(\mathcal{M}_C)$  is equivalent to

$$\xi^{\mathbb{G}}\left[R_k^*\left(\operatorname{ch}_i(v^{p,q})\right)\right] = 0\,,$$

whenever the degree of  $\xi^{\mathbb{G}}(\operatorname{ch}_i(v^{p,q}))$  is zero. This holds again only for curves C, because in this case degree of  $\operatorname{ch}_i(v^{1,0}) = 0$ . We then denote by

 $T_k, R_k$ 

their duals on the homology  $H_*(\mathcal{M}_C)$ . The beauty of this is that we have a natural conformal element

$$\omega = \sum_{v \in B} e^0 \otimes u_{v,1} u_{v^{\vee},1} \,,$$

<sup>†</sup> such that its field

$$Y(\omega, z) = \sum_{k \in \mathbb{Z}} L_k z^{-k-2}$$
(2.4)

recovers  $L_k = T_k + R_k$  whenever  $k \ge -1$  and

$$B = \bigsqcup B_{p,p} \, .$$

I will discuss the necessary modification when there are non (p, p)-classes later.

## **3** Compatibility explained on an Example: Flag and Gr(n,k)

We show the power of rephrasing Virasoro constraints in terms of the vertex algebra proving some new results without having to do many computations.

<sup>&</sup>lt;sup>†</sup>Here I used the notation  $\{v^{\vee}\}$  to denote the dual basis of B with respect to  $\chi_{ssym}$  pretending that it is non-degenerate. This is not true in general, but can be easily fixed.

Gross–Joyce–Tanaka(20') formulated the wall-crossing to take place in the quotient

$$W_* = V_{*+2}/DV_*$$

which carries a Lie algebra structure by

$$[\bar{u},\bar{v}] = \overline{[z^{-1}]Y(u,z)v}$$

for any lifts  $u, v \in V_*$  of  $\bar{u}, \bar{v} \in W_*$ .

They defined for each reasonable choice of stability condition  $\sigma$  and  $\alpha \in K^0$  the homology classes

$$\left[M_{\alpha}(\sigma)\right]^{\mathrm{in}}$$

satisfying the wall-crossing formulae

$$\left[M_{\alpha}(\sigma_{2})\right]^{\mathrm{in}} = \sum_{\vec{\alpha}\vdash\alpha} \mathrm{coeff.}\left[\cdots \left[\left[M_{\alpha_{1}}(\sigma_{1})\right]^{\mathrm{in}}, \left[M_{\alpha_{1}}(\sigma_{1})\right]^{\mathrm{in}}\right], \cdots, \left[M_{\alpha_{k}}(\sigma_{1})\right]^{\mathrm{in}}\right]\right]$$

for some pre-described coefficients depending on the two different stability condition  $\sigma_1, \sigma_2$ . Their main contribution is in showing that when there are no strictly  $\sigma$ -semistables in class  $\alpha$ , then these are the pushforwards of the usual virtual fundamental classes

$$\left[M_{\alpha}(\sigma)\right]^{\mathrm{vir}}$$

under the open embeddings  $M_{\alpha}(\sigma) \hookrightarrow \mathcal{M}_C$ . Recall the map

$$\eta: \mathbb{D}_C \longrightarrow \mathbb{D}_C^{\mathrm{in}}$$
,

from Miguel's talk which was used to define an invariant Virasoro<sup>‡</sup>. Since the dual of  $R_{-1}$  is our translation operator T, the dual  $\lambda = \eta^*$  induces a well-defined lift

$$\lambda: W_* \longrightarrow V_*$$

**Standard result:** If  $L_k$  are given by the Virasoro element as in (2.4), then if

$$L_k \circ \lambda [M_{\alpha_1(\sigma)}]^{\mathrm{in}} = \delta_{k,0} \lambda [M_\alpha(\sigma)]^{\mathrm{in}}$$

is satisfied for all  $i = 1, \dots, k$ , then

$$L_k \circ \lambda [M_\alpha(\sigma)]^{\rm in} = \delta_{k,0} \lambda [M_\alpha(\sigma)]^{\rm in}.$$
(3.1)

The point is to show that the set  $P \subset L$  of *physical states* satisfying precisely the condition (3.1) forms a Lie subalgebra of L.

Unfortunately in this raw form this can not be applied in too many interesting geometries. That is why I selected these examples. Start with any acyclic quiver Q, for example just  $A^k$ 

$$\stackrel{1}{\circ} \longrightarrow \stackrel{2}{\circ} \longrightarrow \stackrel{3}{\circ} \longrightarrow \cdots \longrightarrow \stackrel{k}{\circ}, \tag{3.2}$$

and consider the  $\mu$ -stability, a representation  $V_{\bullet}$  is said to be (semi)-stable if for every  $W_{\bullet} \subset V_{\bullet}$  and the respective dimension vectors  $\vec{e}$  and  $\vec{d}$  we have

$$\sum_{i} e_i \mu_i < (\leq) \sum_{i} d_i \mu_i \, .$$

We only need two cases

<sup>&</sup>lt;sup>‡</sup>Here  $\mathbb{D}_C^{\text{in}} = \ker(R_{-1})$ 

1. In the case when  $\mu_{i+1} > \mu_i$ , where *i* increases in the direction of the arrows we get that the only non-zero classes

$$\left[M_{\vec{d}},(\mu)\right]^{\mathrm{m}}$$

are those for the dimension vectors  $\delta_i$  when  $i = 1, \dots, k$  and these are just the point classes in  $H_*(\mathcal{M}_{\delta_i})$ .

2. In the example (??), we choose real numbers

$$1 \gg \mu_1 \gg \mu_2 \gg \cdots \gg \mu_{k-1}$$

together with  $\mu_k < -\frac{k(k-1)}{2}$ . It is then easy to see that there are no strictly semistables in the class

$$\vec{d} = (1, 2, \cdots, k)$$

and

$$[M_{\vec{d}}(\mu)]^{\mathrm{m}} = [\mathrm{Flag}(\mathbb{C}^k)].$$

Similar arguments can be applied to the case

$$\stackrel{1}{\circ} \longrightarrow \stackrel{2}{\circ}$$

to get  $[\operatorname{Gr}(n,k)] = [M_{(k,n)}(\mu)]^{\operatorname{in}}$ . This proves the claim.

**Remark 3.1.** The statement for flag-varieties was also proved by Miguel purely using geometry.

### 4 Rank reduction for curves

The first difference that appears when moving on to curves is the issue of extending  $L_k$  for  $k \ge -1$  from Miguel's talk to a full field

$$\sum_{k\in\mathbb{Z}}L_kz^{-k-2}$$

The next few lines are not strictly necessary in the setting of working with curves, but they are a precursor for things to come when thinking of higher dimensions. To shorten notation it makes sense to introduce the *holomorphic pairing* 

$$\chi^{\mathrm{H}}(v,w) = (-1)^{p(v)} \int_{X} v \cdot w \cdot \mathrm{td}(X) \,,$$

 $\S$  and its supersymmetrization

$$\chi_{\text{ssym}}^{\text{H}}(v,w) = \chi^{H}(v,w) + (-1)^{\text{deg}(v)}\chi^{\text{H}}(w,v)$$

The usual field of the lattice vertex algebra we mentioned before is

$$Y(u_{v,1},z) = \left\{ \sum_{k>0} u_{v,k} \cdot z^{k-1} + (-1)^{\delta_{\text{odd}}(v)} \sum_{k>0} \sum_{w \in B} \frac{k!}{(k-\delta_{\text{odd}}(v))!} \chi_{\text{ssym}}(v,w) \frac{\partial}{\partial u_{w,k}} z^{-k-\delta_{\text{ev}}(v)} + \chi(v,\beta) z^{-1} \right\}$$

<sup>&</sup>lt;sup>§</sup>Here H stands for Hodge, because it interacts non-trivially with the Hodge grading.

where we use  $\delta_{\text{odd}}(v) \equiv \deg(v)$  (mod 2). Again to make things easier on the eye, we define its modification

$$Y^{\mathrm{H}}(v,z) = \sum_{k>0} u_{v,k} z^{k} - \sum_{k>0,w\in B} \frac{k!}{(k-\delta_{\mathrm{odd}}(v))!} \chi^{H}_{\mathrm{ssym}}(v,w) \frac{\partial}{\partial u_{w,k}} z^{-k-\delta_{\mathrm{ev}}(v)} + \chi(v,\beta) z^{-1}.$$

To obtain the full field, we may now take the sums of normal ordering products

$$L(z) = \sum_{\substack{\in B^{p,q}\\0 \le q \le p \le 1}} : Y^H(v,z) \frac{\partial}{\partial z} Y^H(\hat{v},z) : + \sum_{\substack{\in B^{0,1}\\0 \le q \le p \le 1}} : \frac{\partial}{\partial z} Y^H(v,z) Y^H(\hat{v},z) :$$

by a direct computation.  $\P$  This is again very special for curves and is important for the next result.

#### Lemma 4.1.

$$L_k \circ \lambda \big[ M_{\alpha_1}(\sigma) \big]^{\text{in}} = \delta_{k,0} \lambda \big[ M_\alpha(\sigma) \big]^{\text{in}}$$
(4.1)

is satisfied for all  $i = 1, \cdots, k$ , then

$$L_k \circ \lambda [M_\alpha(\sigma)]^{\mathrm{in}} = \delta_{k,0} \lambda [M_\alpha(\sigma)]^{\mathrm{in}}.$$

*Proof.* The proof is to show again that there is a Lie-subalgebra of physical states. This relies this time on the fact that the *locality condition* of the fields is satisfied:

$$(z-w)^N [L(z), Y([M_\alpha(\sigma)]^{\text{in}}, w)] = 0.$$

This follow, because the field L(z) could be constructed from the usual pairing in the same way instead, which can be observed by showing that the part of the  $T_k$  operator coming from  $v^{0,1}$ , resp.  $v^{1,0}$  vanishes in both scenarios and the rest is unchanged. We then use Don's lemma.

**Lemma 4.2.** If three fields a(z), b(z) and c(z) are mutually local, then

$$: a^{(n)}(z)b^{(m)}(z):$$

is local with respect to c(z).

It is now a simple observation that if we define the operator

$$\mathcal{L} = \sum_{j \ge -1} \frac{(-1)^j}{(j+1)!} D^j L_j : W_* \longrightarrow V_*$$

which is the dual of the one introduced by Miguel, then (4.1) is equivalent to

$$\mathcal{L}[M_{\alpha(\sigma)}]^{\mathrm{in}} = 0.$$

Furthermore, by Miguel's talk this makes our (3.1) equivalent to the geometric Virasoro constraints observed by van Bree/Moreira/MOOP. We will now use the rank-reduction of

<sup>&</sup>lt;sup>¶</sup>This time  $\hat{v}$  is the dual of v with respect to  $\chi^H$ .

Joyce/Thaddeus/Mochizuki together with the pushforward formula of Miguel from his talk to prove that

$$\mathcal{L}[M_{r,\chi}]^{\mathrm{in}} = 0$$

and more! For this we also need to consider the moduli schemes of Bradlow pairs

$$\mathcal{O}_X(-k) \longrightarrow F$$

which we denote by  $P_{r,\chi}^k$  and their classes sit in the homology of a stack of pairs  $\mathcal{P}_C^k$ .

Theorem 4.3. We have

$$\mathcal{L}^{\mathcal{O}_X(-k)} \left[ P_{r,\chi}^k \right]^{\text{in}} = 0 , \quad \mathcal{L} \left[ M_{r,\chi} \right]^{\text{in}} = 0 .$$

$$(4.2)$$

for any choice of  $(r, \chi)$ .

*Proof.* The result follows by starting with

$$P_{0,n} = C^{[n]}, \qquad M_{1,d} = \operatorname{Jac}^d$$

We then use first that there is a wall-crossing

$$[P_{r,\chi}^{k}]^{\text{in}} = \sum_{\substack{(\vec{r},\vec{\chi}) \vdash (r,\chi) \\ \vec{d} \vdash 1 \\ r > 0}} \text{coeff.} \left[ \left[ [P_{r_{1},\chi_{1}}^{k}]^{\text{in}}, [P_{r_{2},\chi_{2}}^{k}]^{\text{in}} \right], \cdots, [P_{r_{k},\chi_{k}}^{k}]^{\text{in}} \right],$$
(4.3)

where additionally  $r_i > 0$  to show that

$$\mathcal{L}^{\mathcal{O}_X(-k)}[P_{r,\chi}^k]$$

holds assuming (4.2) is satisfied for all lower ranks. Then using

$$[M_{r,\chi}]^{\mathrm{in}} = \Omega_{r,\chi} + \sum_{\substack{\vec{r},\vec{\chi} \models (r,\chi) \\ r_i > 0}} \mathrm{coeff.} \left[ \cdots \left[ [M]^{\mathrm{in}}, [M_{\alpha_2}]^{\mathrm{in}} \right], \cdots, [M_{\alpha_k}]^{\mathrm{in}} \right], \tag{4.4}$$

where the first term on the right hand side is defined by

$$\Omega_{r,\chi} = \pi_* \left( [P_{r,\chi}^k]^{\text{in}} \cap e(\mathbb{T}_{\mathcal{P}_C^k/\mathcal{M}_C}) \right)$$
(4.5)

using the projection  $\pi : \mathcal{P}_C^k \to \mathcal{M}_C$ . We therefore only need a stack version of the result discussed by Miguel

**Proposition 4.4.** If  $[P_{r,\chi}^k]^{\text{in}}$  satisfies

$$\mathcal{L}^{\mathcal{O}_X(-k)}[P^k_{r,\chi}]^{tin} = 0$$

then

$$\mathcal{L}\Omega_{r,\chi}^k = 0.$$

This then shows that

$$\mathcal{L}\big[M_{r,\chi}\big]^{\mathrm{in}} = 0$$

"for all ranks.