## Quantum Dynamical Systems, Sec. 4.3-4.4

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## 4.3 Invariant States

**Definition 1.** If  $\tau^t$  is a group of \*-autmorphisms of the C\*-algebra  $\mathfrak{A}$ , a state  $\mu$  on  $\mathfrak{A}$  is callte  $\tau^t$ -invariant if  $\mu \circ \tau^t = \mu$  for all  $t \in \mathbb{R}$ . The set of such invarianct sets is denoted as  $E(\mathfrak{A}, \tau^t) \subset E(\mathfrak{A})$  the set of all states on  $\mathfrak{A}$ .

**Theorem 1.** Let  $\tau^t$  be a group of \*-autmorphisms of the C\*-algebra  $\mathfrak{A}$ . If there exist a state  $\omega$  on  $\mathfrak{A}$  such that the function  $t \mapsto \omega(\tau^t(A))$  is continuous for all  $A \in \mathfrak{A}$  then  $E(\mathfrak{A}, \tau^t)$  is a non-empty, convex and weak-\*-compact subset of  $\mathfrak{A}^*$ . In particular, this holds if  $(\mathfrak{A}, \tau^t)$  is a C\*-dynamical system.

*Proof.* To prove existance we want to construct a net that has a converging subnet with a  $\tau^t$ -invariant limit.

So for all  $A \in \mathfrak{A}$  consider the expression:

$$\omega_T(A) = \frac{1}{T} \int_0^T \omega \circ \tau^s(A) ds.$$
(1)

By assumption, the function  $s \mapsto \omega(\tau^s(A))$  is continuous, thus the integral is well defined and we have  $\omega_T \in E(\mathfrak{A})$  for all T > 0. Since  $E(\mathfrak{A})$  is weak-\*-compact, the net  $(\omega_T)_{T>0}$  has a weak-\*-convergent subnet. The formula

$$\omega_T(\tau^t(A)) = \frac{1}{T} \int_0^T \omega \circ \tau^s(\tau^t(A)) ds$$
<sup>(2)</sup>

$$=\frac{1}{T}\int_{0}^{T}\omega\circ\tau^{s+t}(A)ds\tag{3}$$

$$=\frac{1}{T}\int_{t}^{T+t}\omega\circ\tau^{s'}(A)ds'$$
(4)

$$=\omega_T(A) - \frac{1}{T} \int_0^t \omega \circ \tau^{s'}(A) ds' + \frac{1}{T} \int_T^{T+t} \omega \circ \tau^{s'}(A) ds' \qquad (5)$$

is used for an estimate. Using this we can estimate

$$|\omega_T(\tau^t(A)) - \omega_T(A)| \le 2||A|| \frac{|t|}{T},$$
(6)

from which it follows that the limit of an convergent subnet of  $(\omega_T)T > 0$  is  $\tau^t$ -invariant. Let  $\mu, \omega$  be states in  $E(\mathfrak{A}, \tau)$  and  $\lambda \in (0, 1)$ . Then  $\gamma(A) = \mu(A) + \lambda \omega(A)$  is a state and also  $\tau^t$ -invariant since

$$\gamma(\tau^t(A)) = \mu(\tau^t(A)) + \lambda\omega(\tau^t(A)) = \mu(A) + \lambda\omega(A) = \gamma(A)$$
(7)

and weak-\*-closedness of the set of invariant state is clear.

**Definition 2.** If  $\tau^t$  is a group of \*-autmorphisms of the von Neumann algebra  $\mathfrak{M}$  we denote by  $N(\mathfrak{M}, \tau^t) \equiv E(\mathfrak{M}, \tau^t) \cap N(\mathfrak{M})$  the set of normal  $\tau^t$ -invariant states.

Note that for a  $W^*$ -dynamical system the compactness argument used in this proof breaks down. There is no general existence result for *normal invariant states*. In fact

## 4.4 Quantum Dynamical Systems

**Definition 3.** If  $\mathfrak{C}$  is a  $C^*$ -algebra and  $\tau^t a$  group of \*-automorphisms of  $\mathfrak{C}$  we define

$$\mathcal{E}(\mathfrak{C},\tau^t) \equiv \{\mu \in E(\mathfrak{C},\tau^t) | t \mapsto \mu(A^*\tau^t(A)) \text{ is continuous for all } A \text{ in } \mathfrak{C}\}$$
(8)

If  $\mu \in \mathcal{E}(\mathfrak{C}, \tau)$  we say that  $(\mathfrak{C}, \tau^t, \mu)$  is a quantum dynamical system.

**Example 1.** If  $(\mathfrak{A}, \tau^t)$  is a  $C^*$ -dynamical system then  $\mathcal{E}(\mathfrak{A}, \tau) = E(\mathfrak{A}, \tau)$  and  $(\mathfrak{A}, \tau^t, \mu)$  is a quantum dynamical system for any  $\tau^t$ -invariant state  $\mu$ .

**Example 2.** If  $(\mathfrak{M}, \tau^t)$  is a  $W^*$ -dynamical system then  $N(\mathfrak{M}, \tau) \subset E(\mathfrak{M}, \tau)$ and  $(\mathfrak{M}, \tau^t, \mu)$  is a quantum dynamical system for any  $\tau^t$ -invariant state  $\mu$ .

**Lemma 1.** Let  $(\mathfrak{C}, \tau^t, \mu)$  be a quantum dynamical system and denote the GNS representation of  $\mathfrak{C}$  associated to  $\mu$  by  $(\mathcal{H}_{\mu}, \pi_{\mu}, \Omega_{\mu})$ . Then there exists a unique self-adjoint operator  $L_{\mu}$  on  $\mathcal{H}_{\mu}$  such that

1. 
$$\pi_{\mu}(\tau^{t}(A)) = e^{itL_{\mu}}\pi_{\mu}e^{-itL_{\mu}}$$
 for all  $A \in \mathfrak{C}$  and  $t \in \mathbb{R}$ 

2.  $L_{\mu}\Omega_{\mu} = 0$ 

*Proof.* For a fixed  $t \in \mathbb{R}$  one easily checks that  $(\mathcal{H}_{\mu}, \pi_{\mu} \circ \tau^{t}, \Omega_{\mu})$  is a GNS representation of  $\mathfrak{C}$  associated to  $\mu$ . By unicity of the GNS construction there exists a unitary operator  $U_{\mu}^{t}$  on  $\mathcal{H}_{\mu}$  such that, for any  $A \in \mathfrak{C}$ , one has

$$U^t_\mu \pi_\mu(A)\Omega_\mu = \pi_\mu(\tau^t(A))\Omega_\mu \tag{9}$$

and in particular

$$U^t_\mu \Omega_\mu = \Omega_\mu \tag{10}$$

For  $s,t\in\mathbb{R}$  we have

$$U^{t}_{\mu}U^{s}_{\mu}\pi_{\mu}(A)\Omega_{\mu} = U^{t}_{\mu}\pi_{\mu}(\tau^{s}(A))\Omega_{\mu} = \pi_{\mu}(\tau^{t+s}(A))\Omega_{\mu} = U^{t+s}_{\mu}\pi_{\mu}(A)\Omega_{\mu}, \quad (11)$$

and the cyclic property of  $\Omega_{\mu}$  yields that  $U_{\mu}^{t}$  si a unitary group on  $\mathcal{H}_{\mu}$ . Using an earlier result one can show that  $U_{\mu}^{t}$  is also strongly continuous. By Stone theorem  $U_{\mu}^{t} = e^{itL_{\mu}}$  for some self-adjoint operator  $L_{\mu}$  and property 2 follows from Equation (9). Finally for  $A, B \in \mathfrak{C}$  we get

$$U^{t}_{\mu}\pi_{\mu}(A)\pi_{\mu}(B)\Omega_{\mu} = \pi_{\mu}(\tau^{t}(A))\pi_{\mu}(\tau^{t}(B))\Omega_{\mu} = \pi_{\mu}(\tau^{t}(A))U^{t}_{\mu}\pi_{\mu}(B)\Omega_{\mu}, \quad (12)$$

and property (1) follows from the cyclic property of  $\Omega_{\mu}$ .

**Definition 4.** Given a quantum dynamical system  $(\mathfrak{C}, \tau^t, \mu)$ , we denote by

- $(\mathcal{H}_{\mu}, \pi_{\mu}, \Omega_{\mu})$  its GNS-representation
- $\mathfrak{C}_{\mu} = \pi_{\mu}(\mathfrak{C})''$  the enveloping von Neumann Algebra
- $(\pi_{\mu}, \mathfrak{C}_{\mu}, \mathcal{H}_{\mu}, L_{\mu}, \Omega_{\mu})$  its Normal Form (which exists by the Lemma 1).

**Definition 5.** Two quantum dynamical systems  $(\mathfrak{C}, \tau^t, \mu)$  and  $(\mathfrak{D}, \sigma^t, \nu)$  are isomorphic if there exists a \*-isomorphism  $\phi : \mathfrak{C} \to \mathfrak{D}$  such that  $\phi \circ \tau^t = \sigma^t \circ \phi$  for all  $t \in \mathbb{R}$  and  $\mu = \nu \circ \phi$ .

**Definition 6.** Let  $\omega, \mu$  be states on  $\mathfrak{C}$ .  $\mu$  is called  $\omega$ -normal if  $\mu = \tilde{\mu} \circ \pi_{\omega}$  for some  $\mu \in N(\mathfrak{C}_{\omega})$ . The set of  $\omega$ -normal states on  $\mathfrak{C}$  is denoted by  $N(\mathfrak{C}, \omega)$ .

**Theorem 2.** (Simplified Version of Thm. 2.30 in the Book) Let  $\omega, \mu$  be states on  $\mathfrak{C}$ . Then  $\mu \in N(\mathfrak{C}, \omega)$  if and only if there exists a  $\sigma$ -weakly continuous  $\ast$ morphism  $\pi_{\mu|\omega} : \mathfrak{C}_{\omega} \to \mathfrak{C}_{\mu}$  such that  $\pi_{\mu} = \pi_{\mu|\omega} \circ \pi_{\omega}$ . If this is the case, then there exists a subalgebra  $z_{\mu|\omega}\mathfrak{C}_{\omega} \subseteq \mathfrak{C}_{\omega}$  such that the restriction of  $\pi_{\mu|\omega}$  to  $z_{\mu|\omega}\mathfrak{C}_{\omega}$ is a  $\ast$ -isomorphism.

*Proof.* We only show the first part. Assume such a \*-morphism exists. Write  $\hat{\mu}(A) = \langle \Omega_{\mu}, A\Omega_{\mu} \rangle$  for the extension of  $\mu$  to  $\mathfrak{C}_{\mu}$ . We get

$$\mu = \hat{\mu} \circ \pi_{\mu} = \hat{\mu} \circ \pi_{\mu|\omega} \circ \pi_{\omega}$$

Since  $\tilde{\mu} := \hat{\mu} \circ \pi_{\mu|\omega}$  defines a normal state on  $\mathfrak{C}_{\omega}$  we can conclude that  $\mu \in N(\mathfrak{C}, \omega)$  is  $\omega$ -normal.

For the other direction, assume we have  $\mu = \tilde{\mu} \circ \pi_{\omega}$  for some  $\mu \in N(\mathfrak{C}_{\omega})$ . Let  $(\mathcal{K}, \Phi, \Psi)$  be the GNS-representation of  $\mathfrak{C}_{\omega}$  corresponding to  $\tilde{\mu}$ . Then  $(\mathcal{K}, \Phi \circ \pi_{\omega}, \Psi)$  is a GNS-representation of  $\mathfrak{C}$  corresponding to  $\mu$ . Indeed, we have

$$\mu(A) = \tilde{\mu}(\pi_{\omega}(A)) = \langle \Psi, \Phi(\pi_{\omega}(A))\Psi \rangle$$

for all  $A \in \mathfrak{C}$  which, by density of  $\pi_{\omega}(\mathfrak{C})$  in  $\mathfrak{C}_{\omega}$  proves  $\mu(A) = \langle \Psi, \Phi(A)\Psi \rangle$  for all  $A \in \mathfrak{C}_{\omega}$ . Furthermore,  $\mathcal{K} = \overline{\Phi(\mathfrak{C}_{\omega})\Psi} = \overline{\Phi \circ \pi_{\omega}(\mathfrak{C})\Psi}$  also since  $\pi_{\omega}(\mathfrak{C})$  is dense in  $\mathfrak{C}_{\omega}$ . By the uniqueness of the GNS-representation, there exists a unitary map  $U: \mathcal{K} \to \mathcal{H}_{\mu}$  such that

$$\pi_{\mu}(A) = U\Phi(\pi_{\omega}(A))U^*$$

for all  $A \in \mathfrak{C}$ . Thus, for  $X \in \mathfrak{C}_{\omega}$  we can simply define

$$\pi_{\mu|\omega}(X) := U\Phi(X)U^*$$

and we get  $\pi_{\mu} = \pi_{\mu|\omega} \circ \pi_{\omega}$ .

**Lemma 2.** Let  $\omega \in \mathcal{E}(\mathfrak{C}, \tau^t)$  and  $\mu \in N(\mathfrak{C}, \omega) \cap E(\mathfrak{C}, \tau^t)$ . Then we have  $\mu \in \mathcal{E}(\mathfrak{C}, \tau^t)$  and, the map  $\pi_{\mu|\omega}$  from Thm 2 is an isomorphism between the quantum dynamical systems  $(z_{\mu|\omega}\mathfrak{C}_{\omega}, \hat{\tau}^t_{\omega}, \tilde{\mu})$  and  $(\mathfrak{C}_{\mu}, \hat{\tau}^t_{\mu}, \hat{\mu})$ , where  $\hat{\mu}$  and  $\tilde{\mu}$  are defined the same way as in the previous proof and  $\hat{\tau}^t_{\mu}$  is defined as

$$\hat{\tau}^t_\mu(A) := e^{itL_\mu} A e^{-itL_\mu}$$

for  $A \in \mathfrak{C}_{\mu}$ .  $\hat{\tau}^t_{\omega}$  is defined analogously.

*Proof.* Notice that for all  $A \in \mathfrak{C}$  we have

$$\mu(A^*\tau^t(A)) = \tilde{\mu}(\pi_\omega(A)^*\pi_\omega(\tau^t(A))) = \tilde{\mu}(\pi_\omega(A)^*e^{itL_\omega}\pi_\omega(A)e^{-itL_\omega})$$

which is continuous in t since  $\tilde{\mu}$  is normal.

We still need to show that  $\hat{\mu} \circ \pi_{\mu|\omega} = \tilde{\mu}$  and that  $\hat{\tau}^t_{\mu} \circ \pi_{\mu|\omega} = \pi_{\mu|\omega} \circ \hat{\tau}^t_{\omega}$ . By density, it suffices to show that these equalities holds on  $\pi_{\omega}(\mathfrak{C})$ . We have

$$\hat{\mu} \circ \pi_{\mu|\omega} \circ \pi_{\omega}(A) = \hat{\mu}(\pi_{\mu}(A)) = \mu(A) = \tilde{\mu} \circ \pi_{\omega}(A)$$

where we have used (in this order) Thm 2, a property of the GNS representation and the definition of  $\tilde{\mu}$ . For the other equation, Thm 2 implies

$$\hat{\tau}^t_{\mu} \circ \pi_{\mu|\omega}(\pi_{\omega}(A)) = \hat{\tau}^t_{\mu} \circ \pi_{\mu}(A)$$

and the first property from Thm 1 gives

$$\hat{\tau}^t_\mu \circ \pi_\mu(A) = \pi_\mu \circ \tau^t(A)$$

Applying these two facts again, we get

$$\pi_{\mu} \circ \tau^{t}(A) = \pi_{\mu|\omega} \circ \pi_{\omega} \circ \tau^{t}(A) = \pi_{\mu|\omega} \circ \hat{\tau}_{\omega}^{t} \circ \pi_{\omega}(A)$$

## References

 Stéphane Attal, Alain Joye and Claude-Alain Pillet. Open Quantum Systems I - The Hamiltonian Approach Springer, 2006.